

Progress in understanding of Indian Ocean circulation, variability, air-sea exchange and impacts on biogeochemistry

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37 **Abstract.** Over the past decade, our understanding of the Indian Ocean has advanced through concerted efforts toward
38 measuring the ocean circulation and air-sea exchanges, detecting changes in water masses, and linking physical processes
39 to ecologically important variables. New circulation pathways and mechanisms have been discovered, which control
40 atmospheric and oceanic mean state and variability. This review brings together new understanding of the ocean-
41 atmosphere system in the Indian Ocean since the last comprehensive review, describing the Indian Ocean circulation
42 patterns, air-sea interactions and climate variability. Coordinated international focus on the Indian Ocean has motivated
43 the application of new technologies to deliver higher-resolution observations and models of Indian Ocean processes. As a
44 result we are discovering the importance of small-scale processes in setting the large-scale gradients and circulation,
45 interactions between physical and biogeochemical processes, interactions between boundary currents and the interior, and
46 between the surface and the deep ocean. A newly discovered regional climate mode in the southeast Indian Ocean, the
47 Ningaloo Niño, has instigated more regional air-sea coupling and marine heatwave research in the global oceans. In the
48 last decade, we have seen rapid warming of the Indian Ocean overlaid with extremes in the form of marine heatwaves.
49 These events have motivated studies that have delivered new insight into the variability in ocean heat content and
50 exchanges in the Indian Ocean, and climate variability on interannual to decadal timescales. This synthesis paper reviews
51 the advances in these areas in the last decade.

52 1. Introduction

53 The physical processes taking place in the Indian Ocean and overlying atmosphere underpin the variability evident in
54 monsoons, extreme events, marine biogeochemical cycles, ecosystems, and ultimately human experience. The Indian
55 Ocean rim countries, accounting for one third of the Earth's human population, depend on this ocean for food and
56 resources, and are dramatically impacted by its variability (Hermes et al., 2019). Increasing our understanding of
57 interactions between geologic, oceanic and atmospheric processes that control the complex physical dynamics of the
58 Indian Ocean region is a priority for many national, bilateral, and international programmes including the Indian Ocean
59 Observing System (IndOOS; Beal et al., 2020), the Climate and Ocean: Variability, Predictability and Change
60 (CLIVAR)/Intergovernmental Oceanographic Commission (IOC) - Indian Ocean Region Panel
61 (https://www.clivar.org/sites/default/files/documents/indian/135_IOP5.pdf), and the second International Indian Ocean
62 Expedition (IIOE-2), to name a few. While initiated through IIOE-2, this review draws on the collective results of all of

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69 the programmes and individual efforts. We focus, in particular, on questions about the Indian Ocean circulation, climate
70 variability and change such as: 1) how have the atmospheric and oceanic circulation of the Indian Ocean changed in the
71 past and how will they change in the future; 2) how do these changes relate to topography and connectivity with the Pacific,
72 Atlantic and Southern oceans; and 3) what impact does the circulation, variability, and change have on biological
73 productivity and fisheries.

74 Recent focus on the Indian Ocean has motivated new international efforts in field campaigns and modelling studies, and
75 leveraged advances in global observations that contribute to the Indian Ocean Observing System (IndOOS; Beal et al.,
76 2020). The Argo profiling float array (Roemmich et al., 2012) reached full coverage in the Indian Ocean in 2006, the
77 RAMA moored buoy array (McPhaden et al., 2009) has now delivered multi-year time series of tropical oceanic and
78 atmospheric variability, with some sites dating back to 2000. Satellite systems continue to provide observations vital to
79 interpreting spatial and temporal variability in the in situ observations, and new technology is now enabling high resolution
80 observations of boundary current variability and small scale processes. Thus, since the reviews of Schott and McCreary
81 (2001) and Schott et al. (2009), the spatial coverage of observations and length of time series have increased substantially
82 such that the signals of many previously unresolved processes are now able to be observed.

83
84 These new higher-resolution observations and companion improvements in model simulations have highlighted the
85 importance of small scale processes in setting the large-scale gradients and circulation, interactions between physical and
86 biogeochemical processes, interactions between boundary currents and the interior, and between the surface and the deep
87 ocean. Overlaid on these interior Indian Ocean processes, ocean warming due to increasing greenhouse gas concentrations
88 has been shown to be pervasive and relentless (Wijffels et al., 2016), and extending to abyssal depths (Johnson et al.,
89 2008a; Desbruyeres et al., 2017).

90
91 The Indian Ocean plays a key role in the global climate system, enabling upwelling of the lower cell of the meridional
92 overturning circulation from abyssal to upper-deep and intermediate waters through diffusive mixing (Schmitz, 1995;
93 Lumpkin and Speer, 2007; McDonagh et al., 2008; Talley, 2013; Hernandez-Guerra and Talley, 2016) and exporting the
94 largest poleward heat flux of all Southern Hemisphere basins (Roxby et al., 2014). In recent decades, the upper 700 m of
95 the entire Indian Ocean has warmed rapidly (Desbruyères et al., 2017). In the southern Indian Ocean, the warming was
96 directly linked primarily to heat advection from a strengthened ITC and, secondly, to a decrease in mean air-sea flux
97 cooling (Li et al., 2017b; Zhang et al., 2018a). This coupling between the ocean and atmosphere in the Indian and Pacific
98 Oceans shifted the balance of global warming, accelerating ocean warming and causing a hiatus in the warming of Earth's
99 surface atmosphere (Section 6). Marine heatwaves have emerged as an increasing threat to marine ecosystems as ocean
100 temperatures warm (e.g. Oliver et al., 2018). Increasingly vulnerable populations need more reliable monsoon predictions,

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112 a task complicated by variability across timescales from intraseasonal to interannual, decadal and beyond in a tightly
113 coupled ocean-atmosphere system (Hazra et al., 2017).

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115 The starting point for this synthesis report are the reviews by Schott and McCreary (2001) and Schott et al. (2009),
116 describing the circulation patterns, air-sea interactions and climate variability on timescales from intraseasonal to
117 interannual, and relatively large spatial scales. We begin with a description of the large scale setting that has been well
118 established since Schott et al. (2009) (Section 2). We then consider the structure and propagation of variability in air-sea
119 interactions at seasonal and intra-seasonal scales, including the contribution of the mesoscale and the ocean's role in air-

120 sea interaction (Section 3). Section 4 discusses new advances in understanding of the upper ocean circulation, organised
121 by region (southern basin, equatorial and northern basin). This section includes an update of the near-surface circulation
122 maps of Talley et al. (2011), including recent work on boundary currents around Australia and Madagascar, and a
123 discussion of the biogeochemical variability observed in each region. The interocean connections with the Pacific, Atlantic
124 and Southern Oceans are discussed in Section 5. Section 6 describes the variability of the Indian Ocean circulation with
125 the recent advances in understanding the warming across the basin, climate modes such as the Indian Ocean Dipole,
126 connection with the El Nino-Southern Oscillation (ENSO), and Indian ocean marine heatwaves. Section 7 focuses on
127 multiscale processes in the Bay of Bengal as an "ocean laboratory", since there have been multiple international programs
128 in this Bay in the last decade. Recent advances from the large scales (>100 km) down to sub-mesoscales (100 m to 10 km)
129 and further down to mixing scales (mm) are discussed. We then link back from mixing to large scales via salinity budgets
130 and coupled phenomena such as the Madden-Julian Oscillation (MJO) to understand the complexity of these processes
131 across multiple scales. We end with a short summary and open questions that will need to be addressed over the next
132 decade.

133 2 Large-scale setting

134 The oceanic and atmospheric circulation of the Indian Ocean are unlike those in the Pacific and Atlantic oceans, largely
135 due to geography. The Asian landmass limits the northern extent of the Indian Ocean to around 25°N so that there is no
136 high-latitude cooling of the ocean, and consequently no dense water formation such as that seen in the North Atlantic and,
137 to some extent, the North Pacific. The intense seasonal variation in temperature over Asia drives the seasonal monsoons:
138 the southwest monsoon in boreal summer, and northeast monsoon in boreal winter. The timing of the onset of the monsoon,
139 and associated wet and dry periods in the Indian Ocean rim countries, varies considerably depending on a range of large-
140 scale climate modes and smaller-scale coupled ocean-atmosphere interactions. The seasonally-reversing winds drive
141 seasonally-reversing ocean currents in the northern Indian Ocean (Section 4.4), e.g. the southwest/northeast monsoon
142 current and the Somali Current. Equatorial currents in the Indian Ocean, eastward near the surface above westward

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147 undercurrents (Section 4.3), provide rapid connection between the western and eastern basin and are also subject to
148 monsoon dynamics.

149 In the southern Indian Ocean (Section 4.2), the connection of the Indian and Pacific Oceans through the Indonesian Seas
150 also contributes to the unique circulation patterns. The very warm and fresh ITF water is funneled into the tropical southern
151 Indian Ocean and carried westward by the South Equatorial Current. The warm, fresh waters are much lighter than those
152 further south, creating a north-south density (pressure) gradient that drives near-surface broad, eastward geostrophic
153 currents between 16°S and 32°S and between Madagascar and Australia (Niiler et al., 2003). This pressure gradient also
154 generates the Leeuwin Current, a unique poleward-flowing eastern boundary current (Godfrey and Ridgway, 1985) that is
155 a downwelling region but is also, counter-intuitively, highly productive (Waite et al., 2007). These two features are not
156 found in the southeastern Atlantic and Pacific oceans. There, the eastern basin currents are characterised by a clear
157 subtropical gyre circulation with weak, equatorward flow and upwelling against the coast.

158 The tropical Indian Ocean (Section 4.3) is home to the largest fraction of sea surface temperature (SST) warmer than 28°C
159 (the tropical warm pool), and is therefore a key region for deep atmospheric convection: the upward part of the Walker
160 Circulation that drives cloud formation and precipitation over the tropical Indo-Pacific. Variation in SST is the primary
161 driver of variation in exchanges between the ocean and atmosphere and is thus a key focus in this paper. Sea surface
162 salinity effects on ocean-atmosphere exchanges have become better understood and are discussed throughout and in
163 particular in Section 7.

164 The tropical Indian Ocean sea surface temperature (SST) has warmed faster over the period 1950–2010 than either the
165 tropical Pacific or Atlantic (Han et al., 2014). Most of this warming (1.2°C over 1901–2012) has occurred in the western
166 Indian Ocean, which has been the largest contributor to the overall global SST trend, with implications for primary
167 productivity (Roxy et al., 2014, 2016). The Indian Ocean accounts for 50–70% of the total ocean heat uptake in the global
168 upper (700 m) ocean, over the last decade, associated with anthropogenic warming (Lee et al., 2015). The deeper ocean
169 (700–2000 m) is warming across the globe with a robust signature of anthropogenic warming evident even in the short
170 Argo record since 2005 (Wijffels et al. 2016, Rathore et al. 2020). Warming in the abyss is detectable and widespread,
171 communicated from the surface of the ocean along pathways from Antarctic Bottom Water formation regions (Purkey and
172 Johnson, 2012). Considerable variability in the Indian Ocean climate system exists on the backdrop of this strong, long-
173 term warming trend.

174 An extensive debate erupted in recent years about whether there was hiatus or a reduced rate of global warming
175 (Lewandowsky et al. 2018). However, persistent cold anomalies in the eastern Pacific have been argued to have enhanced
176 oceanic heat uptake, and the strengthened trade winds are consistent with this argument (Kosaka and Xie 2013, England
177 et al. 2014). It has further been argued that the excess heat taken up by the tropical Pacific has been pumped into the Indian

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Ocean via the Indonesian throughflow (Lee et al. 2015). The tropical Indian Ocean is likely affected by the Southern Ocean trends at a rapid timescale of the order of a decade (Yang et al. 2020), and the Indian Ocean warming may accelerate the Atlantic meridional overturning circulation (Hu et al. 2019) and the Pacific response to anthropogenic forcing (Zhang et al. 2019). Based on these oceanic tunnels and atmospheric bridges into and out of the Indian Ocean, one could hypothesise that the Indian Ocean may be acting as the clearinghouse for oceanic warming under anthropogenic forcing.

Variability in the oceanic and atmospheric circulation of the Indian Ocean is the result of complex interactions that are both internal and external to the Indian Ocean. The recent review of the IndoOS plan (Beal et al., 2019, 2020) summarises the major scientific drivers, of which we still have limited understanding (Fig. 1). The over-arching signal is anthropogenic climate change, causing a background trend of ocean warming and increasing acidity due to uptake of heat and carbon dioxide and affecting the nature of large and small scale variability mechanisms.

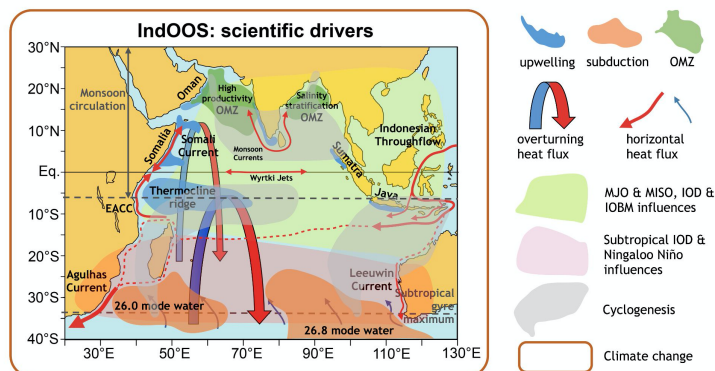


Figure 1: Schematic view of key phenomena in the Indian Ocean (from Beal et al. 2019). The main scientific drivers of the Indian Ocean Observing System, including the Oxygen Minimum Zones (OMZs), upwelling and subduction zones, major heat flux components, the tropical modes of the Madden-Julian Oscillation (MJO), the Monsoon Intra-Seasonal Oscillation (MISO), the Indian Ocean Dipole (IOD) and Indian Ocean Basin Mode (IOBM), the subtropical modes of Ningaloo Niño and subtropical IOD, cyclogenesis, and climate change.

A net poleward flow of heat out of the Indian Ocean is accomplished by a combination of the horizontal circulation along the boundaries, coupled with the Indian Ocean's part of the global meridional overturning circulation (MOC) and shallow

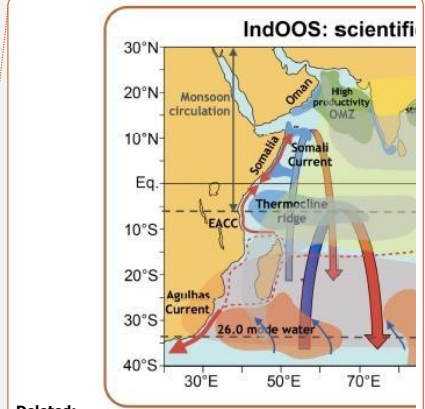
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207 [overturning cells](#). The ITF delivers heat from the Pacific into the Indian Ocean. The Agulhas Current moves heat rapidly
208 southward at surface and intermediate depths (Bryden and Beal, 2001), with 30% of Indian Ocean heat export thought to
209 be carried across 32°S by this gyre circulation (Talley, 2008). The shallow Leeuwin Current makes a smaller direct
210 contribution to the poleward flow of heat (Smith et al., 1991; Feng et al., 2003; Furue et al., 2017) but generates a rich
211 field of mesoscale eddies that carry heat and momentum into the Indian Ocean interior, contributing to heat export across
212 32°S (Domingues et al. 2006, Feng et al., 2007; Dilmahamod et al. 2018).

213 [In the upper ocean, the shallow overturning consists of the cross-equatorial cell \(Miyama et al. 2003; Schott et al. 2004\)](#)
214 [and the subtropical cell \(Schott et al. 2004\). The ascending branches of these cells connect to different upwelling zones in](#)
215 [the southern and northern Indian Ocean and, therefore, play an important role in regulating the climatological mean,](#)
216 [seasonal, and interannual heat balance in the tropical Indian Ocean \(Lee 2004; Lee and McPhaden 2008\).](#) At intermediate
217 depths (500-2000 m), mode waters of varying density enter the Indian Ocean from the Southern Ocean. Along their
218 northward path they mix with lighter waters above, progressively upwelling to the sea surface in a range of locations north
219 of 10°S to then return south in a widespread southward Ekman transport of near-surface waters (Schott et al., 2009). The
220 lower part of the mode water layer mixes with denser waters below and joins the southward flowing deep waters (2000-
221 4000 m). This southward flow also has a contribution from transformed abyssal waters: Antarctic Bottom Water moves
222 northward at abyssal depths, mixing with lighter waters above, progressively upwelling along its path from the Southern
223 Ocean to the Indian Ocean to return southward at shallower depths (Talley, 2013). Cross-equatorial flow is accomplished
224 both at abyssal levels and via the East Africa Coastal Current, seasonally reversing Somali Current (Schott et al., 2009)
225 [and southward Ekman transport \(Schott and McCreary, 2001\).](#)

226 The remaining elements of Fig. 1 refer to oxygen minimum zones (OMZ) in the Arabian Sea and Bay of Bengal and the
227 range of mechanisms that drive strong variations in sea surface temperature leading to shifts in atmospheric convection
228 and precipitation with major effects on rim countries. These mechanisms include: Madden-Julian oscillation (MJO) and
229 Monsoon Intraseasonal Oscillation (MISO), Indian Ocean Dipole (IOD), Indian Ocean Basin Mode, Subtropical IOD, and
230 Ningaloo Niño which are discussed further in Section 6. Cyclogenesis is not discussed in this synthesis. [For discussion of](#)
231 [OMZ, the reader is referred to the review papers of McCreary et al. \(2013\) and Rixen et al. \(2020\).](#)

232 Extreme precipitation in the Bay of Bengal and evaporation in the Red Sea and Arabian Sea lead to strong variability in
233 ocean salinity that in turn impacts ocean circulation and air-sea interaction. The surface salinity gradient in the northern
234 Indian Ocean decreases from the Arabian Sea in the west to the Bay of Bengal in the east. Strong evaporation over the
235 Arabian Sea results in highly saline surface waters (Antonov et al., 2010; Chatterjee et al., 2012), while surface waters in
236 the Bay of Bengal are comparatively fresh and highly stratified as a result of monsoon precipitation and outflow from river
237 systems such as the Ganges-Brahmaputra (Shetye et al., 1996; Vinayachandran et al., 2002). The surface forcing is
238 balanced by the seasonally reversing monsoon currents to maintain the climatological distribution of salinity.

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243 3 Air-sea interactions

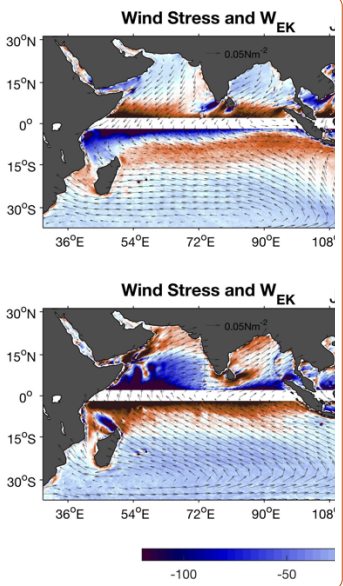
244 The tropical Indian Ocean is highly variable across multiple scales, all of which involve atmosphere-ocean interaction;
245 from the locally intense heat and moisture fluxes that drives tropical cyclones to large-scale convection in the ascending
246 branch of the Hadley circulation, and basin scale ocean heat transport carried by overturning cells that contribute to decadal
247 variability and trends. At intermediate time scales, the intraseasonal oscillations involve strong air-sea coupling (e.g.,
248 Demott et al., 2015). The Indian Ocean Dipole (IOD) is an example of an inherently coupled mode of variability (Saji et
249 al., 1999, Webster et al. 1999, Murtugudde et al. 2000). The monsoonal rainfall around the Indian Ocean is largely fuelled
250 by warm SSTs and strong sea-to-air moisture fluxes. These phenomena emphasise the need to understand the mechanisms
251 of air-sea interaction within the Indian Ocean, with a particular focus on how these processes can be better represented in
252 models to aid predictions of variability in the Earth system.

253 3.1 Seasonal cycle and the monsoons

254 In the open ocean south of 10°S, the wind pattern throughout the year is southeasterly trade winds across the tropics and
255 subtropics and westerlies south of 35°S (Fig. 2). The evaporative cooling of the ocean surface by the trade winds leads to
256 high salinity throughout the subtropics. The curl of the wind stress drives year-round Ekman pumping (downwelling)
257 south of around 15°S (Fig. 2). Downwelling of these denser, high salinity surface waters supplies the downward limb of
258 the shallow Subtropical Cell, STC and Cross-Equatorial Cell, CEC (Schott et al. 2002; Miyama et al. 2003; Schott et al.,
259 2004; Lee 2004; Schott et al., 2009). The subsurface path of the shallow overturning is not well known, and the return to
260 the surface is in any of a number of upwelling zones including the Seychelles-Chagos Thermocline Ridge, for the STC and
261 along Somalia, Oman and the west coast of India for the CEC. North of around 10°S, the winds over the Indian Ocean are
262 characterised by seasonal reversals due to the monsoons (Fig. 2), which in turn cause most of the near-surface currents in
263 these regions to seasonally reverse (Schott et al., 2009; Shankar et al., 2002, Section 4.4).

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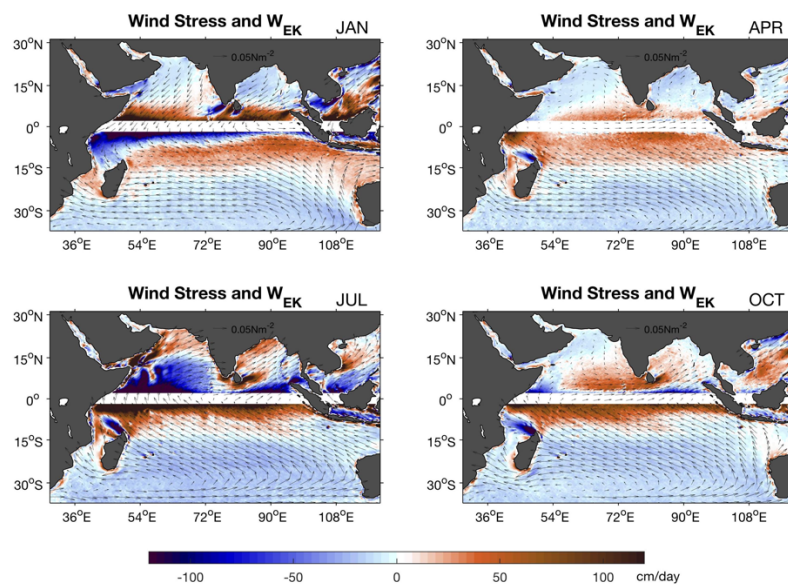


Figure 2: Climatology (2001–2018) of monsoon wind stress (vectors) and Ekman pumping rate (colour shaded) with positive values denoting Ekman suction (upwelling) and negative values Ekman pumping (downwelling) for (a) January - NE monsoon, (b) April, (c) July - SW monsoon, and (d) October. The climatology was constructed by the Objectively Analyzed air-sea Flux High-Resolution (OAFlux-HR) analysis (adapted from Yu 2019).

A strong positive correlation between seasonal net heat fluxes into the ocean and SST variability (Fig. 3) suggests that the seasonal cycle of SST is largely due to the seasonal cycle of winds and cloud cover (Yu et al., 2007). One prominent exception is the Seychelles-Chagos thermocline ridge (located between 5°S and 10°S and east of 50°E), where upwelling and horizontal advection exhibit substantial seasonal variations that in turn contribute to the seasonal cycle of SST (Hermes and Reason, 2008; Foltz et al., 2010). On the equator and to the north, seasonally reversing winds drive complex patterns of upwelling and downwelling that lead to complex SST variability.

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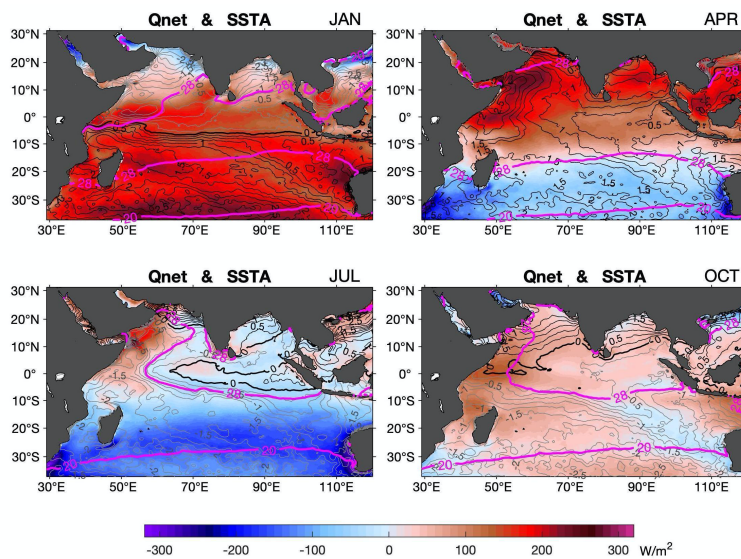
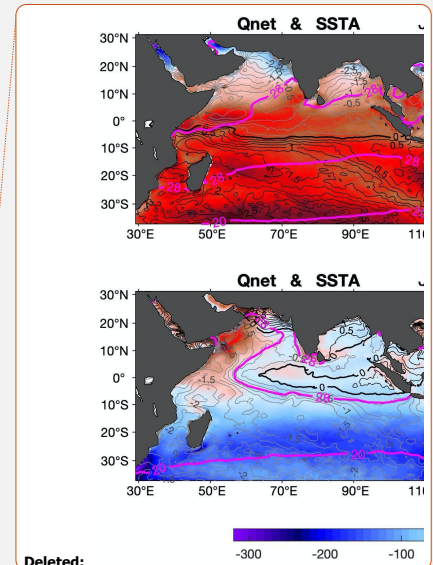


Figure 3: Climatology (2001–2018) of ocean-surface net heat input (colour shaded; positive values denote ocean heat gain and negative values ocean heat loss), SST anomaly (black contours) and 20°C, 28°C SST contours (pink) for (a) January - NE monsoon, (b) April, (c) July - SW monsoon, and (d) October (adapted from Yu 2019). Net heat flux is the sum of solar radiation, longwave radiation, and turbulent latent and sensible heat fluxes. The turbulent heat flux climatology was constructed by the OAFux-HR analysis and surface radiation climatology by the NASA CERES EBAF (Kato et al., 2013).

In the Bay of Bengal and Arabian Sea, surface heat fluxes dominate the seasonal cycle of SST, with the exception of the upwelling zone along the western boundary of the Arabian Sea (Chowdary et al., 2015; Yu et al., 2007). However, salinity effects and subsurface processes (barrier layers, vertical entrainment, variations in the depth of penetration of solar radiation and zonal advection) also influence SST variability (Thangaprakash et al., 2016). Rainfall variability driven by the monsoons creates near-surface salinity variability, most notably in the Bay of Bengal where there is a pronounced annual cycle of sea surface salinity (SSS; Fig. 4, Akhil et al., 2014). Freshwater input at the northern end of the Bay forms a shallow mixed layer stratified by low salinity and is advected southwards along the east coast of India, where it is eventually eroded by vertical mixing (Akhil et al., 2014). The variability in freshwater input contributes to the seasonal



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cycle of barrier layer thickness in the Bay of Bengal (Howden and Murtugudde, 2001; Thadathil et al., 2007), which in turn modulates how strongly SST responds to surface forcing (Li et al., 2017). The seasonally reversing currents that connect the salty Arabian Sea and fresh Bay of Bengal also strongly influence sea surface salinity patterns (Section 4.1.3).

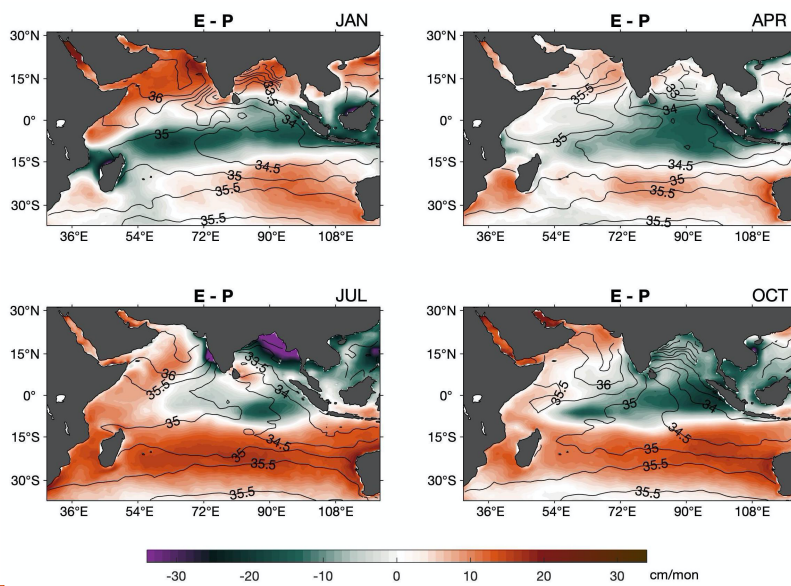
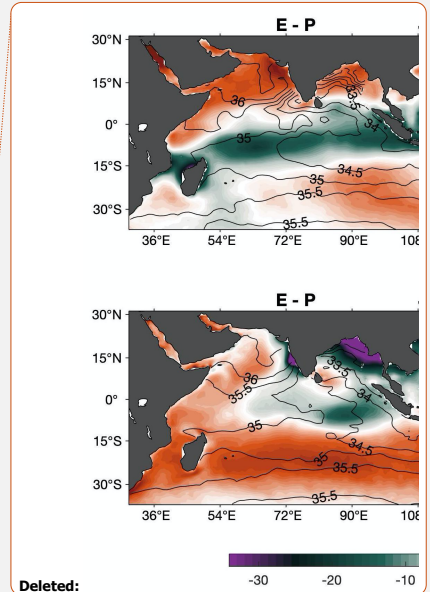


Figure 4: Climatology (2001–2018) of evaporation minus precipitation (colour shaded; positive values denote freshwater leaving the ocean and negative values addition of fresh water to the ocean) and sea surface salinity (black contours) for (a) January - NE monsoon, (b) April, (c) July - SW monsoon, and (d) October (adapted from Yu 2019).

The seasonal cycles in the atmosphere and ocean circulation strongly influence the biological productivity of the near-surface Indian Ocean (Wiggert et al. 2006). Fig. 5 shows the seasonal cycle of satellite chlorophyll *a* and surface currents. The dramatically low productivity in the subtropics, where wind stress curl drives large-scale downwelling (Fig. 2), and highly productive coastal boundaries where wind-driven upwelling occurs, highlights the impact of the circulation and atmosphere-ocean interaction on biological productivity. In turn, the chlorophyll *a* distribution has important implications for air-sea interaction, since higher concentrations of phytoplankton lead to increased absorption of solar radiation (e.g.,



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Morel and Antoine, 1994; Murtugudde et al. 2002; Giddings et al. 2021). Organisation of chlorophyll *a* at intraseasonal timescales has also been reported (Section 3.2.1).

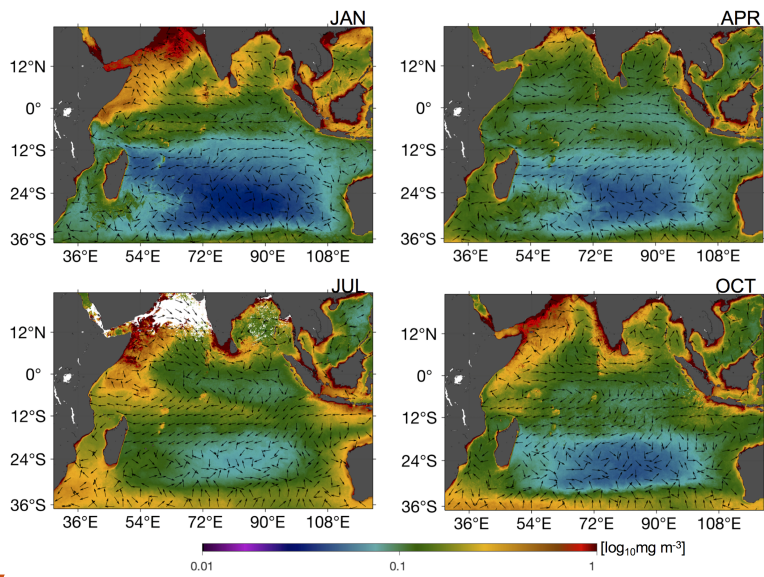
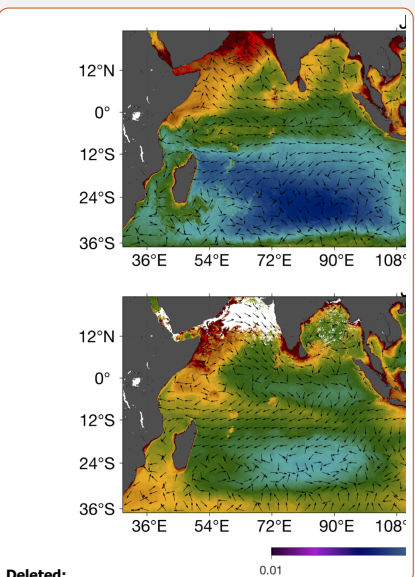


Figure 5: Climatology (2002-2018) of chlorophyll-a concentrations (colormap) and current velocities (arrows) for (a) January (b) April (c) July (d) October. Chlorophyll a climatology was obtained from the MODIS-Aqua product and current velocities were obtained from the third-degree Ocean Surface Current Analysis Real-time (OSCAR) product.

3.2 Intraseasonal air-sea interaction

3.2.1 Madden-Julian Oscillation - MJO

The Madden-Julian Oscillation (MJO; Madden and Julian, 1972, 1971) is the dominant mode of variability in the Indian Ocean at subseasonal time scales. The MJO (Fig. 6) is characterised by eastward-propagating features of enhanced and reduced convection over distances of more than 10,000 km and with a periodicity of around 30–60 days (Zhang, 2005).



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The MJO propagates slowly ($\sim 5 \text{ m s}^{-1}$) through the portion of the Indian and Pacific Oceans where the sea surface is warm, constantly interacting with the underlying ocean and influencing many weather and climate systems. Within the large-scale envelopes of enhanced convection, smaller-scale clusters of clouds propagate westward, and can produce local extremes in rainfall. Air-sea interaction is believed to sustain, and perhaps amplify, the patterns of enhanced and reduced convection as the MJO propagates eastward (Demott et al., 2015). Indo-Pacific warming trends are warping the life cycle of the MJO, which is spending less time over the Indian Ocean, more time over the Pacific and altering mean rainfall trends in parts of the globe (Roxy et al, 2019).

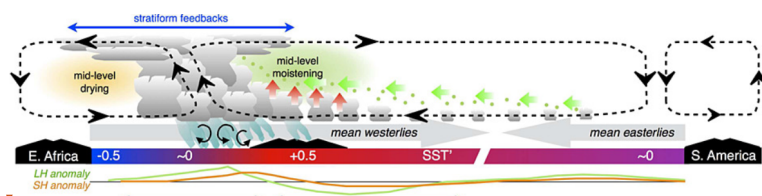
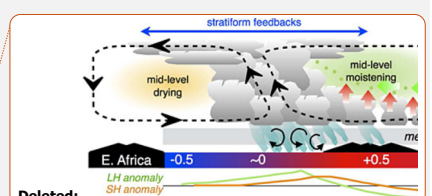


Figure 6: Schematic depiction of Indian and Pacific Ocean feedbacks to the MJO when convection (gray cloud elements) is maximized in the eastern Indian Ocean. Rainfall (aquamarine), circulation anomalies (black dashed cells), convective downdrafts (black rotor arrows), mean winds (faint gray arrows), and moistening by convective detrainment (small green dots) and horizontal and vertical advection (thick green and red arrows, respectively) are overlaid. Net moistening (drying) is shaded green (orange). Positive (red) and negative (blue) SST anomalies for a strong event are shaded, while latent (sensible) heat flux anomalies are shown with green (orange) curves. Central and East Pacific spatial scale is compressed relative to the Warm Pool. Adapted from DeMott et al. (2015).

The MJO-related pattern of winds results in anomalous westerly (easterly) winds to the west (east) of the region of convergence, convection and enhanced rainfall (Fig. 6). These winds generate Kelvin and Rossby waves along the Equator. The Kelvin waves generated by the MJO have been hypothesised (Bergman et al. 2001) to trigger ENSO events in the Pacific. In the Indian Ocean, there is a distinctive sequence of basin-scale ocean waves generated by the MJO. Eastward-propagating equatorial ocean Kelvin waves strike the coast of Sumatra, where they generate coastally-trapped Kelvin waves that propagate northward and southward away from the generation site. Kelvin waves also propagate into the Indonesian seas where they affect the ITF (Pujiana and McPhaden, 2020). Westward-propagating equatorial ocean Rossby waves are also formed, either due to direct intraseasonal wind forcing or through reflection of Kelvin waves at the eastern boundary (Oliver and Thompson, 2010; Webber et al., 2010; Nagura and McPhaden, 2012; Pujiana and McPhaden, 2020). These waves influence local upwelling and currents; they have been linked to variability in coastal currents around the Bay of Bengal (Vialard et al., 2009), to enhancement of the spring Wyrki jets in the eastern equatorial Indian Ocean (Prerna et al., 2019), to changes in subsurface equatorial currents in the central Indian Ocean (Iskandar and McPhaden,

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375 2011) and to changes in upwelling and chlorophyll a concentration in the off-equatorial central Indian Ocean (Webber et
 376 al., 2014). [Such waves also propagate energy downwards into the deep ocean \(e.g., Pujiana and McPhaden, 2020\),](#)
 377 [contributing to deep ocean variability at multiple time scales \(e.g., Matthews et al., 2007\).](#) Downwelling Rossby waves in
 378 the western Indian Ocean create positive SST anomalies through a combination of reduced entrainment of cooler water
 379 from below and zonal advection (Rydbeck et al., 2017; Webber et al., 2012b). These waves therefore act as a triggering
 380 mechanism for new MJO events (Rydbeck and Jensen, 2017; Webber et al., 2010, 2012b, 2012a), and may also play a role
 381 in amplifying existing MJO events.

382 MJO-related winds also lead to variability in mixing within and at the bottom of the mixed layer. Westerly wind bursts
 383 generate zonal currents that create strong vertical current shear (Moum et al., 2014). These currents and the associated
 384 mixing persist after the passage of the atmospheric disturbance. Cooler waters from below the surface are mixed with
 385 surface waters, leading to a reduction in available ocean heat content for the next MJO event and thus reducing its potential
 386 amplitude (Moum et al., 2016). By examining the causes of SST variability in two separate MJO events, McPhaden and
 387 Foltz (2013) showed that the presence or absence of barrier layers may play a crucial role in determining how strongly
 388 mixing and vertical entrainment influence SST. They also found that zonal advection plays a relatively stronger role when
 389 a barrier layer is present. Chi et al. (2014) confirmed the importance of barrier layers in influencing the turbulent heat flux,
 390 but found that thin barrier layers can be eroded by strong current shear that occurs during active phases of the MJO. Wind
 391 mixing and surface heat and freshwater fluxes both contribute in roughly equal proportions to intraseasonal variability in
 392 mixed layer depth (Keerthi et al., 2016).

393 Various studies have investigated the relative importance of surface heat fluxes and subsurface ocean processes for the
 394 evolution of SST at intraseasonal time scales. [The Seychelles-Chagos Thermocline Ridge \(SCTR\), is a region of high](#)
 395 [intraseasonal SST variability \(Saji et al. 2006, Hermes and Reason, 2008\).](#) Several observational studies have concluded
 396 that the SST variability [here](#) is predominantly generated by variability in surface heat fluxes (Jayakumar et al., 2011;
 397 Vialard et al., 2008), while Drushka et al. (2012) suggest this finding applies across most of the tropical Indian Ocean.
 398 Such studies, however, typically exhibit large uncertainty in the subsurface ocean terms. [The shallow thermocline and](#)
 399 [strong high frequency winds in the SCTR region enhance near-inertial waves and lead to strong mixing at the base of the](#)
 400 [mixed layer as well as in the thermocline \(e.g. Cuypers et al. 2013; Sabu et al. 2021\).](#) Modelling studies have shown that
 401 ocean dynamics play an important role in generating SST variability (Halkides et al., 2015; Han et al., 2007). For example,
 402 Fig. 7 from the study of Halkides et al. (2015) shows the relative contribution of modelled ocean dynamical processes and
 403 thermodynamical processes (i.e., surface heat fluxes) in forcing intraseasonal SST variability. Fig. 7a shows that ocean
 404 dynamical processes (green shading), including horizontal and vertical advection, are the dominant source of intraseasonal
 405 SST variability on the equator and in upwelling regions off Indonesia, Sri Lanka and along the western boundary. The
 406 ocean dynamical processes are in turn dominated by horizontal advection along the equator and tropical coastlines (Fig.
 407 7b, pink shading), and vertical advection (blue shading) in the off-equatorial ocean interior.

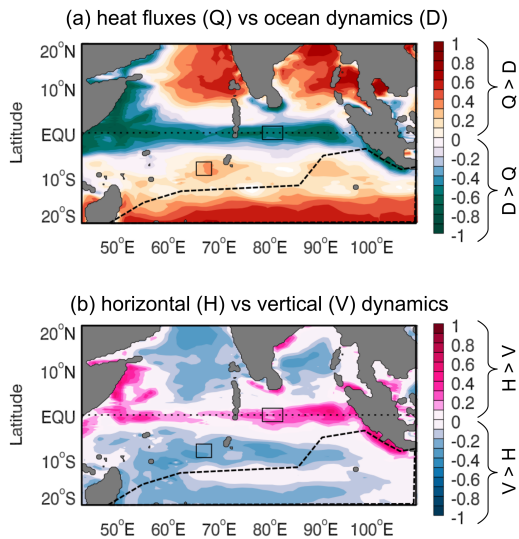
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410 **Figure 7: Modelled balance of processes driving intraseasonal SST variability. (a) Relative role of heat fluxes (Q)**
 411 **and ocean dynamics (D) in driving SST variability, with red (green) colours implying Q (D) dominates forcing, (b)**
 412 **Relative role of horizontal (H) and vertical (V) processes in the dynamical forcing, with pink (blue) colours implying**
 413 **that H (V) processes dominate. All fields are derived from the ECCO-JPL ocean state estimate. The dotted line**
 414 **marks the Equator, dashed line in the southern hemisphere outlines a region in which the model does not fully**
 415 **resolve the ocean heat budget, and the boxes on the Equator and at 10°S mark regions for further analysis not**
 416 **described here. Modified from Halkides et al. (2015).**

417 Organisation of chlorophyll a at intraseasonal timescales has also been reported, with model studies indicating potential
 418 biophysical feedbacks due to the variability of penetrative radiation into the water column (Waliser et al. 2005, Jin et al.
 419 2013a; Giddings et al., 2021). In the Bay of Bengal, the proportion of incoming solar radiation absorbed within the mixed
 420 layer varies between 60% and 97% due to a combination of variability in chlorophyll a concentration and mixed layer
 421 depth (Lotliker et al., 2016) and an increase in chlorophyll of 0.3 mg/m³ can lead to SST increase of up to 0.35°C on
 422 intraseasonal time scales (Giddings et al., 2021). Representing the seasonal cycle of chlorophyll a concentration in the
 423 Arabian Sea in a coupled model led to substantial changes in the simulated SST and monsoon rainfall over India (Turner

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426 et al., 2012), suggesting that incorporating this process into coupled models may be important to improve simulation of
427 monsoon rainfall and circulation around the Indian Ocean.

428 Figure 8 illustrates propagation of surface patterns in an MJO composite constructed by Jin et al. (2013a). In each panel
429 the peak in outgoing longwave radiation (OLR, a proxy for convection) is indicated by a red diagonal line. The MJO
430 generates substantial surface heat flux anomalies that create a pattern of surface heat fluxes and SST anomalies such that
431 warm (cool) SSTs lead enhanced (reduced) convection by a quarter of a phase (e.g., Shinoda et al., 1998). The MJO also
432 leads to low-frequency rectifications in the mean state of physical and ecosystem responses (Fig. 8, Waliser et al. 2003,
433 Jin et al. [2013a,b](#)), in particular semi-annual variability can rectify into mean flows along the equator (Nagura and
434 McPhaden, 2014).

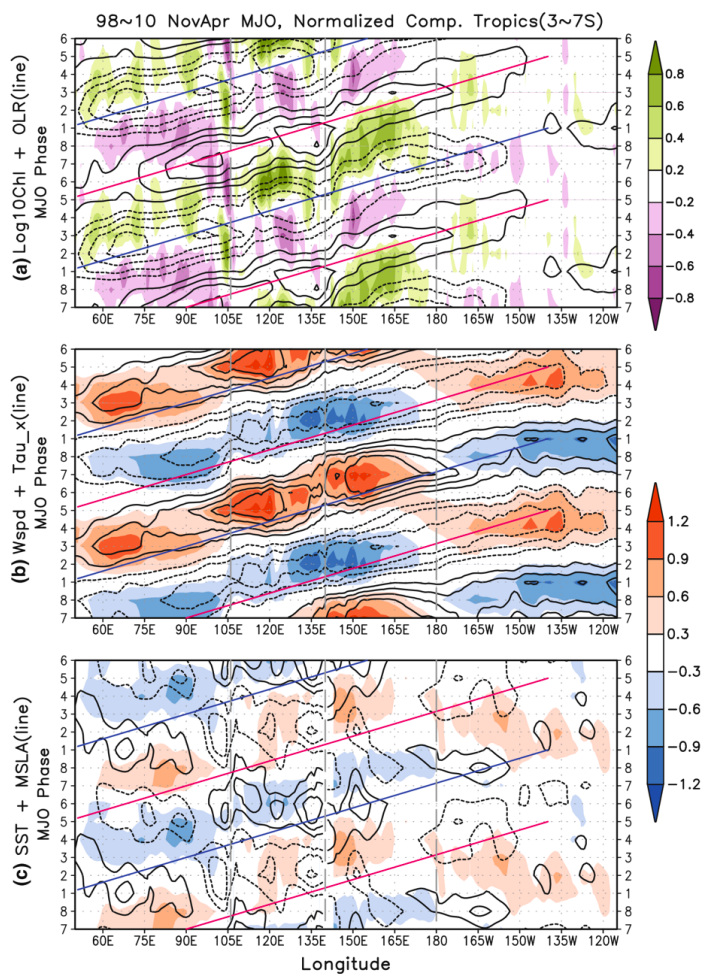
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Figure 8: MJO composite evolution for the Boreal winter (Nov-Apr) averaged over latitudes 3°–7°S for the period of November 1st, 1997 to October 31st, 2010, of a) Log₁₀Chl from SeaWIFS satellite observations (shaded) and satellite-derived outgoing longwave radiation (contour), b) wind speed (shaded) and zonal wind stress (contour), both from the cross-calibrated multi-platform (CCMP) dataset, and c) NOAA-OI satellite SST anomalies (shaded) and AVISO mean sea level anomaly (contour). All contour intervals match shading levels in (c), and solid (dash) line indicates positive (negative) values. All variables are normalised, and the same MJO composite is repeated for two cycles for convenience. There are between 127–227 events in the composite for each MJO phase. Red diagonal lines indicate peak signals of positive OLR, and blue lines indicate negative OLR peak, so these are guides for the MJO propagation. The relative location of each propagation line in all panels is the same. Left and center gray vertical dashed line indicates the western and eastern boundary of the Maritime Continent, and the right gray line is on the Dateline, where chlorophyll a propagation stops. From Jin et al. (2013a).

3.2.2 Monsoon Intraseasonal Oscillation - MISO

While the MJO dominates intraseasonal variability during October to April, during May to September (boreal summer, southwest monsoon), the Monsoon Intraseasonal Oscillation (MISO; Goswami, 2012; Suhas et al., 2013) dominates. MISOs can be seen as low pressure systems laden with moisture which deliver rain from atmospheric instabilities (Fig. 9). The MISO is also known as the Boreal Summer Intraseasonal Oscillation (BSISO; Lau and Waliser, 2012; Lee et al., 2013). The MISO oscillations are dynamically linked to the equatorial MJO (e.g., Sperber and Annamalai, 2008), but exhibit northeastward and northwestward propagating features, with the main centre of action being the Bay of Bengal. These northward-propagating bands of enhanced and reduced rainfall exhibit a similar relationship with SST to the MJO: warm SST leading increased rainfall (cool SST leading reduced rainfall) that then determine the wet/dry (or active/break) cycles of the South Asian monsoon (Vecchi and Harrison, 2002; Roxy et al., 2013; Suhas et al., 2013; Zhang et al., 2018). These SST anomalies are primarily forced by variations in surface heat fluxes in the Bay of Bengal (Girishkumar et al., 2017; Vialard et al., 2012), while variations in wind-induced mixing, Ekman pumping and entrainment drive SST variability in the Arabian Sea (Duncan and Han, 2012; Vialard et al., 2012).

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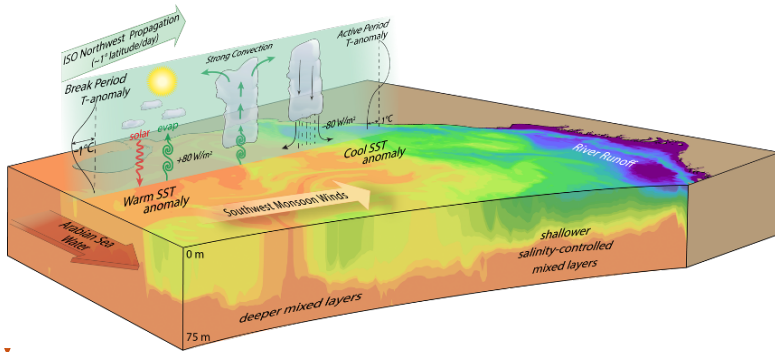
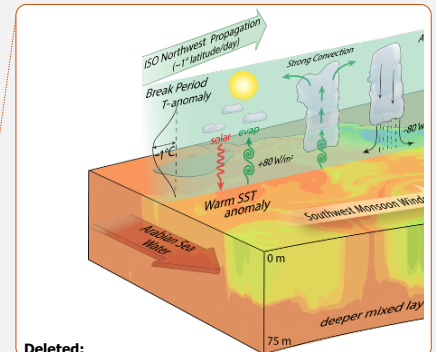


Figure 9: A schematic of the Monsoon Intraseasonal Oscillation (MISO) in the Bay of Bengal, showing the coupled ocean-atmosphere 30–60 day mode northwestward propagation and associated processes in the atmosphere and the ocean. (From Mahadevan et al., 2016a).

Simulations of the MISO are still generally poor in state-of-the-art coupled models (e.g., Goswami et al., 2013; Jayakumar et al., 2017; Sabeerali et al., 2013; Sharmila et al., 2013) and re-analysis products (e.g. Sanchez-Franks et al., 2018). Evidence exists from observations of low-level convergence and OLR, as well as from forced atmospheric and coupled ocean-atmosphere model experiments, that both MJOs and MISOs are phenomena that require coupling between the ocean and atmosphere to exist. This is even though the scales of SST anomalies tend to be an order of magnitude smaller than the scales of the propagating atmospheric systems (Waliser et al., 1999; Zhou and Murtugudde, 2009). Including air-sea coupling in simulations of the MISO has been identified as key to improving simulation of this oscillation in some models (e.g., Jayakumar et al., 2017; Li et al., 2018; Roxy et al., 2013; Sharmila et al., 2013), and has been shown to improve aspects of simulation in others (e.g., Bellon et al., 2008; Peatman and Klingaman, 2018).

While new theories continue to be proposed for MJOs (e.g., Wang et al., 2016), MISOs have not received similar attention likely due to their more local nature compared to the global impacts of MJOs (e.g. their impact on ENSO). The mechanism that causes the northward propagation of the MISO is still a topic of research. The most recent theory for MISOs proposed by Zhou et al. (2017a, b) invokes an explicit coupling between the ocean and the atmosphere in a so-called Central Indian Ocean mode. Zonal winds at intraseasonal timescales over the Indian Ocean are argued to be coupled to SSTs to produce a barotropic instability in the meridional gradient of the zonal winds. The horizontal atmospheric eddy fluxes generated by the barotropic instability are invoked to explain the northward propagation and the advection of momentum and moisture as a coupled phenomenon. Key questions remain about the oceanic and air-sea interaction processes that



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reorganise the SSTs in the Central Indian Ocean mode as well as the respective roles of the vertical and horizontal shears in driving northward propagation of MISOs.

Observations and models indicate that MISOs may be slowing down because of the warming in the Indian Ocean (Sabeerali et al., 2013), which needs to be understood better for providing reliable monsoon predictions and projections in this climate vulnerable region. This is underscored by the observational evidence that climate variability and change are increasing the frequency of dry spells and the intensity of wet spells in the Indian summer monsoon, which are directly related to MISO (Singh et al., 2014).

3.2.3 Intraseasonal drivers of heavy rainfall

As the MJO season begins to wind down in April, northward propagating MISOs begin to become dominant in the northern Indian Ocean, north of around 5°N. While the southwesterlies produce some of the strongest coastal upwelling off Somalia and cool the Arabian Sea, the Bay of Bengal remains warm and largely above the convective threshold (28°C) owing to the freshwater input from rainfall as well as rivers discharging into the Bay (Roxy and Tanimoto, 2007). The freshwater input creates a shallow density stratification (barrier layer) within the temperature mixed layer and thereby weakens the upwelling of cold water from the thermocline. MISOs deliver rain from atmospheric instabilities, but what controls the rainfall at intraseasonal timescales during the summer can be expected to be region specific with moisture supply determining the rainfall variability over land (Pathak et al., 2017).

Over the ocean, the largely evaporative Arabian Sea is relatively cool but the southwesterlies begin to slow down as they approach the Western Ghats mountain range on the west coast of India, leading to maximum rainfall there during the boreal summer monsoon season (Xi et al., 2015). Rather counterintuitively, the warm SST in the Bay of Bengal remains above the convective threshold (Gadgil et al., 1984; Roxy, 2013) and yet, the ocean is not in direct control of the intraseasonal rainfall events. Once the SSTs are warm enough to support atmospheric convection, it is baroclinic instabilities, and not static instabilities induced by warm SSTs, that drive the majority of rainfall over the Bay of Bengal (Xi et al., 2015). These findings should provide a proper paradigm for understanding the role of SST in monsoon and MISO in terms of the ocean dynamics and air-sea interaction processes that matter most.

3.3 Ocean internal variability impacts on air-sea interaction

Mesoscale eddies are ubiquitous in the ocean. In the tropical Indian Ocean, however, linear dynamics dominate and the impacts of eddies are (or seem) small. While the Indian Ocean has the largest SST variability occurring at seasonal timescales, strong mesoscale variability is also observed along the Somali coast where the western boundary current crosses the Equator. The slope of the East African coastline and the equatorial crossing of the low-latitude jet produce multiple eddies (Nof and Olson, 1993), which are shown to generate strong air-sea coupling at mesoscales (Schott and

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535 McCreary, 2001; Schott et al., 2009; Vecchi et al., 2004; Seo et al., 2008). Some intraseasonal oscillations in the ocean
536 were reported in the southwestern tropical Indian Ocean (Kindle and Thompson, 1989) but generally, the impact of ocean
537 internal variability on SSTs in the ~~tropical Indian Ocean has not been widely studied. At the eastern boundary of the~~
538 ~~subtropical Indian Ocean, instability of the poleward Leeuwin Current generates a rich field of mesoscale eddies that carry~~
539 ~~heat into the Indian Ocean interior, contributing to air-sea exchange of heat and the oceanic interior poleward heat transport~~
540 ~~(Domingues et al. 2006, Feng et al., 2007; Dilmahamod et al. 2018). In the subtropical southeast Indian Ocean, mesoscale~~
541 ~~eddies, and possibly annual and semiannual Rossby waves propagating from the eastern boundary, were found to influence~~
542 ~~the seasonal variation of the surface layer heat balance through horizontal advection (Cyriac et al. 2019).~~

543 Low-frequency internal variability is also possible. Jochum and Murtugudde (2004) performed forced ocean model
544 experiments with climatological forcing alone to demonstrate that significant low-frequency variability at interannual
545 timescales is generated in the Indian Ocean ~~by~~ mesoscale eddies and other types of nonlinearity. The role of internal
546 variability in regional coupled climate variability as well as ecosystem and biogeochemistry remain interesting problems
547 for this already warm ocean, ~~in which~~ even small SST anomalies can be important for generating large-scale ocean
548 atmosphere interactions (Palmer and Mansfield, 1994).

549 **4 Upper Ocean Circulation and Biogeochemical Variability**

550 **4.1 Overview**

551 The near surface circulation in the Indian Ocean consists of the monsoon-dominated, seasonally reversing currents north
552 of around 10°S, and the steady currents to the south, as illustrated in Fig. 10a for the southwest monsoon (July-August)
553 and Fig. 10b for the northeast monsoon (January-February). This figure has been updated from Talley et al. (2011) to
554 recognise recent advances in understanding of circulation patterns. In the northern Indian Ocean, additions are a revision
555 of the Red Sea circulation (Menezes et al., 2019). In the southern Indian Ocean, moving in an anti-clockwise direction
556 from the Maritime Continent, additions are: 1) seasonally reversing flows in the Java Sea; 2) the Holloway Current along
557 Australia's Northwest Shelf (Holloway and Nye, 1985; Holloway, 1995; Brahmanpour et al., 2016); 3) revised position
558 of the ~~salinity-driven Eastern Gyrar Current that flows eastward from around 90°E along approximately 15°S, recirculating~~
559 Indonesian Throughflow Water from the South Equatorial Current ~~and supplying~~ the poleward-flowing Leeuwin Current
560 (Meyers et al., 1995; Domingues et al., 2007; Menezes et al., 2013, 2014); 4) the ~~near-surface~~ South Indian Countercurrent
561 with 3 distinct branches, northern, central and southern, flowing from the southern tip of Madagascar to Australia where
562 they merge with the poleward-flowing Leeuwin Current (Menezes et al., 2014 ~~and references therein~~); and 5) the splitting
563 of the Flinders Current near 110°E, with one branch recirculating back toward Australia, and the other a westward
564 continuation of the Flinders Current, previously not shown (Duran et al., 2020).

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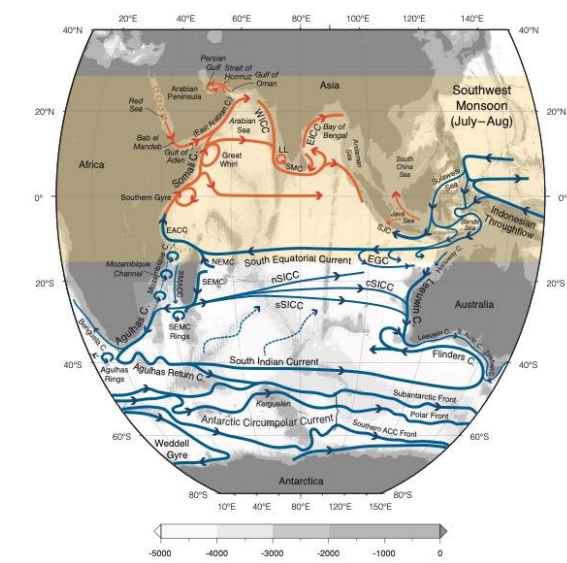
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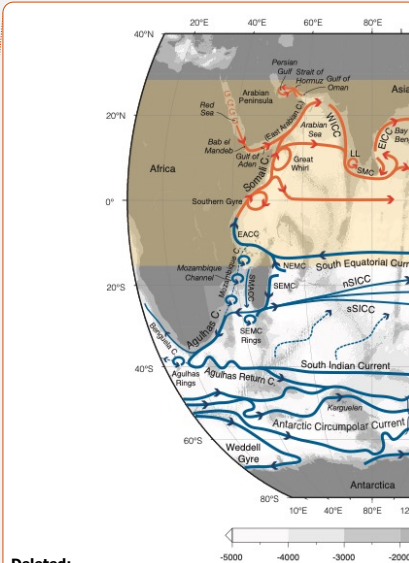
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573 **Figure 10a: Schematic near-surface circulation during the Southwest Monsoon (July-August). Blue: year-round**
574 **mean flows with no seasonal reversals. Orange: monsoonally reversing circulation (after Schott & McCreary,**
575 **2001). The ACC fronts are taken directly from Orsi, Whitworth, and Nowlin (1995). Acronyms: EACC, East**
576 **African Coastal Current; NEMC, Northeast Madagascar Current; SEMC, Southeast Madagascar Current;**
577 **SMACC, Southwest MADagascar Coastal Current; WICC, West Indian Coastal Current; EICC, East Indian**
578 **Coastal Current; LH and LL, Lakshwadeep high and low; SJC, South Java Current; EGC, Eastern Gyral Current;**
579 **SICC, South Indian Countercurrent (south, central and southern branches); NEC, Northeast Monsoon Current.**
580 **Updated from Talley et al. (2011), originally based on Schott and McCreary (2001). The light gray shading shows**
581 **seafloor bathymetry.**



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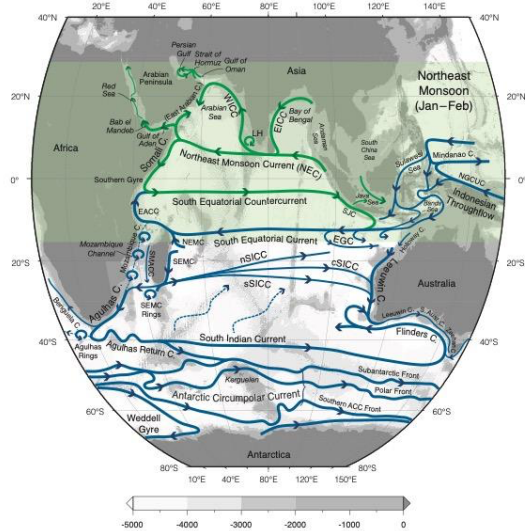


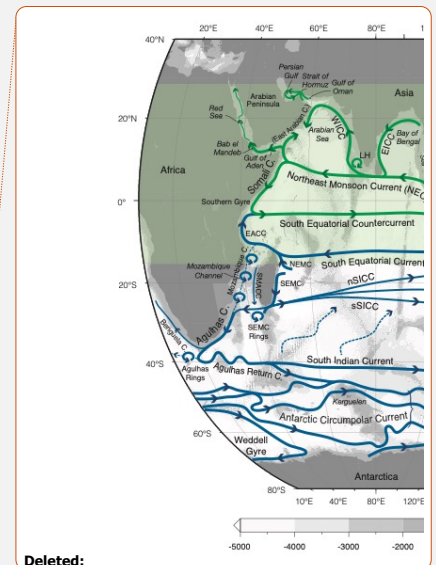
Figure 10b: Schematic near-surface circulation during the Northeast Monsoon (January-February). Details as for Fig. 10a.

The intermediate and deep circulation and overturning cells will not be examined in this synthesis. The reader is referred to Talley et al. (2011) and references therein, and in addition, Nagura and McPhaden (2018) who used Argo and CTD data to map out the circulation and water masses in density classes associated with the shallow overturning circulation, with emphasis on the southern hemisphere.

4.2 Southern Indian Ocean

4.2.1 South Equatorial Current

The South Equatorial Current (SEC), the northern limb of the southern Indian Ocean subtropical gyre, carries Indonesian Throughflow (ITF) waters into the interior Indian Ocean, flowing westward between 10–20°S (Fig. 10a and 10b). Upon reaching the northern tip of eastern Madagascar, it bifurcates and supplies the Northeast Madagascar Current (NEMC);



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Schott and McCreary, 2001; Song et al., 2004; Valsala and Ikeda, 2007) and the Southeast Madagascar Current (SEMC) and contributes to the development of Mozambique Channel eddies. The mean flow through the Mozambique Channel is weak (Song et al., 2004), although there is an indication from ocean model results that the eddy-dominated flow contributes on the order of 20 Sv southward (Durgadoo et al., 2013). The Mozambique Channel eddies, eddies from the SEMC and recirculation combine to feed into the Agulhas Current (Schott and McCreary, 2001).

Between 50 and 80°E the SEC is coincident with the southern half of the Seychelles-Chagos Thermocline Ridge (SCTR, Vialard et al., 2009). The SCTR is characterized by a relatively shallow thermocline and thin mixed layer (~30m) across the southern tropical Indian Ocean in the latitude band 5-15°S. Between 50 and 80°E the SCTR/SEC is a region of significant upwelling (Hermes and Reason, 2008; Vialard et al., 2009; Resplandy et al., 2009; Dilmahamod, 2014), which affects biogeochemistry, and even fisheries (Resplandy et al., 2009; Robinson et al., 2010; Dilmahamod, 2014).

In the eastern IO, the intraseasonal variation of the SEC is mostly attributed to the baroclinic instability of the mean current (Feng and Wijffels, 2002), which is important for the meridional heat transport in the region and contributes to the demise of Indian Ocean Dipole events (Ogata and Masumoto, 2011; Yang et al. 2015). Barotropic instability of the SEC has also been proposed to be a key mechanism for generating intraseasonal variability (Yu and Potemra, 2006). These intraseasonal signals propagate westward as Rossby Waves, influencing the SEC variability in the western Indian Ocean (Zhou and Murtugudde, 2008).

Interannual variability in the ITF due to ENSO, IOD and other influences is communicated into the interior Indian Ocean along the SEC and via Kelvin and Rossby waves (Godfrey, 1989, 1996; Meyers et al., 1995; Meyers, 1996; Wijffels and Meyers, 2004). Pressure anomalies associated with ENSO and IOD are communicated through the Indonesian seas as Kelvin and Rossby waves. These anomalies propagate westward into the Indian Ocean as Rossby waves. At the same time the pressure anomalies drive variations in ITF and SEC transport and induce temperature/salinity variability via advection. Geostrophic transport variability in the long-time repeat XBT line IX1 shows that the SEC is stronger during La Niña and positive Indian Ocean Dipole events (Meyers, 1996; Liu et al., 2015). Similarly, the Pacific Decadal Oscillation alters the SEC and ITF transports and associated water properties (Section 6.1). During the climate change hiatus period of 2000-2011, the enhanced heat transport of the SEC/ITF was a key mechanism for the fast warming trend in the southern subtropical Indian Ocean (Section 6.1).

4.2.2 Western Boundary

The Agulhas Current (Fig. 10) has long been known as one of the strongest western boundary currents in the global oceans, with an average transport of 75 Sverdrups and current speeds in excess of 2 m s^{-1} (Beal et al., 2015; Beal et al., 2011).

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Moved down [10]: Vertical sections of the SEC/SCTR region

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811 The Agulhas Current plays a vital role in the global thermohaline circulation, advecting warm, salty, subtropical water
 812 southwards, following the continental shelf of South Africa and meandering less than 150 km offshore (Gründlingh, 1983;
 813 Lutjeharms 2006). The strength and warmth of the Agulhas Current influences atmospheric storm tracks and storm
 814 development. The large moisture source of the warm Agulhas Current region contributes significantly to the frequency
 815 and strength of African precipitation, which significantly impacts rain-fed subsistence farming (Hermes et al. 2019 and
 816 references therein).

817 South of the tip of Africa, the Agulhas Current retroflects eastwards into the South Indian Ocean (Fig. 10). This
 818 retroflection area is highly variable, occluding rings that propagate into the South Atlantic Ocean. The Agulhas variability
 819 is linked upstream to modes of variability including ENSO (Elipot and Beal, 2018, Trott et al., 2021) and downstream
 820 with the Atlantic meridional overturning circulation, providing an essential link between the Pacific, Indian and Atlantic
 821 Oceans (Beal et al., 2011). Estimates of the rate of mass and heat exchange carried by Agulhas leakage south of Africa
 822 (and the number of rings shed per year) vary and are difficult to verify reliably (Weijer et al., 2014). Daher et al (2020)
 823 recently used a combination of drifters and Argo floats to derive an estimate of Agulhas leakage of 20 Sv. van Sebille et
 824 al. (2011) and le Bars et al. (2014) suggested upstream variability of the Agulhas Current has an effect on inter-ocean
 825 exchange between the South Indian and South Atlantic oceans, primarily by influencing the frequency of ring shedding at
 826 the Agulhas retroflection. However, a few recent papers suggest instead that its variability is driven by the Southern
 827 Hemisphere Westerlies (Durgadoo et al. 2013; Loveday et al., 2014; Elipot and Beal, 2015).

829 The Agulhas Current has a seasonal cycle and is strongest in summer (Krug and Tournadre, 2012; Beal and Elipot, 2016)
 830 and tied to a baroclinic adjustment of near-field winds (Hutchinson et al. 2018). Seasonal changes in the Agulhas
 831 retroflection region (Lutjeharms and van Ballegooyen, 1988; Quartly and Srokosz, 1993) and in the southwest Indian
 832 Ocean (Ffield et al., 1997) have been suggested from hydrographic and satellite data (Krug et al., 2012), but with weak
 833 statistical significance due to a lack of sufficiently long time series.

835 Although long term observations in this region are limited there are numerous recent studies that have further elucidated
 836 our understanding of the Agulhas Current. Beal and Elipot (2016) used 3 years of in situ data to show that, contrary to
 837 expectations, the Agulhas Current has not intensified since the early 1990s. Instead, it has broadened as a result of more
 838 eddy activity, driven by intensifying winds. Variability in the path and strength of the Agulhas Current has mostly been
 839 attributed to solitary Agulhas meanders within the Current system (also known as Natal pulses) which drive upwelling
 840 and cross-shelf transports, affecting marine productivity, fisheries and recruitment over the Agulhas Bank (Beal and
 841 Bryden, 1999; Roberts et al., 2010, Elipot and Beal, 2015). Recent work has highlighted the importance of submesoscale
 842 eddies in the Agulhas Current frontal region driving an inshore edge flow reversal which can have important consequences
 843 on fisheries (Krug et al., 2017).

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854
855 The advance in models has also helped improve our understanding of the Agulhas Current, which is generally not well
856 represented in global ocean models. Hutchinson et al. (2018) used idealized models to expose a link between the
857 seasonality of the Agulhas Current and propagation of first baroclinic mode Rossby waves communicating the wind stress
858 signal across the western portion of the Southern Indian Ocean, with the signal from winds further east having little effect.

859 4.2.3 Interior flows

860 In the central-eastern South Indian Ocean between 20°S and 30°S, the surface geostrophic flow is generally eastward,
861 opposite to the prediction of both the Ekman and Sverdrup theories (Sharma 1976; Sharma et al., 1978; Godfrey and
862 Ridgway, 1985; Schott et al., 2009). This flow is driven by the large-scale, poleward drop in the dynamic height (steric
863 height) near the sea surface (Godfrey and Ridgway, 1985; Schott et al., 2009) related to the meridional transition from the
864 very fresh and warm SEC waters to the increasingly cooler, saltier and denser waters to the south. The flow generally
865 extends from the sea surface to ~200–300 m (Domingues et al., 2007; Palastanga et al., 2007; Divakaran and Brassington,
866 2011; Menezes et al., 2014). The mechanisms that determine the vertical extent of the interior eastward flow remains
867 unclear, although this depth coincides with the depth of the shelf break at the eastern boundary and the bottom of the
868 Leeuwin Current along that boundary. This correspondence may be achieved by the westward propagation of baroclinic
869 Rossby waves (Weaver and Middleton, 1989; Furue et al., 2013). Below the near-surface eastward flows, the flow is
870 weakly westward (Domingues et al., 2007; Schott et al., 2009; Furue et al., 2017).

871
872 Embedded in this general eastward flow are narrower eastward jets (Maximenko et al., 2009; Divakaran and Brassington,
873 2011; Menezes et al., 2014), collectively known as the South Indian (Ocean) Countercurrent (SICC; Palastanga et al. 2007;
874 Siedler et al. 2006; Menezes et al., 2014). They start out as a single jet emanating from the southern tip of Madagascar
875 around 25°S, possibly fed by a partial retroflection of the SEMC (Palastanga et al., 2007; Siedler et al., 2006, 2009) and
876 divide into separate jets around the Central Indian Ridge (65°E–68°E) (Menezes et al., 2014). Eastward flows exist in
877 similar latitude bands in the North and South Pacific and North and South Atlantic (Yoshida and Kidokoro, 1967; Merle
878 et al., 1969; Takeuchi, 1984; Kubokawa, 1999; Qiu and Chen, 2004; Kobashi and Kubokawa, 2012). However, the jets in
879 these basins are weaker and shallower than the SICC and do not extend all the way to the eastern boundary (Menezes,
880 2015).

881
882 Three main jets (Fig. 10a) are evident in geostrophic velocity calculated from both altimetric sea surface height and
883 hydrography and are captured in OGCMs (Maximenko et al., 2009; Divakaran and Brassington, 2011; Menezes et al.,
884 2014). The stronger southern jet (3–4 Sv) crosses the basin around 26°S and has an associated thermal front at depths
885 around 100–200 m (Sharma 1976; Siedler et al., 2006; Menezes et al., 2014; Palastanga et al., 2007). This front suggests
886 that the southern SICC has physics similar to the Subtropical Countercurrents (STCCs) of the Pacific Ocean (Kubokawa,

Moved down [12]: Chlorophyll *a* concentrations and production rates in Mozambique Channel surface waters are generally low (< 0.4 mg/m³ and < 700 gC m² d⁻¹, Fig. 5), and not significantly different in cyclonic and anticyclonic eddies (Lamont et al., 2014; Barlow et al., 2014;

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Moved down [14]: observed between 25 and 125 m depth depending on the proximity to the shelf and the influence of mesoscale eddies (Barlow et al., 2014; Lamont et al., 2014). Eddies in the Mozambique Channel also have a strong influence on the lateral transport of nutrients and chlorophyll from the coasts of Madagascar and Africa. Indeed, enhanced phytoplankton production within both cyclonic and anticyclonic eddies in the Mozambique Channel often occurs in response to lateral nutrient inputs into the euphotic zone by horizontal advection from the coasts of Madagascar and Africa rather than

Moved down [15]: eddy induced upwelling and downwelling (José et al., 2014; Lamont et al., 2014; Roberts et al., 2014). In

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Moved down [16]: 2014). In contrast, in the Southeast Madagascar Current, topographically-induced coastal upwelling brings cold, nutrient-rich water up to the surface, which supports high rates of primary production (Lutjeharms and Machu, 2000; Ho et al., 2004; Quartly and Srokosz, 2004). This upwelling and its impacts are observed in both the austral summer and winter (Ho et al., 2004).

The Agulhas Current itself is warm and oligotrophic with sources derived from low nutrient and low chlorophyll surface waters from the Mozambique Channel, Southeast Madagascar Current and the southwestern tropical Indian Ocean (Fig. 5; Lutjeharms, 2006). Chlorophyll *a* concentrations and production rates in Agulhas Current surface waters are particularly low during austral summer (< 0.2

Moved down [17]: 5, 6 and 20 in Hood et al., 2017). The Agulhas Current can drive upwelling and elevate primary production

Moved down [18]: 5, 6 and 20 in Hood et al., 2017). This enhancement is most pronounced in austral summer and further

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Deleted: favourable (easterly) winds and the aforementioned topographically-induced upwelling.

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1999; Kobashi and Kubokawa, 2012, Menezes et al., 2014). The location and strength of the SICC vary between studies, from well-defined jets (Siedler et al. 2006, Palastanga et al. 2007, Divakaran and Brassington 2011, Menezes et al. 2014) to a mean velocity structure (Jia et al., 2011a), or even absence of the SICC (Srokosz et al. 2015). Depending on the region and time in which its characteristics were determined, the SICC varies from a weak mean current of 2–3 cm/s (Jia et al., 2011a) to a strong jet of 50 cm/s eastward flow (Siedler et al., 2006).

The eastward flowing Eastern Gyral Current (EGC) is part of an anticyclonic recirculation centred at the Indonesian-Australian basin (5°S–20°S and 100°E–125°E) (Domingues et al., 2007; Menezes et al., 2013, and references therein). Part of the northern SICC merges with the EGC around 15°S, 100°E (Fig. 10a). The EGC supplies ITF-origin water to the Leeuwin Current (LC) and is an essential component of the LC dynamics (Domingues et al., 2007; Benthuyssen et al., 2014; Lambert et al., 2016; Furue et al., 2013, 2017; Yit Sen Bull and van Seville, 2016). The geostrophic flow of the EGC is controlled by the meridional salinity gradient, making its dynamics distinct from the temperature dominated SICC (Menezes et al., 2013). This salinity front is formed by the encounter of the fresh Indonesian Throughflow Water carried westward by the SEC, and the salty subtropical underwater formed at the Southern Indian Ocean subtropical salinity maximum. The seasonal cycles of the EGC and the SICC are also distinct: the EGC is stronger in austral winter (3–5 Sv) and weaker (<0.5 Sv) in summer with the cycle in phase with the Leeuwin Current (Feng et al., 2003; Menezes et al., 2013; Furue et al., 2017). The SICC is overall stronger in spring-summer and weaker in winter (Palastanga et al., 2007; Jia et al., 2011a; Menezes et al., 2014) and experiences strong interannual variability, which peaks at biennial timescales and is decadal modulated (Menezes et al., 2016).

The multiple jets of the SICC are embedded in a zone of high eddy kinetic energy, with eddies generated by instabilities of the Leeuwin Current and of the SICC itself (Palastanga et al., 2007; Divakaran and Brassington, 2011; Huhn et al., 2012; Jia et al., 2011a, 2011b; Menezes et al., 2014, 2016; Siedler et al., 2006). By co-locating Argo floats and satellite data, Dilmahamod et al. (2018) described the passage of surface and subsurface South Indian Ocean eddies (SIDDIES). These westward-propagating, long-lived features (>3 months) originate in areas of high evaporation in the eastern Indian Ocean and prevail over a preferential latitude band, forming a permanent structure linking the eastern to the western Indian Ocean (the “SIDDIES Corridor”). This corridor of eddy passage allows the advection of water masses and biogeochemical properties across the basin (Dilmahamod et al., 2018).

4.2.4 Eastern Boundary

Unlike any other eastern boundary current, the Leeuwin Current (LC; Figs. 10 and 11a) flows poleward, along the shelf break of the west coast of Australia (Smith et al., 1991). Figure 11a presents the long-term average volume transport of the LC System from an observational climatology with similar structure found in a 1/10° ocean general circulation model

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The central branch of the SICC is found east of 75°–80°E, between 22° S and 24°S, and has a volume transport between 0.5 Sv and 3–5 Sv (Menezes et al., 2014). The shallowest and weakest northern SICC jet (1 Sv) has an equatorward orientation and flows on the southern flank of the northern cell of the subtropical gyre (Palastanga et al., 2007; Menezes et al., 2014). Part of the northern SICC merges with the

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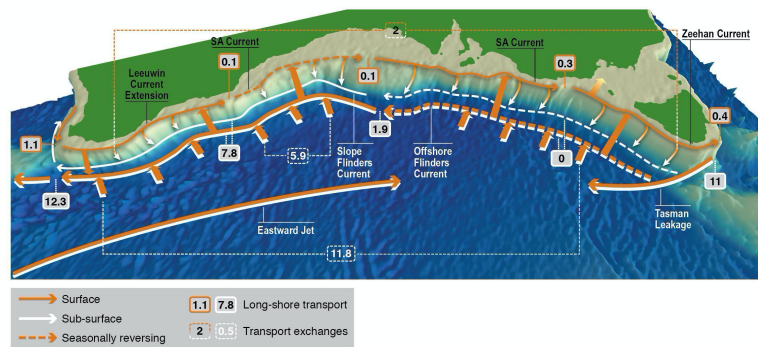
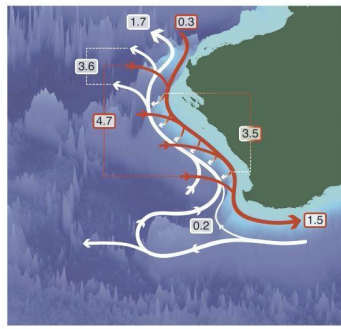
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(Furue et al., 2017). The primary source waters for the LC are the interior eastward flows (Section 4.2.3) that turn southeastward as they approach the coast and merge with the LC (Fig. 11a; Domingues et al., 2007; D'Adamo et al., 2009; Menezes et al., 2013, 2014; Furue et al., 2017, 2019). On average, the LC carries 0.3 Sv southward at 22°S, gains 4.7 Sv from the Indian Ocean interior, loses 3.5 Sv through downwelling to the layer beneath, and carries 1.5 Sv at its southern limit. The LC is approximately 200–300 m deep, extends from 22°S (North West Cape) to 34°S (Cape Leeuwin) and exists throughout the year despite significant seasonality (Feng et al., 2003; Ridgway and Godfrey, 2015; Furue et al., 2017). The Holloway Current, which flows southwestward on the North West Shelf (D'Adamo et al., 2009; Bahmanpour et al., 2016), is another weaker source to the LC from the north. Inshore of the LC, there exist seasonal equatorward flows that recirculate waters of distinct watermass properties influenced by air-sea interaction over the continental shelf (Woo et al., 2006).

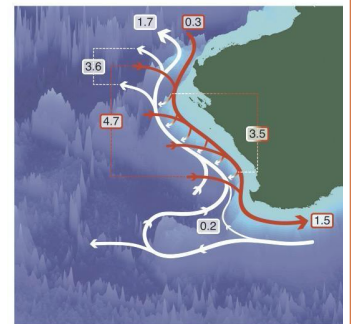


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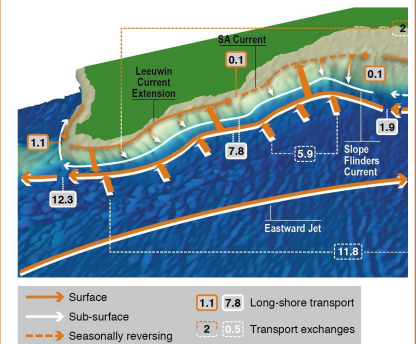
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Figure 11: a) Schematic summary of Australia's Leeuwin Current System three-dimensional transports (Sv). The red arrows and red-outline numbers represent the upper-layer (0–200 m) meridional transport of the poleward Leeuwin Current and meridionally-integrated zonal transport of the shallow eastward flows. The white arrows represent the lower-layer (200–900 m) flows of the Leeuwin Undercurrent. Taken from Furue et al. (2017). © American Meteorological Society. Used with permission. b) Schematic summary of the Southern Australia Current System three-dimensional transports (Sv). Long-shore transport for the Shelf Break Currents and Flinders Current in grey box with orange and white outlines, respectively. Integrated vertical and onshore flow transport in dashed outline box. Reprinted from Duran et al. (2020), with permission from Elsevier, Progress in Oceanography. Both schematics are based on a geostrophic calculation in the CARS 1/4-degree climatology.

The mean state of the LC is driven by the meridional pressure gradient in the upper ocean (e.g., Godfrey and Ridgway, 1985; Godfrey and Weaver, 1989, 1991), evident as a large poleward decrease in SSH balanced by an eastward surface geostrophic current (Section 4.2.3.1). The eastward flow approaches the eastern boundary, inducing downwelling and a surface poleward current (Fig. 11a; Godfrey and Ridgway, 1985; McCreary et al., 1986; Thompson, 1987; Weaver and Middleton, 1989, 1990; Furue et al., 2013; Benthuyssen et al., 2014), in opposition to the prevailing southerly winds. As a poleward boundary current, the LC waters are relatively fresh and warm from tropical origins (Rochford, 1969; Andrews, 1977; Legeckis and Cresswell 1981; Domingues et al., 2007; Woo and Pattiaratchi, 2008). Saltier Indian Central Water, the surface water of the Subtropical South Indian Ocean, joins the LC as it flows poleward (Section 4.2.3) increasing the mean density of the LC. Surface cooling along the poleward path also contributes to the increase in density (Woo and Pattiaratchi, 2008; Furue, 2019).

The LC flows around the southwestern corner of Australia and continues to flow eastward along the shelf break of the south coast of Australia to reach the southern tip of Tasmania near 42°S, 140°E (Fig. 11b; Oliver et al., 2016; Oke et al., 2018; Duran et al., 2020). This 5500-km long boundary current was first documented as a continuous flow by Ridgway and Condie (2004). When the longshore current is weak, however, it tends to be somewhat fragmentary (Oke et al., 2018; Duran et al., 2020) and sometimes even reverses in places (Duran et al., 2020). For this reason, and additionally because of the scarcity of observational sampling, the current is not traditionally regarded as a single current. Along southern Australia, the boundary currents can be described following Ridgway and Condie's (2004) naming convention. The current's western sector is called the Leeuwin Current Extension, the central part, to the south of the Great Australian Bight, is called the South Australian Current, and the easternmost part along Tasmania is called the Zeehan Current. They are collectively known as the Shelf-Break Currents (SBCs) of the Southern Australia Current System (Duran et al., 2020). It is not clear whether the SBCs along the south coast of Australia are, dynamically, an extension of the LC. The SBCs are at least consistent with the local northward Ekman drift (Ridgway and Condie, 2004; Duran et al., 2020) and hence would exist without the LC.

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Deleted: Beneath the LC, the Leeuwin Undercurrent exists in a depth range of 200–900 m, just offshore of the LC (Furue et al., 2017). It is fed in the south by a weak northward deflection of the Flinders Current, and a stronger northeastward recirculation of the Leeuwin Undercurrent around topography. The Leeuwin Undercurrent carries 0.2 Sv northward at its southern end, receives 3.5 Sv downwelled from the layer above, loses 3.6 Sv to westward flow into the Indian Ocean interior, and carries 1.7 Sv northward at its northern end near 22°S (Furue et al., 2017).

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Moved up [31]: The LC flows around the southwestern corner of Australia and continues to flow eastward along the shelf break of the south coast of Australia to reach the southern tip of Tasmania near 42°S, 140°E (Fig.

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Moved up [30]: As a poleward boundary current, the LC waters are relatively fresh and warm from tropical origins (Rochford, 1969; Andrews, 1977; Legeckis and Cresswell 1981; Domingues et al., 2007; Woo and Pattiaratchi, 2008).

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The LC's mean state

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On seasonal timescales, the LC transport generally tends to be strongest in austral autumn and weakest in austral summer (McCreary et al., 1986; Smith et al., 1991; Feng et al., 2003; Furue et al., 2017). There are two theories to explain this seasonality. In one, the local winds, which generally induce an offshore Ekman drift and therefore tend to weaken the LC, reach their annual maximum or minimum when the LC transport reaches its minimum or maximum, respectively (McCreary et al., 1986; Furue et al., 2013). In the other, a seasonal pressure anomaly originates in the Gulf of Carpentaria and propagates counterclockwise along the shelf break, driving the seasonality of the LC and of the SBCs to the south of Australia (Ridgway and Godfrey, 2015). Like the LC, the SBCs tend to be strongest in austral autumn and weakest in austral summer (Ridgway and Condie, 2004; Oke et al., 2018; Duran et al., 2020). In particular, the eastern part of the South Australian Current is seen to reverse in summer (Duran et al., 2020). This variability is consistent with the counterclockwise propagation of pressure anomaly shown by Ridgway and Godfrey (2015) and also with the seasonality of the wind stress along the south coast of Australia, with onshore (offshore) Ekman drift tending to drive eastward (westward) shelf-break flow (Duran et al. 2020).

On interannual time scales, the LC is modulated by the El Niño Southern Oscillation owing to the steric height anomalies in the western equatorial Pacific Ocean propagating through the Indonesian Seas and along Western Australia (Feng et al., 2003). During El Niño and La Niña periods, the LC transport weakens and strengthens, respectively, and is correlated with Fremantle sea level (Feng et al., 2003). During the strong 2010–2011 La Niña event, the LC reached record strength speeds (Feng et al., 2013) and the consequences of the unprecedented marine heat wave that resulted are described in Section 6.4.

On multidecadal timescales, the major boundary currents around Australia, including the LC, are reported to have strengthened during 1979–2014 in an eddy-resolving OGCM, consistent also with observations (Feng et al., 2016; see Section 6.1 for associated changes). At intraseasonal timescales, winds or heat anomalies on the North West Shelf region due to MJO events lead to intraseasonal variability of the Holloway Current on the North West Shelf and then of the LC (Marshall and Hendon, 2014; Marin and Feng, 2019).

The LC is accompanied by mesoscale eddies that cause the LC to meander energetically (Pearce and Griffiths, 1991; Feng et al., 2005; Waite et al., 2007; Meuleners et al., 2008). Those eddies are, at least partially, generated by barotropic, baroclinic, or mixed instability of the LC itself (Pearce and Griffiths, 1991; Feng et al., 2005; Meuleners et al., 2008). The eddy kinetic energy is greatest when the LC transport is strongest, in May–June (Fang and Morrow, 2003; Feng et al., 2005). Some of these eddies cause a large meander of the LC: a large anti-cyclonic eddy often forms at 28°–29°S and at 31°–32°S (Feng et al., 2003; Feng et al., 2007) steering the LC offshore to return to the continental shelf further south. This state typically starts during May–June and ends in July–August (Feng et al., 2007). Similarly, it is suggested that the eastern part of the SBCs becomes unstable in boreal autumn and winter, generating eddies, which subsequently propagate westward south of Australia (Oke et al., 2018). Turbulent mixing has been found to be enhanced in anticyclonic eddies

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1395 near the surface, and in cyclonic eddies at deeper levels (500-1000 m) due to the interaction of the eddies and near-inertial
1396 waves, which has implications for watermass modifications and the meridional overturning circulation (Cyriac et al. 2021).

1397
1398 Just below the Leeuwin Current is the equatorward Leeuwin Undercurrent (LUC; Thompson, 1984; Church et al., 1989;
1399 Smith et al., 1991; Fig. 11a). The LUC hugs the continental slope and extends from 200 m to 900 m (Furue et al., 2017).
1400 The LUC begins at Cape Leeuwin (34°S, 114°E) and is fed by a northward bend of a small fraction of the Flinders Current
1401 (FC; Fig. 10, 11; Furue et al., 2017). The remaining part of the FC continues westward but another small fraction of it
1402 appears to retroflect eastward and join and augment the LUC (Duran, 2015; Furue et al., 2017). Near 22°S, most of the
1403 LUC volume leaves the continental slope and flows offshore (Duran, 2015), apparently following the southern flank of
1404 the Exmouth Plateau although its bottom at 900 m is much shallower than the topographic feature (Fig. 11a).

1405
1406 To the south of Australia, an undercurrent has been recently identified below the Zeehan Current in a numerical simulation
1407 (Oke et al., 2018) and in a geostrophic calculation based on a gridded T-S climatology (Duran et al., 2020). Traditionally
1408 this flow was identified as a branch of the FC (Cirano and Middleton, 2004; Rosell-Fieschi et al., 2013; Feng et al., 2016)
1409 because the former flows in the same direction as the latter, but the FC as the northern boundary current of the subtropical
1410 gyre cannot exist on an eastern boundary (Anderson and Gill, 1975; Philander and Yoon, 1982; McCreary et al., 1992)
1411 and it jacks the vertical structure of an undercurrent (Duran et al., 2020). This northwestward- or westward-flowing
1412 undercurrent appears to exist all the way from the west coast of Tasmania to Cape Leeuwin (the southwestern tip of
1413 Australia) but its separation from the FC is less clear to the south of the Great Australian Bight and further west, where
1414 the FC accelerates and tends to overwhelm the undercurrent (Duran et al., 2020). Below, we call this current “slope FC”
1415 following Duran et al. (2020).

1416
1417 The mechanisms responsible for the LUC and undercurrent off southern Australia remain an open question, although
1418 models have been developed to investigate potential processes. The linear, continuously stratified models of McCreary et
1419 al. (1986) and Kundu and McCreary (1986) produce a surface poleward and a subsurface equatorward current, resembling
1420 the LC and LUC, along the eastern boundary. This class of model, however, requires large vertical diffusivity to produce
1421 a realistic LC and LUC (McCreary, 2013, personal communication). Along a continental slope, alongshore and cross-shelf
1422 buoyancy advection cause a shelf break front, forming a surface intensified poleward current, like the LC, and an
1423 equatorward undercurrent by thermal wind shear (Benthuyssen et al., 2014). Analytical shelf models have been extended
1424 to include cross-shelf buoyancy gradients to derive a poleward undercurrent like the LUC (Schloesser, 2014). These
1425 process-based analytical theories have not been tested in an eddy-resolving model.

1426
1427 The LUC and the slope FC are connected to the LC and the SBCs, respectively, by downwelling (Fig. 11; Furue et al.,
1428 2017; Duran et al., 2020), suggesting a common, but as yet unexplained, dynamics. Note, however, that for the LC–LUC

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pair, the mean downwelling appears to occur along isopycnal surfaces, and hence the LC water mass is not found in the LUC (Furue, 2019). For the SBCs and the slope FC, the nature of the downwelling is not known. The seasonality of these undercurrents are not well known. No systematic seasonal variability of the LUC was evident in a hydrographic climatology and ocean general circulation model (Furue et al., 2017).

4.2.5. Biogeochemical Variability

The ITF impacts both ocean currents and basin-scale biogeochemistry (Talley and Sprintall, 2005; George et al., 2013; van Sebille et al., 2014). Talley and Sprintall (2005) mapped silicate on the 31.96 potential density surface, revealing a striking silicate maximum associated with the SEC that extends westward to at least 60°E, highlighting the broad reach of ITF nutrient influence into the Indian Ocean. Ayers et al. (2014) estimated the depth- and time-resolved nitrate, phosphate, and silicate fluxes at the three main exit passages of the ITF that feed into the SEC: Lombok Strait, Ombai Strait, and Timor Passage. They found that the nutrient flux is significant relative to basin wide new production, and that the majority of ITF nutrient supply to the Indian Ocean via the SEC is to thermocline waters, where it is likely to support primary production and significantly impact biogeochemical cycling.

Satellite chlorophyll and primary production estimates suggest that values in the SEC are considerably higher than those found in the southern hemisphere subtropical gyre to the south, with Chla from ~0.10 to 1.0 mg/m³ and primary production from ~400 to 1000 mgC m⁻² d⁻¹ (Fig. 5; Figs. 5 and 6 in Hood et al., 2017). The highest concentrations and rates in the SEC are observed in the Eastern Indian Ocean in July and August during austral winter, associated with the ITF nutrient sources and upwelling off Java. The lowest chlorophyll concentrations and rates are observed in January (austral summer).

Model results and satellite observations show that the SEC/SCTR region exhibits an annual cycle in surface chla concentration and primary production, with the highest values in austral winter (June-August; > 0.20 mg/m³ and >600 mgC m⁻² d⁻¹, respectively) due to the strong southeasterly winds that increase wind stirring and induce upwelling (Resplandy et al., 2009; Dilmahamod, 2014; Fig. 5; Figs. 5 and 6 in Hood et al., 2017). Vertical sections of the SEC/SCTR region also reveal a deep chla maximum (George et al., 2013). Along 65°E this maximum shoals from > 100 m at 16°S to ~50 m at 10°S due to upwelling. The increases in surface Chl-a concentrations in austral winter are associated with decreases in the subsurface chla maximum (Resplandy et al., 2009; Dilmahamod, 2014). Surface freshening associated with the core of the SEC also influences the chla distribution in the SCTR region by modulating the static stability and mixed layer depth (George et al. (2013)).

The SEC provides relatively oligotrophic (low nutrient, low chlorophyll and low primary production) tropical source waters that feed into the EACC, NEMC, SEMC and the Mozambique channel. Chlorophyll *a* concentrations and

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production rates in Mozambique Channel surface waters are generally low ($< 0.4 \text{ mg/m}^3$ and $< 700 \text{ gC m}^{-2} \text{ d}^{-1}$, Fig. 5), and not significantly different in cyclonic and anticyclonic eddies (Lamont et al., 2014; Barlow et al., 2014; Figs. 5, 6 and 20 in Hood et al., 2017). Deep chlorophyll maxima are observed between 25 and 125 m depth depending on the proximity to the shelf and the influence of mesoscale eddies (Barlow et al., 2014; Lamont et al., 2014). Eddies in the Mozambique Channel also have a strong influence on the lateral transport of nutrients and chlorophyll from the coasts of Madagascar and Africa. Indeed, enhanced phytoplankton production within both cyclonic and anticyclonic eddies in the Mozambique Channel often occurs in response to lateral nutrient inputs into the euphotic zone by horizontal advection from the coasts of Madagascar and Africa rather than through eddy induced upwelling and downwelling (José et al., 2014; Lamont et al., 2014; Roberts et al., 2014). In contrast, in the Southeast Madagascar Current, topographically-induced coastal upwelling brings cold, nutrient-rich water up to the surface, which supports high rates of primary production (Lutjeharms and Machu, 2000; Ho et al., 2004; Quartly and Srokosz, 2004). This upwelling and its impacts are observed in both the austral summer and winter (Ho et al., 2004).

The Agulhas Current itself is warm and oligotrophic with sources derived from low nutrient and low chlorophyll surface waters from the Mozambique Channel, Southeast Madagascar Current and the southwestern tropical Indian Ocean (Fig. 5; Lutjeharms, 2006). Chlorophyll *a* concentrations and production rates in Agulhas Current surface waters are particularly low during austral summer ($< 0.2 \text{ mg/m}^3$ and $< 500 \text{ mgC m}^{-2} \text{ d}^{-1}$) with higher concentrations and rates in the austral winter (Machu and Garcon, 2001; Figs. 5, 6 and 20 in Hood et al., 2017). The Agulhas Current can drive upwelling and elevate primary production in the coastal zone through meandering and topographic interactions, but it can also dramatically suppress primary production when it impinges onto the shelf (Schumann et al., 2005).

In general, chlorophyll concentrations and primary production are elevated in the coastal zone of southeast Africa along the inshore side of the Agulhas Current (Fig. 5; Machu and Garcon, 2001; Goschen et al., 2012; Figs. 5, 6 and 20 in Hood et al., 2017). This enhancement is most pronounced in austral summer and further southward downstream, and it is associated with upwelling-favourable (easterly) winds and the aforementioned topographically-induced upwelling.

The near-surface eastward flows are generally associated with very low (oligotrophic) nutrient and chlorophyll-*a* (Chl-*a*) concentrations ($< 0.1 \text{ mg/m}^3$) and also very low primary production ($< 500 \text{ mgC m}^{-2} \text{ d}^{-1}$; Fig. 5 Figs. 5 and 6 in Hood et al., 2017). A well-defined deep Chl-*a* maximum is observed between 50 and 150 m during the austral fall along 55°E between 20 and 30°S (Coles et al., unpublished data). An exception to this, however, is the South-East Madagascar bloom (SMB). The SMB occurs in near-surface waters off the southeastern coast of Madagascar in the late austral summer/fall (Jan-April). It was first described as a dendroid bloom by Longhurst (2001), owing to its branching shape that projects eastward (Fig. 12). The bloom can extend over a 2,500 km² area with Chl-*a* concentrations reaching 2-3 mg/m³ (Longhurst, 2001), making it a 'hot spot' for primary production in an otherwise oligotrophic region. Fig. 12 illustrates the bloom's

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large spatial variability, with high Chl-*a* filaments apparently co-occurring and being transported with mesoscale and submesoscale eddies and jets.

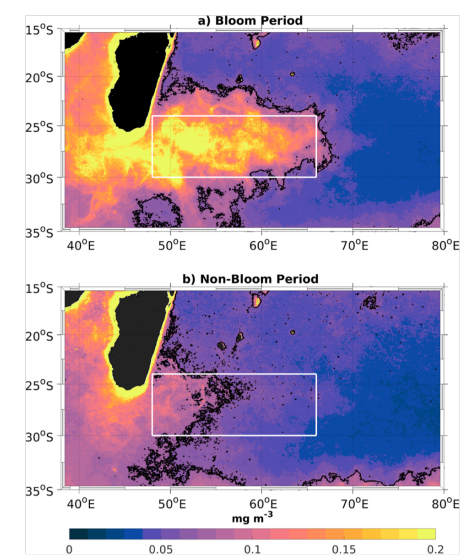


Figure 12: (a) Spatial maps of mean Chl-*a* concentration (mg/m³) during months of maximum austral summer bloom. (b) Same as (a) but during January of minimum Chl-*a* concentration in austral summer. The black contour denotes the 0.07 mg/m³ threshold used to distinguish between bloom and non-bloom years. From Dilmahamod et al. (2019).

Why the SMB flourishes in late austral summer is unclear. Longhurst (2001) attributed SMB development to mixed layer deepening and entrainment of nutrients by the vigorous mesoscale eddy field. These nutrients could stimulate phytoplankton growth in the photic zone, with the eddies shaping the eastward propagation of the enhanced surface Chl-*a* concentrations. However, Uz (2007), Srokosz and Quartly (2013), and Dilmahamod et al. (2019) subsequently showed that the bloom occurs within a warm (> 26.5°C), shallow mixed layer (~30 m) overlying a strong pycnocline. Furthermore, they suggested that diazotrophs known to inhabit the region (Poulton et al., 2009) might introduce new nitrogen (N) from

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N₂ fixation that could support the enhanced Chl-*a* concentration as observed elsewhere (Mulholland et al., 2014; Hood et al. 2004; Coles et al., 2004). Subsequent studies also highlight the role of mesoscale eddies (Fig. 12), that could advect, disperse and co-mingle nutrients and/or phytoplankton biomass (Dilmahamod et al. 2019; Huhn et al., 2012; Raj et al., 2010; Srokosz and Quartly, 2013; Srokosz et al., 2004, 2015; Uz, 2007).

A different explanation of the SMB and its eastward projection was proposed by Srokosz et al. (2004). In their proposed mechanism, the bloom initiates off Madagascar due to coastal processes that bring limiting nutrients to the photic zone and phytoplankton are transported horizontally by mesoscale eddies, resulting in an eastward propagation of the bloom. Dilmahamod et al. (2020) extend this further using a model to suggest that, from a nutrient flux analysis, horizontal advection of low-salinity nutrient-rich Madagascan coastal waters can indeed trigger a phytoplankton bloom. Alternatively, the apparent eastward propagation of the SMB has recently been attributed to advection by the SICC (Fig. 10; Dilmahamod et al., 2019; Huhn et al., 2012; Wilson and Qiu, 2008). Indeed, Huhn et al. (2012) further suggested that the bloom is shaped by a meridional barrier of jet-like Lagrangian coherent structures associated with the SICC.

At the eastern boundary, the tropical source waters and downwelling tendency of the Leeuwin Current combine to create a warm, oligotrophic current with low productivity. Chl-*a* concentrations are usually < 30 mgChl *a* m⁻² and rates of primary production rates generally do not exceed 500 mgC m⁻² d⁻¹ (Koslow et al., 2008; Lourey et al., 2006; Lourey et al., 2013). Productivity in the Leeuwin Current is lowest during austral summer, when the water column is stratified. During summer, subsurface chlorophyll maxima are found between 50 and 120 m depth (Hanson et al., 2007) as observed in open ocean subtropical oligotrophic waters (e.g., Venrick, 1991). However, rates of primary production in near shore upwelling regions (e.g., off of the North West Cape during summer) can sometimes attain very high levels (3000–8000 mgC m⁻² day⁻¹) as observed in other eastern boundary upwelling zones (Furnas, 2007).

In all seasons, meanders in the Leeuwin Current give rise to warm core, anticyclonic eddies that carry moderately high chlorophyll coastal water offshore. The elevated chlorophyll concentrations in these eddies is due to the presence of coastal diatom communities. These diatoms are transported offshore into cooler oligotrophic waters that are dominated by much smaller open ocean phytoplankton species (Waite et al., 2007a; Paterson et al., 2008; Waite et al., 2016). These eddies, which can extend to more than 2000 m depth, are unusual because they are downwelling (anticyclonic) circulations that should inhibit the input of new nutrients from depth. Nonetheless, these eddies, and the elevated chlorophyll concentrations that are associated with them, persist for months (Feng et al., 2007; Moore et al., 2007; du Fois et al., 2014). It has been hypothesized that the diatom communities in these eddies are supported by internal nutrient recycling and/or lateral supply (Waite et al., 2007a; Paterson et al., 2013; Thompson et al., 2007, 2011).

Generation of these warm (and cold) core eddies by the Leeuwin Current is prolific between 20° and 35° S (Gaubert et al., 2013). Most of these eddies move directly westward and some may be very long-lived (Feng et al., 2005; Feng et al.,

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2007; Moore et al., 2007; Gaube et al., 2013; du Fois et al., 2014). The persistence and potential biogeochemical/ecological impacts of these eddies in the open ocean have not been investigated fully.

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4.3 Equatorial regime

4.3.1 Wyrтки Jets

Owing to the seasonally reversing monsoon winds, the equatorial Indian Ocean (EIO) exhibits unique characteristics and is in contrast with the equatorial Atlantic and Pacific Oceans. Unlike the other basins, the annual winds along the EIO are very weak and mostly meridionally oriented except during the two intermonsoon seasons between boreal winter (April-May) and summer (Oct-Nov) when strong westerly wind bursts prevail along the EIO (see Schott and McCreary, 2001 and references therein). The semi-annual cycle in the zonal wind is well known observationally and was shown to be due to the meridional advection of easterly momentum by the cross-equatorial monsoon winds (Ogata and Xie, 2011). The westerly winds force strong eastward jets in the top 100 m along the equator that are known as spring and fall Wyrтки Jets, respectively (Wyrтки, 1973). These surface jets are usually confined within the top 100 m of the water column (Han et al., 1999; Iskander et al., 2011) and deepen (shoal) the thermocline and elevate (lower) the sea level in the east (west) (Rao et al., 1989; Schott and McCreary, 2001; Nagura and McPhaden, 2010a). These jets play a major role in zonal redistribution of mass, heat, salt and other water properties at the Equator and in off-equatorial basins (Reppin et al., 1999; Murtugudde and Busalacchi, 1999; Han et al., 1999; McPhaden et al., 2015; Chatterjee et al., 2017). Long term ADCP observations from the RAMA equatorial mooring suggest that the fall jet in the central EIO is usually stronger with a maximum transport of ~19.7 Sv compared to the spring jet which shows maximum transport of ~14.9 Sv with comparable standard deviations (McPhaden et al., 2015).

These eastward surface zonal currents tend to propagate westward during spring and eastward during fall (Nagura and McPhaden, 2016). The westward phase propagation speed during spring is estimated to be on average between 0.7-1.5 m s⁻¹ (Qiu et al., 2009; Nagura and McPhaden, 2010a) and driven primarily by the westward propagating surface zonal winds associated with atmospheric deep convection that moves from the Maritime Continent to the northern Bay of Bengal during spring (Nagura and McPhaden, 2010b; Nagura and McPhaden, 2016). Equatorial Rossby waves may also contribute to this westward propagation (Nagura and McPhaden, 2010a). In contrast, during fall, as the deep convection moves southeastward, the surface equatorial zonal winds, and thus surface currents, propagate eastward.

The spring and fall Wyrтки Jets also show considerable intraseasonal and interannual variability. While the intraseasonal variability of the Wyrтки Jets has been shown to be influenced by their own instability (Sengupta et al., 2001, 2007; Han et al., 2004) and local winds (Masumoto et al., 2005; Sengupta et al., 2007; Iskander et al., 2009; Prerna et al., 2019), the

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interannual variability of the Wyrski Jets is mainly caused by the anomalous wind forcing along the EIO associated with ENSO (Murtugudde et al., 2000; Gnanaseelan et al., 2012; Joseph et al., 2012) and IOD (Nagura and McPhaden, 2010b; Nyadjro and McPhaden (2014); Prerna et al., 2019): IOD weakens (strengthens) the equatorial zonal winds during its positive (negative) phase. While IOD modulates the zonal winds along the entire equator, the influence of ENSO is primarily limited to the eastern part of the EIO (Gnanaseelan et al., 2012). Moreover, it has been shown that these climate modes affect the boreal fall jet more significantly than the boreal spring jet. Recent modelling studies suggest that MJO convection can lead to a stronger spring Wyrski jet particularly in the eastern EIO. The interannual variability of MJO can, therefore, contribute to the observed interannual variability of this equatorial jet as well (Deshpande et al., 2017; Prerna et al., 2019).

4.3.2 5-30 Day Ocean Waves and Instabilities

Meridional velocity along the equator shows prominent high frequency variability at all depths, in the periodic band of 10-20 days with a peak at ~15 days (referred to as biweekly variability) and in the 20-30 days band with a peak at ~25 days (Masumoto et al., 2005; David et al., 2011; Chatterjee et al., 2013; Smyth et al., 2014). This variability is attributed to Yanai waves, first discovered in the atmosphere (also referred to as mixed Rossby-Gravity waves; Yanai and Maruyama 1966; Arzeno et al., 2020; Pujana and McPhaden, 2021). Unlike Kelvin and Rossby waves, Yanai wave phases can propagate westward or eastward depending upon their frequency, but their group velocity is always eastward (Miyama et al., 2006). These waves lead to convergent meridional heat flux into the equatorial regime (Shinoda, 2010; Smyth et al., 2014). While these waves were first observed in the ocean in the late 1990s, the establishment of the equatorial RAMA moorings (McPhaden et al., 2009) over the last two decades has provided more insight into these processes. Bi-weekly (10-20 day) is shown to be forced by the direct meridional wind stress (Sengupta et al., 2004) and to some extent by the meridional gradient of the zonal wind stress (Miyama et al., 2006). The 20-30-day band can be excited by off-equatorial barotropic/baroclinic instabilities in addition to direct wind forcing. A detailed review of the biweekly variability is provided in Schott et al. (2009) and hence, we focus on the 20-30 day variability in this review.

While the 20-30-day oscillation in meridional velocity is reported near the surface in the central EIO (David et al., 2011), in the eastern EIO these variabilities are seen only in subsurface layers (100-200 m depth) of the water column (Masumoto et al., 2005). This indicates a possible downward energy propagation of a vertical beam that carries energy to deeper depths. In the central EIO, these 20-30-day Yanai waves are excited by horizontal shear between the westward-flowing South Equatorial Current and the eastward-flowing Southwest Monsoon Current during IOD events (Fig. 10a; David et al., 2011). In the western EIO, these waves are primarily driven by cross equatorial meridional winds (Chatterjee et al., 2013). During early boreal summer (June/July), when the Somali current begins to cross the Equator along the western boundary of the basin, it bends offshore to conserve potential vorticity (Schott and McCreary, 2001) and forms a gyral circulation known as the Southern Gyre (Fig. 10a). Subsequently, these swift currents turn barotropically unstable and

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The existence of the high frequency variability in the meridional velocity is observed all along the EIO and across the entire depth range

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1636 generate eddy flow that is advected southward to the Equator near the western boundary i.e. at ~50-55°E. They generate
1637 a westward propagating cross-equatorial flow with a wavelength set by the eddy field which is similar to the wavelength
1638 of 20-30 day Yanai waves and thus excite these frequencies efficiently (Chatterjee et al., 2013).

1639 The ocean response to convectively coupled Kelvin waves (CCKW) in the atmosphere was investigated using ocean glider
1640 measurements from the CINDY/DYNAMO field experiment (Webber et al., 2014; Matthews et al., 2014). CCKW are
1641 atmospheric weather systems that propagate eastward along the Equator and are an important constituent of the MJO
1642 convection (Baronowski et al., 2016). CCKW enhance surface wind speed and latent heat flux during their passage
1643 suppressing the diurnal cycle of SST and leading to sustained decrease in bulk SST of around 0.1°C, one third of the SST
1644 anomaly due to a single, average MJO event, suggesting the oceanographic impact could have a strong feedback on the
1645 MJO cycle (Baronowski et al., 2016). Using RAMA moored measurements of upper ocean and surface atmosphere
1646 variability, Pujiana and McPhaden (2018) demonstrated that CCKW force oceanic Kelvin waves, affect surface heat fluxes
1647 and generate upper ocean turbulence.

1648 4.3.3 Equatorial Upwelling and Downwelling

1649 In the Pacific and Atlantic Oceans, permanent easterlies drive permanent equatorial upwelling due to Ekman
1650 divergence, but in the Indian Ocean, where the mean winds are weak and westerly, permanent upwelling does not
1651 exist (Schott and McCreary, 2001). Mean westerly winds along the Equator are downwelling favorable, driving
1652 surface convergence and thermocline divergence, which has been observed and described with Argo and RAMA data
1653 (Wang and McPhaden, 2017). Instead of upwelling along the equator, coastal upwelling along the coasts of Sumatra
1654 and Java is prominent. During June-October, south-easterly trade winds blow close to the Equator and drive the
1655 offshore Ekman transport away from the Sumatra-Java coast (Quadfasel and Cresswell, 1992; Sprintall et al., 1999;
1656 Susanto et al., 2001). The associated wind-driven upwelling intensifies as the monsoon progresses, reaching its peak
1657 by August and finally weakening by October as the monsoon winds wane. Recent studies suggest that when the
1658 easterly winds prevail during summer, upwelling favourable Kelvin waves also contribute to intensifying the
1659 equatorial upwelling (Iskander et al., 2009; Chen et al., 2016). During boreal winter-early spring (December-March),
1660 an intermittent/weaker subsurface thermocline shoaling is evident (Chen et al., 2016). Subsequently, the prevalence
1661 of westerly winds, which drive downwelling Kelvin waves, depress the thermocline in the east (Susanto et al., 2001;
1662 Prerna et al., 2019). Apart from this seasonal cycle, interannual climatic variability associated with ENSO and IOD
1663 events (Saji et al., 1999; Vinayachandran et al., 1999; Nyadjiro and McPhaden, 2014) also influences the upwelling
1664 intensity in this region (Section 6.2).
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Moved down [32]: 4.3.4 Equatorial Undercurrents

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Moved down [33]: a much weaker and seasonally varying transient feature driven by seasonally reversing monsoon winds (Reppin et al., 1999; Schott and McCreary, 2001). The equatorial RAMA moorings have recorded an eastward EUC with a core within the thermocline during boreal winter and spring (Chen et al., 2015, 2019) and occasionally in summer and fall at a depth of 90-170 m water column (Iskandar and McPhaden, 2011). During

Deleted: water column (Iskandar and McPhaden, 2011). During winter, the eastward EUC is forced by the upwelling Kelvin and Rossby waves that are in turn forced by the easterly winds along the equator; during summer

Moved down [34]: the westward EUC is primarily forced by the eastward pressure gradient generated by the downwelling reflected Rossby waves off the eastern boundary of the basin. On intraseasonal timescales of 30-70-days, the EUC variability is dominated by that of Kelvin and Rossby waves of lower order baroclinic modes (Iskander and McPhaden, 2011).

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4.3.4 Equatorial Undercurrents

In the Pacific and Atlantic, easterly winds produce an eastward mean undercurrent in the thermocline but in the Indian Ocean westerly winds do not produce a mean westward undercurrent. The reason is that nonlinear momentum advection drives mean eastward currents in the thermocline that flow up the zonal pressure gradient (Nagura and McPhaden, 2014). The Indian Ocean Equatorial Undercurrent (EUC) is, therefore, a much weaker and seasonally varying transient feature driven by seasonally reversing monsoon winds (Reppin et al., 1999; Schott and McCreary, 2001). The equatorial RAMA moorings have recorded an eastward EUC with a core within the thermocline during boreal winter and spring (Chen et al., 2015, 2019) and occasionally in summer and fall at a depth of 90-170 m (Iskandar and McPhaden, 2011). During winter, the eastward EUC is forced by the upwelling Kelvin and Rossby waves that are in turn forced by easterly winds along the equator in that season. During summer, the westward EUC is primarily forced by the eastward pressure gradient generated by the downwelling reflected Rossby waves off the eastern boundary of the basin. On intraseasonal timescales of 30-70-days, the EUC variability is dominated by that of Kelvin and Rossby waves of lower order baroclinic modes (Iskander and McPhaden, 2011). The undercurrents also undergo significant interannual variations related to the IOD. These variations are important in the mass and heat balance on IOD time scales, with significant impacts on upwelling and SST (Zhang et al., 2014; Nyadjiro and McPhaden, 2014)

4.3.5 Cross-Equatorial Circulation

The cross-equatorial circulation in the upper ocean is achieved by the Cross-Equatorial Cell (CEC), driven by southern hemisphere southeasterly winds and the seasonally-reversing monsoon winds in the northern hemisphere (Miyama et al. 2003; Schott et al. 2002, 2004, 2009). Thermocline waters subducted in the subtropical southeast Indian Ocean move equatorward and enter the northern hemisphere via the western boundary to upwell off Somalia and Oman. The return across the Equator, the surface branch of the CEC, is via the near-surface meridional flow in the interior Indian Ocean that is southward in the mean at nearly all longitudes (Miyama et al., 2003; Lee, 2004). This cell is unique to the Indian Ocean and is consistent with Sverdrup dynamics, being driven by the predominantly negative wind stress curl (Godfrey et al., 2001; Miyama et al., 2003; Wang and McPhaden, 2017). It carries most of the cross-equatorial transport of mass and heat (Schott and McCreary, 2001) and helps to moderate the seasonal climate of the region. The seasonal cross-equatorial mass flux is oppositely directed along the western boundary and in the interior (Beal et al., 2013). Flow in the interior is directed from the summer to the winter hemisphere (Horii et al., 2013; Wang and McPhaden, 2017) consistent with monsoon wind forced Ekman and Sverdrup dynamics as proposed in the model study of Miyama et al. (2003).

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In OGCMs, the southward flow of the CEC was found to occur just below the surface, beneath a northward surface current (Wacogne and Pacanowski, 1996; Miyama et al. 2003). This “equatorial roll”, also unique to the Indian Ocean, is only of order 100 m depth and so has little impact on the cross-equatorial heat transport of the CEC. Horii et al. (2013) and Wang and McPhaden (2017) presented the first observational evidence for the equatorial roll.

The spatial structure and time evolution of the cross-equatorial circulation is difficult to depict due to its dependence on the fluctuating monsoon winds. Consequently, the flow patterns obtained from an Eulerian average as in Fig. 10 cannot capture the monsoon-dependent streamlines that a flow will follow at a given moment. Lagrangian methods based on ocean drifter velocities (Laurindo et al. 2017) and real and simulated surface drifter trajectories identify pathlines that connect the monsoonal Indian Ocean, revealing three cross-equatorial gyre pathways that connect the Somali Current with the interior flow north and south of the Equator (Fig. 7 in l’Hegaret et al., 2018).

4.3.6 Biogeochemical Variability

Much of the current understanding of biogeochemical variability in the equatorial zone of the Indian Ocean is based on satellite ocean color observations and models, augmented by some additional, relatively sparse, in situ measurements. Seasonal climatologies of near-surface chlorophyll concentrations and primary production show a significant seasonality in equatorial waters that is clearly associated with monsoon forcing (Fig. 5. Wiggert et al., 2006; Strutton et al., 2015; Figs. 5 and 6 in Hood et al., 2017). In general, Chl-*a* concentrations and primary production increase northward from the equator with the lowest concentrations ($< 0.1 \text{ mg m}^{-3}$) and rates ($< 800 \text{ mg C m}^{-2} \text{ d}^{-1}$) occurring during the boreal spring intermonsoon period. During the southwest monsoon, Chl-*a* concentrations and rates of primary production increase in western equatorial waters in response to monsoon-forced mixing and upwelling. However, concentrations and rates in the central and eastern equatorial waters stay relatively low ($< 0.5 \text{ mg m}^{-3}$, $< 800 \text{ mg C m}^{-2} \text{ d}^{-1}$, respectively). Island wake effects can be seen advecting high chlorophyll water ($> 0.5 \text{ mg m}^{-3}$) along the equator from the Chagos-Laccadive ridge at 73°E eastward during the autumn intermonsoon period and westward during spring (see Fig. 1 in Strutton et al., 2015).

Well-developed deep Chl-*a* maxima have been observed in the equatorial Indian Ocean along 65°E centered at about 50 m depth in November-December (George et al., 2013) and along 80°E centered at about 75 m in August-September (Sorokin et al., 1985). It is unknown whether or not this subsurface Chl-*a* maximum exists along the equator throughout the year, but it is probably present whenever the water column is stratified. Models predict the presence of a subsurface (60 m) Chl-*a* maximum in eastern Indian Ocean equatorial waters along 87°E (Wiggert et al., 2006) that is present throughout the year, except during the southwest monsoon when high chlorophyll surface water is advected into the region. Physical processes at time scales from intraseasonal to interannual (i.e., Wyrki Jets, MJO and IOD) have been shown to influence biogeochemistry. For example, IOD events can significantly increase chlorophyll concentrations and primary

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1775 production in eastern Indian Ocean equatorial waters (Wiggert *et al.*, 2009). In addition, relaxation of an IOD can deplete
1776 upper ocean nutrients, decreasing biological productivity (Kumar *et al.*, 2012). Biogeochemical responses to the IOD also
1777 have significant higher trophic level impacts (Marsac and Le Blanc, 1999).

1778 Satellite observations and biophysical model simulations show how chlorophyll concentrations and primary production
1779 near the Seychelles-Chagos thermocline ridge, can be increased by MJO-induced wind mixing and nutrient entrainment
1780 (Resplandy *et al.*, 2009). They also concluded that IOD-driven interannual variability of thermocline depth influences the
1781 biogeochemical response to MJO; the deepened nutricline following IOD events inhibits nutrient input into the mixed
1782 layer, and thus decreases the biogeochemical response to MJO.

1783 In model simulations, Wyrтки jets depress the thermocline and nitracline along the equator on the eastern side of the basin
1784 and, as a result, lower equatorial primary production when they arrive in the spring and autumn (Wiggert *et al.* 2006). This
1785 pattern was observed in a 25 day time series study on the equator at 80.5°E in late 2006 that showed a deepening of the
1786 surface layer, nutracline and subsurface Chl-*a* maximum during the autumn Wyrтки jet period (Kumar *et al.*, 2012).

1787 Finally, Strutton *et al.* (2015) examined time-series measurements of near-surface chlorophyll concentration from a
1788 mooring deployed in 2010 at 80.5°E in the equatorial Indian Ocean. These data revealed at least six spikes in chlorophyll
1789 from October through December, separated by approximately 2-week intervals, and coinciding with the development of
1790 the fall Wyrтки jets. The chlorophyll pulses were associated with increases in eastward surface winds and eastward currents
1791 in the mixed layer, and inconsistent with upwelling dynamics because eastward winds that cause intensification of the
1792 Wyrтки jet should drive downwelling. Strutton *et al.* (2015) concluded that the chlorophyll spikes could be explained by
1793 two alternative mechanisms: (1) turbulent entrainment of nutrients and/or chlorophyll from across the base of the mixed
1794 layer by wind stirring or Wyrтки jet-induced shear instability or (2) enhanced southward advection of high chlorophyll
1795 concentrations into the equatorial zone associated with wind-forced biweekly Yanai waves.

1796 4.4 Northern Indian Ocean

1797 The two main basins of the northern Indian Ocean, the Bay of Bengal (BoB) and the Arabian Sea (AS), are characterized
1798 at the surface by remarkably contrasting sea surface salinity with differences of the order of 3 psu (e.g. Chatterjee *et al.*
1799 2012, Gordon *et al.* 2016, Hormann *et al.* 2019) decreasing from west to east (Fig. 4). The fresh surface layer of the BoB
1800 is maintained by large freshwater input deriving from direct rainfall over the ocean and river runoff, especially during the
1801 South Asian monsoon. The salt balance of the BoB is maintained by the subsurface supply of salt water via the Southwest
1802 Monsoon Current (Fig. 10, Vinayachandran *et al.*, 2013). The saltier SSS of the AS is the consequence of an evaporative
1803 regime (e.g., Rao & Sivakumar, 2003; Sengupta *et al.*, 2006). A reversing monsoonal near-surface circulation (Fig. 10 a,b)

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Resplandy *et al.* (2009) used a combination of satellite
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1827 plays a central role in the exchanges of freshwater and heat between the BoB and the AS (McCreary et al. 1993, Hormann
1828 et al. 2019).

1829 Recent multi-year deployments of satellite tracked surface drifters drogued at 15 m depth (Wijesekera et. al, 2016,
1830 Centurioni et al. 2018) have helped to better constrain the amplitude and structure of the circulation and the exchange
1831 processes between the two basins, and to refine the findings reported by other authors (e.g. Schott and McCreary, 2001).
1832 Additionally, implementation of a moored buoy network along the slope and shelf of the Indian coast has helped
1833 significantly in enhancing our understanding of the east India Coastal Current (EICC) and west India Coastal Current
1834 (WICC) (Fig. 10, Mukherjee et al., 2014; Amol et al., 2014; Mukhopadhyay et al., 2020; Anya et al. 2020).

1835 **4.4.1 Bay of Bengal**

1836 In a climatological sense, the main features of the near-surface circulation of the western BoB (Figs. 10a and 10b) are the
1837 reversing EICC, the Southwest/Northeast Monsoon Current (SMC/NMC) and the seasonally variable Sri Lanka Dome.
1838 The eastern side of the BoB, extending into the Andaman Sea, is characterised by a sluggish circulation.

1839 **4.4.1.1 Southwest/Northeast Monsoon Currents**

1840 During the boreal summer SW monsoon, the Southwest Monsoon Current (SMC, Fig. 10a) flows eastward around the
1841 Indian subcontinent supplying salty water from the Arabian Sea to the fresher Bay of Bengal (e.g., Jensen, 2001; Jensen
1842 et al., 2016; Vinayachandran et al., 2013; Wijesekera et al., 2015, 2016). During the winter monsoon, the Northeast
1843 Monsoon Current (NMC, Fig. 10b) reverses the flow carrying fresher water into the Arabian Sea. Figure 13 provides a
1844 snapshot from an operational forecast system, the Coupled Ocean-Atmosphere Mesoscale Prediction System (COAMPS),
1845 of the NMC flow and route for freshwater to enter the Arabian Sea (Wijesekera et al., 2015).

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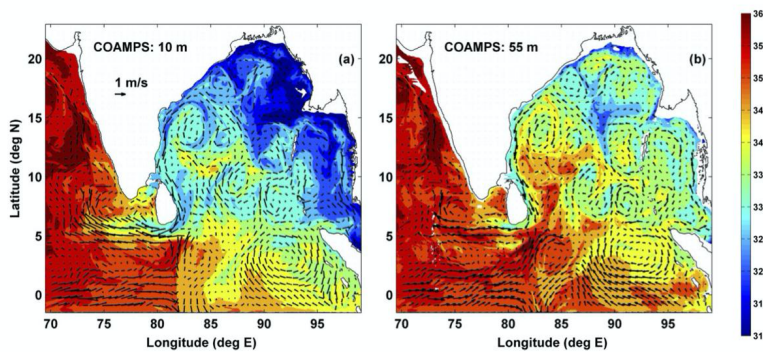
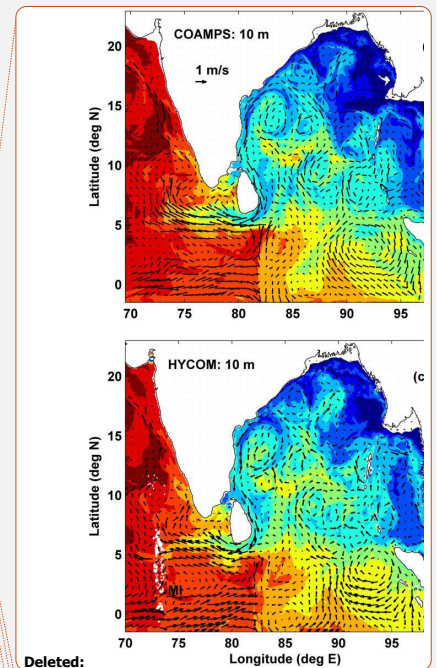


Figure 13: COAMPS velocity vectors (arrows) and salinity (psu, color shading) at (a) 10 m and (b) 55 m on 18 December 2013. Modified from Wijesekera et al. (2015).

A more recent study has found that the origins of the Arabian Sea high salinity water are specifically from the western Arabian Sea and western Equatorial Indian Ocean, and they reach the Bay of Bengal via a combination of the Indian Ocean EUC and the SMC (Sanchez-Franks et al., 2019; Section 8.2). Changes in the supply of salty water to the Bay of Bengal varies interannually due to the strength in the equatorial currents, forced by the local wind field and ENSO (Sanchez-Franks et al., 2019), and is expected to influence the salinity budget of the Bay of Bengal (Vinayachandran et al., 2013) and thus modulate SST variability (Fig. 10a, Jensen, 2001; Jensen et al., 2016; Li et al., 2017; Vinayachandran et al., 2013, 2018; Webber et al., 2018).

4.4.1.2 East Indian Coastal Currents (EICC)

The EICC forms the western boundary current of the Bay of Bengal and plays an important role in the basin-scale heat and salt budget of the Indian Ocean, and hence in determining the local climate (Shenoi et al, 2002), biological processes (Madhupratap et al, 2003; Vinayachandran et al, 2005; Naqvi et al, 2006; Dileepkumar, 2006; McCreary et al, 2009) and marine fisheries (Vivekananda and Krishnakumar, 2010) of this region. It reverses its direction seasonally north of 10°N in response to a combination of local alongshore winds, remote alongshore winds in the eastern BoB, remote forcing from the equatorial Indian Ocean and the interior Ekman pumping of the basin (Shankar et al., 1996; McCreary et al., 1996; Vinayachandran et al., 1996; Mukherjee et al., 2017). The EICC is generally equatorward south of 10°N throughout the



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1875 year. While local winds dominate the EICC forcing during summer and winter, remote forcing dominates during the inter-
1876 monsoon periods (Shankar et al., 1996; McCreary et al., 1996; Suresh et al., 2013).

1877 Climatological ship-drift and hydrographic data suggest the EICC flows poleward during February-September (Shetye et
1878 al., 1993) and turns equatorward during November-January (Shetye et al., 1996; Fig. 10). While the annual cycle is driven
1879 by local alongshore winds and interior Ekman pumping, the semiannual cycle is the result of asymmetry in the monsoon
1880 and equatorial forcing (Mukherjee et al., 2018). During boreal spring (March-May), the EICC is strongest with a magnitude
1881 exceeding 1 m/s with unidirectional currents to about 150 m, forming the western boundary current of a cyclonic basin-
1882 wide gyre of the BoB. The local alongshore winds are weakest and the stronger EICC is primarily forced by the interior
1883 anticyclonic Ekman pumping over the basin (McCreary et al., 1996; Shankar et al., 1996; Vinayachandran et al., 1996;
1884 Mukherjee et al., 2017). During boreal summer, the EICC is weaker and is restricted to within the top 70 m of the water
1885 column. The poleward flow is generally limited to the central part of the coast between 10-18°N and often switches to
1886 short pulses of poleward currents along the coast (Mukherjee et al., 2018; Francis et al., 2020). The poleward flow is
1887 driven by local winds, but the response of the interior cyclonic Ekman pumping and equatorial winds driving an opposite
1888 flow along the coast causes a weaker poleward EICC in summer than in spring (McCreary et al., 1996; Vinayachandran
1889 et al., 1996; Shankar et al., 2002). The basin-scale gyre also disappears in summer and the EICC then consists of several
1890 eddies along the coast. The EICC turns equatorward during November-January (Shetye et al., 1996).

1891 Near-surface alongshore currents also display significant 120 day and intraseasonal variability. The magnitude of the 120
1892 day variability is generally weaker than the semiannual period, particularly in the southern part of the coast. As for the
1893 annual period, upward phase propagation along the coast is also evident for the semiannual and 120 day period, except at
1894 Cuddalore where downward phase propagation is common during summer and winter months (Mukherjee et al., 2014;
1895 Mukhopadhyay et al., 2020). Further, unlike annual and semiannual periods, the 120 day and intraseasonal variability
1896 decorrelate along the coast indicating that these high frequencies are dominated by local responses rather than remote
1897 forcing (Mukherjee et al., 2018; Mukhopadhyay et al., 2020a).

1898 4.4.1.3 Undercurrents

1899 ADCP observations suggest that during summer and winter, when the near-surface current is shallow, the EICC often
1900 exhibits undercurrents along the continental slope. As the EICC is deeper in the north, the undercurrent is observed at a
1901 depth of 100-150 m and can extend up to 700 m. However, in the south undercurrents are seen at relatively shallow depths
1902 of about 70-75 m (Francis et al., 2020). While these undercurrents are observed throughout the coast, they are much more
1903 prominent and more frequent at Cuddalore, the southernmost station of the coast (Fig. 14, Mukherjee et al., 2014;
1904 Mukhopadhyay et al., 2020).

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Deleted: Along the east coast of India, ADCP observations reveal that variability in alongshore currents is dominated by seasonal timescales Mukhopadhyay et al. (2010a)

Deleted: The seasonal reversal of the EICC is apparent all along the east coast, suggesting that the canonical directions of the EICC are a robust feature and explaining their existence in ship-drift climatologies (Cutler and Swallow, 1984; Mariano et al., 1995). In the annual band, upward phase propagation is evident all along the coast, suggesting the dominance of remote forcing over local alongshore winds in driving the EICC. ¶

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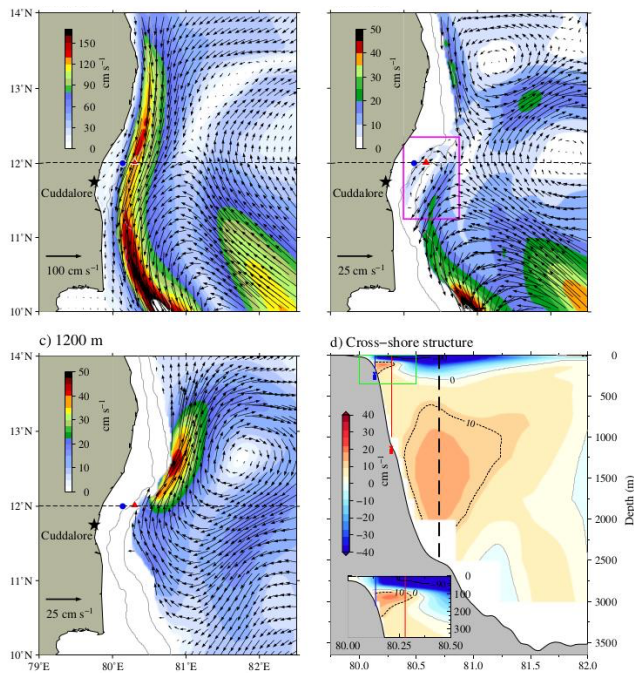
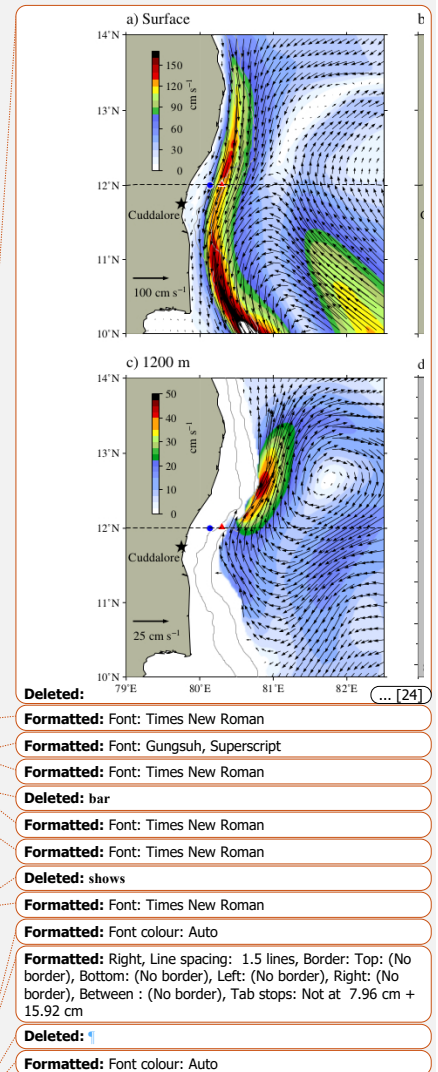


Figure 14: Circulation pattern in the southwestern Bay of Bengal at (a) surface (b) 200 m and (c) 1200 m on 15 November 2014. Vectors show the current direction, and overlaid is the current magnitude (cm s^{-1}). Note that the scales of current vectors and color bars are different at each subplot. Blue circle (red triangle) represents the location of ADCP deployed on the shelf (slope) off Cuddalore. Dashed black line represents the 12°N latitude. Continuous gray lines represent the 100 m and 1000 m bathymetric contours. Rectangular box (magenta) indicates the subsurface eddy near the shelf break. (d) Cross-shore structure of alongshore currents across 12°N . Dashed black vertical line shows the core of the undercurrent, and red (blue) vertical lines show the location of ADCP mooring on the slope (shelf). Inset plot is the zoomed view of shelf break region indicated by green box (Reproduced from Francis et al., 2020).

The prominent upward phase propagation of the annual signal in the subsurface layers, particularly in the southern stations, suggests downward propagation of energy and is thereby attributed as one of the main causes of the undercurrents



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1950 (Mukherjee et al., 2014). A recent modelling study suggests that the wintertime undercurrent off Cuddalore consists of
1951 two separate subsurface anticyclonic eddy circulations: a shallow small scale circulation at a depth range of 100-200 m
1952 and a broader and deep flow below 500 m depth off the continental slope (Francis et al., 2020). The shallow subsurface
1953 anticyclonic eddy was found to spin off from the zonal shear of the mean near-surface EICC along the shelf break (Fig.
1954 14). These eddies exhibit high frequency fluctuations and have 20-30 km length scales. Since the zonal share of the EICC
1955 is primarily linked to the strength of the EICC itself, the variability and strength of this undercurrent is also linked with
1956 the EICC.

1957 **4.4.1.4 Sri Lanka Dome**

1958 The Sri Lanka Dome (Vinayachandran et al. 1999; Schott and McCreary, 2001; Wijesekera et. al, 2016, Cullen and
1959 Shroyer, 2019) is mainly visible as a closed anticyclonic (clockwise) eddy in the near-surface geostrophic current velocity
1960 field starting in May and lasting through October (Fig. 15c). It is a recurring upwelling dome that forms east of Sri Lanka
1961 between 5-10°N, 83–87°E. The SLD is embedded within the Southwest Monsoon Current (SMC) system (Gadgil, 2003)
1962 and enhances the SMC exchange from the Arabian Sea to the Bay of Bengal (Anutaliya et al., 2017). Upwelling associated
1963 with the SLD influences the vertical exchange of water properties, enhances biological productivity, and cools sea surface
1964 temperature (SST) which affects local atmospheric convection (Vinayachandran et al., 2004; de Vos et al., 2014).

Deleted: Role of EICC in salt transport

The EICC starts flowing southward north of 15°N as early as June (Fig. 10a) and eventually develops into a strong coastal current during the winter monsoon (November through January) that extends all the way south of Sri Lanka and connects the northernmost part of the bay, where large amounts of freshwater are discharged by the Indian river system, with the Arabian Sea (AS; Fig. 4, Fig. 13). For this reason, the EICC is regarded as one of the possible pathways responsible for the exchange of salt between the AS and the BoB (Hormann et al. 2019), along with the SMC (Fig. 10a; Jensen, 2001; Jensen et al., 2016; Vinayachandran et al., 2013, 2018; Webber et al., 2018), which transports saline water north into the BoB during June-September. Some of this saline water is transported directly around the tip of India (Fig. 10a), yet there is also a role for transport along the equator from the western tropical Indian Ocean and the western AS (Sanchez-Franks et al., 2019). In February, the off-shore current south of Sri Lanka is still flowing westward into the AS but the EICC north of 10°N starts flowing northward, reaches its maximum strength in April and stops south of 10°N by June, joining the offshore limb of a cyclonic circulation further north that fills the northern Bay from June through August.

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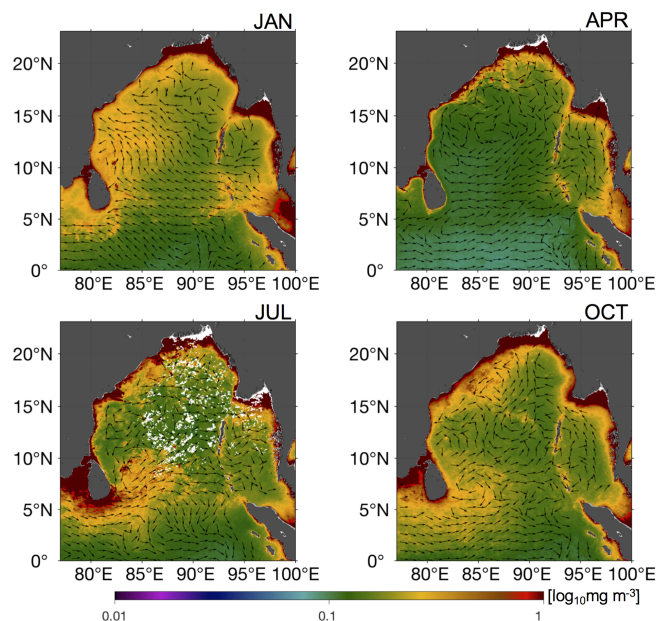


Figure 15: Climatology (2002–2018) of chlorophyll-a concentrations (colormap) and current velocities (arrows) in the Bay of Bengal for (a) January (b) April (c) July (d) October. Chlorophyll climatology was obtained from the MODIS-Aqua product and current velocities were obtained from the third-degree Ocean Surface Current Analysis Real-time (OSCAR) product.

4.4.2 Arabian Sea

Like the Bay of Bengal (BoB), the Arabian Sea (AS) near surface circulation is also driven primarily by the seasonally reversing monsoon winds. The AS is connected to the BoB through the passage between the southern tip of India and the equatorial wave guide, and to the southern hemisphere by the cross equatorial flow via the Somali current system. The Somali current (Fig. 10) forms one of the western boundary currents of the AS. Another major boundary current system

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Moved down [44]: during all seasons whenever and wherever wind forcing and/or currents are insufficiently strong to upwell or

Deleted: These nutrients stimulate diatom blooms (Sasamal et al., 2005) that give rise

Moved down [40]: to significant increases in Chl-*a* concentration (~ 30–100 mgChl*a* m⁻²) and production (~ 0.55–1 gC

Moved down [41]: When the EICC meanders seaward from the Indian coast, it leads to offshore increases in high chlorophyll water.

Moved down [42]: caused by wind-driven entrainment, not only of subsurface nutrients, but also of phytoplankton from the

Moved down [43]: 14). Presumably, this is due to a combination of the downwelling-favorable currents and winds.

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2218 is along the west coast of India, the WICC (Fig. 10), which transports heat and salt from the northern Arabian Sea to the
2219 BoB, and vice-versa. Recent observations (Chatterjee et al., 2012) and modelling studies (Shankar et al., 2015; Vijith et
2220 al., 2016) indicate that the northern extent of the WICC reaches up to 20°N during the winter monsoon, carrying fresher
2221 BoB water to the northern latitudes and modulating the wintertime convection there.

2222 4.4.2.1 Somali current System

2223 The Somali Current is a seasonally reversing western boundary current and is often composed of discontinuous non-linear
2224 eddy driven flows. During summer it flows poleward and the upwelling here is nearly as large as for the eastern boundary
2225 upwelling regimes of the Pacific and Atlantic Oceans (See Schott and McCreary, 2001 for a detailed review).
2226 Unfortunately, owing to piracy, direct in-situ observations are very rare in this region and mostly date back to the early
2227 1960s and 1970s. Hence, the scientific community has mostly relied on numerical model simulations to enhance
2228 understanding of this region over the last few decades.

2229 Recent modelling studies suggest that during the summer monsoon, unlike other western boundary currents, the Somali
2230 Current system can be divided into three dynamically distinctive regions (Wang et al., 2018; Chatterjee et al., 2019):
2231 northern (north of 8°N), central (3-8°N) and the southern (south of 3°N) part. The northern and southern parts are driven
2232 by the large anticyclonic gyres called the Great Whirl (GW) and the southern Gyre (SG), respectively (Fig. 10a). Local
2233 southwesterly alongshore winds known as the Findlater Jet (Findlater, 1969) drives Ekman transport all along the Somali
2234 coast (Schott and McCreary, 2001) with varied magnitude which is strongest in the southern part, and significantly
2235 weakens northward (Chatterjee et al., 2019). The wind stress forcing leads to Ekman Pumping in the central Arabian Sea,
2236 setting up a bowl-shaped mixed layer and warming at the 100 m level. Ekman downwelling velocities are strongest in the
2237 northern part and likely contribute to the formation of the Great Whirl front which upwells cold subsurface water in this
2238 part of the coast. The central part, in contrast, is mainly driven by the local winds and remotely forced Rossby waves. In
2239 fact, the annual Rossby waves radiated out of the southwestern coast of Sri Lanka seem to play a major role in the reversal
2240 of currents to poleward flow in the northern part of the Somali coast as early as mid April. This reversal likely initiates
2241 the generation of the Great Whirl (Beal and Donohue, 2013; Vic et al., 2014), a month before the strong northeastward
2242 Findlater Jet commences along the Somali coast. As the monsoon progresses, these downwelling favourable Rossby waves
2243 oppose the coastal Ekman upwelling and thereby start to weaken the upwelling all along the coast. Moreover, as the
2244 alongshore winds peak, this favours enhanced mixing at the bottom of the mixed layer, which deepens the thermocline
2245 further. This process is more conspicuous in the central part of the coast, where the depth of the 22°C isotherm deepens
2246 by about 30-40 m from June to August (Chatterjee et al., 2019). By this time, the upwelling becomes limited to the northern
2247 part of the coast along the Great Whirl front of the Somali region.

Deleted: In this section we will first discuss some recent developments of the Somali current system and then the WICC.

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2259 4.4.2.2 West India Coastal Current (WICC)

2260 The WICC reverses its direction annually: flowing equatorward (upwelling favourable) during the summer monsoon (May
2261 to September; Fig. 10a) and poleward (downwelling favourable) during the winter monsoon (November to February; Fig.
2262 10b). The equatorward flow during the summer characterises the WICC as a classical eastern boundary current (Shetye
2263 and Shenoi, 1988). Interestingly, as the monthly mean alongshore winds off the west coast of India are always equatorward,
2264 throughout the year, the surface currents flow against the winds during the winter, driven by coastally-trapped Kelvin
2265 waves forced remotely in the BoB and along the east coast of India (McCreary et al., 1993; Shankar and Shetye, 1997;
2266 Shankar et al., 2002; Suresh et al., 2016). Recent observations based on satellite data and alongshore ADCP moorings
2267 reveal strong interannual variability of this seasonal cycle. Vialard et al. (2009), based on a short ADCP record during
2268 2006–2008, reported an absence of seasonal cycle off Goa and they attributed this absence to the radiation of Rossby waves
2269 south of the critical latitude. As the longer record of ADCP data became available, a clear seasonal cycle in the WICC
2270 became evident with weaker amplitudes in the south, stronger poleward (Amol et al., 2014).

2271 The WICC also shows significant intraseasonal variability at times, particularly during boreal winter, exhibiting much
2272 stronger energy in the intraseasonal band than in the seasonal band (Vialard et al., 2009; Amol et al., 2014). Unlike the
2273 seasonal cycle, intraseasonal variability is stronger in the south and weakens poleward. Vialard et al. (2009) attributed this
2274 intraseasonal variability to the atmospheric MJO forcing. Recently, a modelling study suggested that interception of the
2275 intraseasonal equatorial Rossby waves by the southern tip of India and Sri Lanka excites coastal Kelvin waves which
2276 contribute significantly (~60–70%) to the intraseasonal variability along the west coast (Suresh et al., 2013). A satellite
2277 sea level study by Dhage and Sturb (2016) confirmed the model-based findings of Suresh et al. (2013) and revealed that
2278 large-scale winds from the south of India and Sri Lanka also contribute to the coastal signals along the west coast of India.

2279 Another striking feature observed in these ADCP data is the clear signature of upward phase propagation in all timescales
2280 during both monsoon seasons. This upward phase propagation is more conspicuous for the seasonal period than for the
2281 intraseasonal. As a result, the phase of the surface currents often tends to be opposite that in the subsurface layers (Amol
2282 et al., 2014). Moreover, it is found that the strength of this undercurrent intensifies northward along the west coast with
2283 strongest undercurrent off Mumbai and the weakest off Kanyakumari (southernmost point of Indian mainland), indicating
2284 a possible downward propagation of energy along the ray path as suggested earlier by Nethery and Shankar (2007). Since
2285 the ray angle (θ) depends on the frequency (σ) and stratification (N_b) according to $\theta = \sigma / N_b$ (McCreary, 1984; Nethery and
2286 Shankar, 2007) the angle the beam makes from the horizontal is deeper for the intraseasonal band than for the seasonal.
2287 As a result, intraseasonal beams propagate energy deeper into the water column. Therefore, while the WICC shows some
2288 coherence along the coast in the seasonal time scale, it completely decorrelates horizontally for the intraseasonal period.

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4.4.3 Biogeochemical Variability

In the Bay of Bengal, the large freshwater input gives rise to enhanced stratification that inhibits upwelling and wind-mixing and therefore nutrient supply to surface waters (Kumar et al., 2002; Vinayachandran et al., 2002; Madhupratap et al., 2003; Vinayachandran, 2009). Nonetheless, increased productivity is observed along the coast primarily in association with riverine nutrient inputs (Vinayachandran, 2009). These nutrients stimulate diatom blooms (Sasamal et al., 2005) leading to significant increases in Chl-*a* concentration ($\sim 30\text{--}100\text{ mgChla m}^{-2}$) and production ($\sim 0.55\text{--}1\text{ gC m}^{-2}\text{ d}^{-1}$) near the coast (Gomes et al., 2000; Fig. 15). This high Chl-*a* river water flows either along the coast or offshore, up to several hundred kilometers, depending on the coastal current pattern (Vinayachandran, 2009). Along the Indian coast, the flow of Chl-*a*-rich water is determined by the EICC, which flows northward during the spring intermonsoon period and Southwest Monsoon and southward during the autumn intermonsoon and Northeast Monsoon (Fig. 15). When the EICC meanders seaward from the Indian coast, it leads to offshore increases in high chlorophyll water. During the spring intermonsoon and Southwest Monsoon the northward-flowing EICC is upwelling favorable, which may contribute to increases in Chl-*a* concentration and primary production along the coast (Hood et al., 2017).

Elevated productivity is observed further offshore in the southwestern Bay of Bengal during the Northeast Monsoon (Vinayachandran and Mathew, 2003; Vinayachandran, 2009). Modeling studies suggest that this is caused by wind-driven entrainment, not only of subsurface nutrients, but also of phytoplankton from the subsurface chl_a maximum that is present during the autumn intermonsoon period (Vinayachandran et al., 2005). In contrast, productivity near the coast is suppressed during the Northeast Monsoon when the EICC flows southward (Fig. 14). Presumably, this is due to a combination of the downwelling-favorable currents and winds. However, primary production over the shelf in the northern part of the Bay increases during the Northeast Monsoon (Gomes et al., 2000; Fig. 15), possibly due to river nutrient inputs (Vinayachandran, 2009) and / or wind-stress and buoyancy-driven nutrient entrainment as is observed in the northern Arabian Sea during the Northeast Monsoon (Wiggert et al., 2000; 2005; Hood et al., 2017).

Subsurface Chl-*a* maxima are observed in the Bay of Bengal during all seasons whenever and wherever wind forcing and/or currents are insufficiently strong to upwell or entrain them into the surface layer (Sarma and Aswanikumar, 1991; Murty et al., 2000; Sarjini and Sarma, 2001; Kumar et al., 2007). During the intermonsoon periods the Bay of Bengal transitions to more oligotrophic conditions with relatively low surface chlorophyll concentrations ($\leq 0.6\text{ mg/m}^3$; Fig. 15) and production rates ($< 700\text{ mgC m}^{-2}\text{ d}^{-1}$; see Fig. 6 in Hood et al., 2017). *Trichodesmium erythraeum* blooms have been observed during the intermonsoon periods along with high abundances of *Synechococcus* and heterotrophic dinoflagellates (Sarjini and Sarma, 2001; Jyothibabu et al., 2008). In offshore waters subsurface chlorophyll maxima are generally located between 40 and 70m in autumn and 60 and 90m in spring (Kumar et al., 2007). These deep Chl-*a* maxima tend to shoal near the coast (Sarma and Aswanikumar, 1991; Murty et al., 2000) and their depth and chlorophyll concentrations are strongly influenced by eddies (Kumar et al., 2007).

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Strong upwelling also occurs along the southern coast of Sri Lanka during the Southwest Monsoon (Vinayachandran, 2004; 2009; de Vos et al., 2014). Satellite SST and chlorophyll images reveal dramatic eastward advection of cool ($< 28^{\circ}\text{C}$) chlorophyll rich upwelled water by the SMC (Vinayachandran, 2004; 2009; de Vos et al., 2014). Chlorophyll-rich waters from the southwestern coast of India are also advected by the SMC towards Sri Lanka during the Southwest Monsoon (Vinayachandran, 2004; 2009; Strutton et al., 2015). Surface chlorophyll concentrations and rates of primary production along the southern coast of Sri Lanka during the Southwest Monsoon can exceed 10 mgChl a m^{-3} (de Vos et al., 2014) and $1000\text{ mgC m}^{-2}\text{ d}^{-1}$ (Fig. 6 in Hood et al., 2017), respectively, compared to much lower concentrations and rates during the Northeast Monsoon when the NMC flows westward (de Vos et al., 2014; Hood et al., 2017). Vinayachandran (2004; 2009) attribute the productivity response during the Southwest Monsoon to nutrient enrichment from coastal upwelling driven by monsoon winds. Presumably, these high chlorophyll concentrations and production rates are associated with diatom blooms. This elevated productivity extends to the east of Sri Lanka during the peak of the Southwest Monsoon (Vinayachandran et al., 1999; Vinayachandran, 2004; 2009). This eastward extension into the southern Bay of Bengal occurs along the path of the SMC (Vinayachandran et al., 1999) and is associated with upward Ekman pumping east of Sri Lanka. This Ekman pumping also leads to the formation of the aforementioned Sri Lanka Dome (Vinayachandran and Yamagata, 1998).

The western side of the northern Indian Ocean transitions during the southwest monsoon to a eutrophic coastal upwelling system in response to the upwelling favorable winds and currents (Wiggert et al., 2005; Hood et al., 2017 and references cited therein; Fig. 5; Figs. 5 and 6 in Hood et al., 2017). These changes can be seen in ocean color data as substantial increases in chl-a concentrations along the coasts of Somalia, Yemen and Oman (e.g., Brock and McClain, 1992; Banse and English, 2000; Kumar et al., 2000; Lierheimer and Banse, 2002; Wiggert et al., 2005; George et al., 2013; Hood et al., 2017). Chlorophyll-*a* concentrations in the western Arabian Sea can exceed 40 mgChl a m^{-2} during the southwest monsoon with production rates $> 2.5\text{ gC m}^{-2}\text{ d}^{-1}$ (Marra et al. 1998; Fig. 6 in Hood et al., 2017). However, the environmental conditions vary significantly between the eutrophic coastal zones to the west and the oligotrophic open ocean waters offshore that are influenced by wind-curl induced downwelling to the southwest of the Findlater Jet (Lee et al., 2000). The surface nitrate and Chl-*a* concentrations decline dramatically from > 10 to $< 0.02\text{ }\mu\text{M}$ and from > 1.0 to $< 0.2\text{ mgChl a m}^{-3}$, respectively, from the west coast to open ocean in the Arabian Sea (Brown et al., 1999; Wiggert et al., 2005; Hood et al., 2017). In general, the phytoplankton community structure transitions to larger cells (diatoms) during the southwest monsoon in the western Arabian Sea (Brown et al., 1999; Tarran et al., 1999; Shalapyonok et al., 2001). However, small primary producers remain important, even in areas strongly influenced by coastal upwelling (Brown et al., 1999). In contrast, during the oligotrophic spring and fall intermonsoon periods, surface waters in the western Arabian Sea are dominated by picoplankton (Garrison et al., 2000). Subsurface Chl-*a* maxima are observed between 40 and 140 meters in the central southeastern Arabian Sea during all seasons (Gunderson et al., 1998; Goericke et al., 2000; Ravichandran et

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al., 2012), at times occurring in layers below the oxyclines of the oxygen minimum zone (Georick et al., 2000). These features are strongly influenced by mesoscale features (Gundersen et al., 1998).

During the southwest monsoon off Oman and Somalia, the presence of the topographically-locked eddies generate strong offshore flows that advect high nutrient, high Chl-*a* concentrations and coastal phytoplankton communities hundreds of kilometers offshore (Keen et al., 1997; Latasa and Bidigare, 1998; Manghnani et al., 1998; Gundersen et al., 1998; Hitchcock et al., 2000; Lee et al., 2000; Kim et al., 2001). These advective effects can be seen, for example, in association with the Great Whirl off the coast of northern Somalia (Hitchcock et al., 2000) and in the filaments that develop off the Arabian Peninsula during the southwest monsoon (Wiggert et al. 2005; Hood et al., 2017). In contrast, during the northeast monsoon, the circulation and winds transition to downwelling favourable. During the northeast monsoon, cold dry northeasterly winds from southern China and the Tibetan Plateau flow across the northern Arabian Sea. The shear from these winds, combined with surface cooling and buoyancy-driven convection, drive mixing and entrainment of nutrients that, in turn, promote modest increases in chlorophyll and primary production over the northern Arabian Sea (Wiggert et al., 2000; Wiggert et al., 2005; Fig. 5; Figs. 5 and 6 in Hood et al., 2017). These increases in Chl-*a* have been associated with increased diatom abundance (Banse and McClain, 1986; Sawant and Madhupratap, 1996). In the last decade, however, there appears to have been a shift in the composition of winter phytoplankton blooms in the northern and central Arabian Sea from diatom dominance to blooms of a large, green mixotrophic dinoflagellate, *Noctiluca scintillans* (Gomes et al., 2014; Goes et al., 2020).

During the southwest monsoon, the upwelling-favorable WICC induces upwelling along the west coast of India, which increases Chl-*a* concentrations by more than 70% compared to the central Arabian Sea (Kumar et al., 2000; Naqvi et al., 2000; Luis and Kawamura, 2004; Hood et al., 2017). The increased Chl-*a* concentrations near the coast are associated with increases in diatom abundance (Sawant and Madhupratap, 1996). However, these increases in Chl-*a* and their offshore extent are modest compared to the western Arabian Sea (Fig. 5; Fig. 5 in Hood et al., 2017). In contrast, during the northeast monsoon the WICC is downwelling-favorable and tends to suppress primary production off the southwestern coast of India. The depletion of nutrients in this region during the northeast monsoon coincides with blooms of *Trichodesmium* and dinoflagellate species (Parab et al., 2006; Matondkar et al., 2007) resulting in the extremely high rates of nitrogen fixation (Gandhi et al., 2011, Kumar et al., 2017). However, as discussed above, further north and offshore, nutrient entrainment enhances phytoplankton biomass and primary production during the northeast monsoon (Wiggert et al., 2000; McCreary et al., 2001; Luis and Kawamura, 2004; Gomes et al., 2014; Goes et al., 2020; Fig. 5). Near-surface Chl-*a* and primary production off the west coast of India (estimated from satellite ocean color measurements) increases from ~9 to 24 mgChl *a* m⁻² and from ~1 to 2.25 g C m⁻² d⁻¹, respectively, from winter to the summer monsoon (Luis and Kawamura, 2004; Fig. 5; Figs. 5 and 6 in Hood et al., 2017). The elevated productivity during the southwest monsoon is

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modulated by the coastal Kelvin waves that originate from the Bay of Bengal and propagate along the West Indian Shelf, modifying circulation patterns and upwelling (Luis and Kawamura, 2004).

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5 Inter-ocean exchange

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Satellite ocean color observations suggest that Chl-*a* concentrations remain low all year-round in the Laccadive Sea in spite of the transition from upwelling circulations (anti-cyclonic) during the southwest monsoon to downwelling circulations (cyclonic) during the northeast monsoon (Lierheimer and Banse, 2002; Fig. 5[6]). This suggests that the upwelling-enhanced productivity off the west coast of India during the southwest monsoon is largely restricted to the coastal zone. There are, however, sporadic zonal increases in Chl-*a* concentration that extend into the Laccadive Sea (Lierheimer and Banse, 2002). It is not known whether or not there are cryptic subsurface phytoplankton blooms during the southwest monsoon associated with the Laccadive Low.¶

5.1 Indonesian Throughflow

5.1.1 General features

The Indonesian Throughflow (ITF) transfers low-salinity tropical waters from the Pacific to the Indian Ocean via the Indonesian seas (Fig. 10). The ITF is the only tropical oceanic pathway that links ocean basins and plays an important role in the global ocean circulation and climate system (Sprintall et al., 2014; 2019). The simultaneous measurements in the exit channels of the ITF from the International Nusantara Stratification and Transport (INSTANT) program during 2004–2006 (Gordon et al., 2008; Sprintall et al., 2009) suggested that the ITF has a mean transport of 15 Sv into the Indian Ocean. The ITF pathway is composed of many narrow channels within the Indonesian seas, among which about 80% of the total ITF is through the Makassar Strait (Fig. 10, Gordon et al., 2008, 2010). The remaining passages include the Maluku Sea, Lifamatola Passage, Karimata Strait and Sibutu Passage (Fang et al., 2010; Gordon et al., 2012; Susanto et al., 2013).

5.1.2 Variability, dynamics and influence

The interannual variability of the ITF is mainly dictated by the ENSO-related wind forcing through the Pacific waveguide with stronger transport during La Niña years (Meyers, 1996; England and Huang, 2005; Hu and Sprintall, 2016), but the IOD occasionally offsets the Pacific ENSO influences through the Indian Ocean wind variability and Indian Ocean waveguide (Sprintall and Révelard, 2014; Liu et al. 2015; Feng et al., 2018). For the strong negative IOD event in 2016, the Indian Ocean influence overwhelmed that of the Pacific leading to record low ITF volume transports because of the reduction in the interbasin pressure gradient (Pujiana et al., 2019). Strong wind forcing over the equatorial Indian Ocean triggers equatorial Kelvin waves and influences the ITF variability on intraseasonal, semi-annual and interannual time scales (Drushka et al., 2010; Pujiana et al., 2013; Shinoda et al., 2012). Kelvin waves through the Indian Ocean waveguide are suggested to influence the interannual variability in the tropical Pacific Ocean (Yuan et al., 2013; Pujiana and McPhaden, 2020).

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The ENSO cycle also influences the outflowing ITF transport through the salinity effect in the downstream buoyant pool, contributing about 36% of the total ITF interannual transport variation (Hu and Sprintall, 2016; Section 6.1). Fresh anomalies in the buoyant pool during La Nina years can be as large as 0.2 in practical salinity averaged over the upper 180

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m of the water column (Phillips et al. 2005). Such salinity anomalies can strengthen the volume transport of the LC through an increase in the zonal density gradient driving stronger southward flow (Feng et al., 2015a). The Inter-decadal Pacific Oscillation/Pacific Decadal Oscillation (IPO/PDO), through modulations of decadal wind stress in the tropical Pacific, has also directly influenced the strength of the ITF (Feng et al., 2011; Hu et al., 2015; Mayer et al., 2018). This has, in turn influenced heat and freshwater transports, causing upper ocean heat content to increase in the southern Indian Ocean (Feng et al., 2010; Schwarzkopf and Böning, 2011; Nidheesh et al., 2013; Sprintall, 2014; Lee et al., 2015; Nieves et al., 2015; Du et al., 2015; Ummenhofer et al., 2017) and produced interhemispheric contrasts in sea surface temperature (Dong and McPhaden, 2016). During the negative IPO phase, such as during the hiatus in warming of the globally averaged surface atmosphere (1998-2012), enhanced trade winds in the Pacific strengthened the ITF volume and heat transport into the Indian Ocean, driving a rapid warming trend in the Southern Indian Ocean (England et al., 2014; Nieves et al., 2015; Lee et al., 2015; Liu et al., 2015; Zhang et al., 2018). Contributions from air-sea exchanges (Jin et al. 2018a,b) have also been suggested to be important, as has a reduction in the oceanic heat exported from the Indian Ocean at its southern boundary (Lisa Beal, personal communication).

Using a combination of theory, ocean reanalyses, OGCM simulations, and coupled climate model simulations, Jin et al. (2018a,b) found eastern and western Indian Ocean heat content to be affected by remote Pacific forcing through two distinct mechanisms: oceanic influences transmitted through the ITF and the atmospheric bridge. The intensified freshwater input within the Maritime Continent during the past decade was found to strengthen the ITF and its heat and freshwater transports into the Indian Ocean, causing significant warming and freshening trends and accelerated sea-level rise in the eastern Indian Ocean (Hu and Sprintall, 2017a, 2017b; Zhang et al., 2018; Jyoti et al., 2019). The decadal enhancement of the ITF transport has increased upper ocean heat content anomalies in the southeast Indian Ocean and increased the likelihood of marine heatwaves off the west coast of Australia (Feng et al., 2015b; Section 6.4).

5.2 Agulhas Leakage

5.2.1 General features

At the tip of Africa, the southward-flowing Agulhas Current retroflects with most of the flow heading eastwards along the northern edge of the ACC, recirculating back into the Indian Ocean (Fig. 10, Section 4.2.2). Around 20-30% of the Agulhas Current enters the Atlantic Ocean as Agulhas leakage in the form of Agulhas rings and cyclones (van Sebille, 2010a). Agulhas leakage estimates are sensitive to the definition used to calculate the leakage, ranging roughly between 10 and 20 Sv (van Sebille et al., 2010b; Beron-Vera et al., 2013; Cheng et al., 2016; Holton et al., 2017). Bars et al. (2014) proposed an algorithm to measure Agulhas leakage anomalies using absolute dynamic topography data from satellites.

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Deleted: The Agulhas Current, the western boundary current of the southern Indian Ocean gyre, (Fig. 10), flows southward along the South African coastline.

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2508 The division of flow between Agulhas Leakage and Agulhas retroflection can be influenced by the upstream Agulhas
 2509 Current. In a Lagrangian particle tracking experiment, van Sebille et al. (2009) found that a weaker Agulhas Current,
 2510 detaching farther downstream and generating anti-cyclonic vorticity, potentially leads to more Agulhas leakage and larger
 2511 Indian-Atlantic inter-ocean exchange. However, eddy-resolving model results suggest that as model resolution increases,
 2512 the sensitivity of the leakage to Agulhas Current transport anomalies is reduced (Loveday et al., 2014). In addition, the
 2513 ITF potentially influences the Agulhas leakage (Le Bars et al., 2013), as model outputs suggest that the Indian Ocean
 2514 contributes 12.6 Sv to the Agulhas leakage, half of which is from the ITF_v (Durgadoo et al., 2017).

2515 5.2.2 Variability, dynamics and influence on climate

2516 The magnitude of the Agulhas leakage is controlled by wind forcing including the trade winds and the Southern
 2517 Hemisphere Westerlies (e.g., Durgadoo et al., 2013). The poleward shift in the Southern Hemisphere westerlies associated
 2518 with anthropogenic forcing induced a clear increase in the Agulhas leakage during 1995-2004 as shown in numerical
 2519 simulations (Biaostoch et al., 2009; Biaostoch and Böning, 2013). Increased wind stress curl in the South Indian Ocean
 2520 associated with the southward shift of westerlies led to significant warming in the Agulhas Current system since the 1980's
 2521 (Rouault et al., 2009); however further work showed that this is due to an increase in eddies leading to a broadening of the
 2522 current as opposed to intensification (Beal and Elipot, 2016). Given the non-linear nature of Agulhas leakage, the difficulty
 2523 of observing it and ocean model biases in the region, quantifying Agulhas leakage is very challenging (Holton et al., 2017).
 2524 At seasonal time scales, the Agulhas leakage variability is controlled by eddies, however recent studies have shown that
 2525 eddies might not contribute as significantly to leakage as was thought and the non-eddy leakage transport is likely to be
 2526 constrained by large-scale forcing at longer time scales (e.g., Cheng et al., 2018). A recent study shows that the subsurface
 2527 signal from the ENSO cycle influences the Agulhas leakage through Rossby waves with a time lag of 2 years (Paris et al.,
 2528 2018).

2529 The Agulhas leakage carries warm and saline water from the Indo-Pacific Ocean into the Atlantic Ocean. The Agulhas
 2530 leakage has been suggested to influence the Atlantic Meridional Overturning Circulation strength (AMOC; Beal et al.
 2531 2011; Weijer and van Sebille, 2014; Biaostoch et al. 2015) and modify the AMOC convective stability (e.g., Haarsma et
 2532 al., 2011; Caley et al., 2012; Castellanos et al., 2017). It is suggested that the increases in the Agulhas leakage due to
 2533 anthropogenic warming during the past decades would act to strengthen the Atlantic overturning circulation (e.g., Beal et
 2534 al., 2011).

2535 The Agulhas leakage is an important source of decadal variability in the AMOC through Rossby waves (Biaostoch et al.,
 2536 2008; 2015). Source waters from the Agulhas Current take more than four years and mostly one to four decades to arrive
 2537 in the North Atlantic Ocean (van Sebille et al., 2011; Rühls et al., 2013). The increased Agulhas leakage during 1995-2004
 2538 has contributed to the salinification of the South Atlantic thermocline waters (Biaostoch et al., 2009). Hindcast experiments

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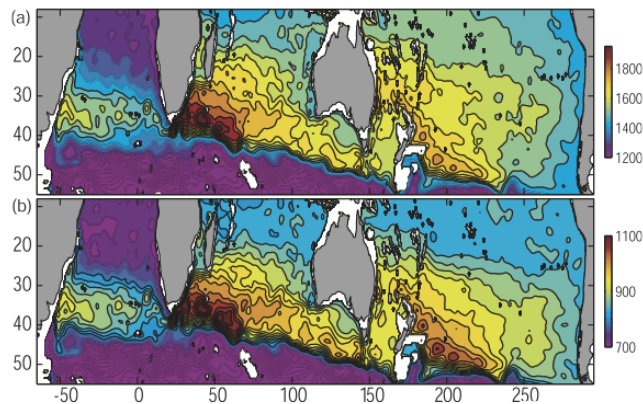
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2547 suggest that the Agulhas leakage **increased** by about 45% during **the 1960s-2000s**, **leading** to the observed warming trend
2548 in the upper tropical Atlantic Ocean (Lübbecke et al., 2015).

2549 5.3 Supergyre connection to the South Pacific

2550 The extreme strong westerly wind stress in the Southern Hemisphere gives rise to a wide and energetic subtropical
2551 supergyre (Figure 16), the Southern Hemisphere supergyre, **that connects three ocean basins** (e.g., Ridgway and Dunn,
2552 2007; Speich et al., 2007; Lambert et al., 2016; Maes et al., 2018; Cessi, 2019). Although the near-surface circulation is
2553 eastward across the **southern Indian Ocean**, there are subsurface westward flows beneath (Section 4.2.3; Schott and
2554 McCreary 2001; Domingues et al. 2007; Furue et al. 2017), and the depth-integrated circulation reveals the westward
2555 return flow of the equatorward side of the **Indian Ocean's** anti-clockwise subtropical gyre. In Figure 16, the southern side
2556 of the Indian Ocean subtropical gyre extends eastward south of Australia to connect with the western Pacific subtropical
2557 gyre. The return flow is accomplished via a pathway that includes the East Australian Current, the South Pacific's western
2558 boundary flow; the Tasman Leakage, a westward flow south of Tasmania that carries Pacific Ocean water back to the
2559 Indian Ocean (distinct from the Flinders Current that hugs the continental slope, Duran et al 2020; Section 4.2.4); and
2560 northwestward flow in the eastern Indian Ocean to close the circulation. **The ITF and Leeuwin Current are also part of the**
2561 **supergyre, connecting the Indian and Pacific Oceans through the Indonesian seas (e.g. Ridgway and Dunn, 2007).**



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Moved down [49]: Figure 16: The interbasin supergyre system for the Pacific and Indian Oceans as shown by the depth-

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Figure 16: The interbasin supergyre system for the Pacific and Indian Oceans as shown by the depth-integrated steric height (a) $P_{0/2000}$, and (b) $P_{400/2000}$, derived from the CARS climatological temperature and salinity fields. The contour interval in (a) is 50 m^2 and in (b) is 25 m^2 . Taken from Ridgway and Dunn 2007.

The supergyre is the subtropical gyre of the southern hemisphere. As such, its flow is primarily determined by the westward integration of wind stress curl from the eastern boundaries as determined by Sverdrup dynamics. The latitudinal position of the Subtropical Front at the southern edge of the supergyre is found to be controlled by strong bottom pressure torque due to the interaction between the ACC and the ocean floor topography (De Boer et al., 2013). According to one analysis in SODA (Simple Ocean Data Reanalysis), the water masses in the supergyre became cooler and fresher and shifted southward by about 2.5° due to changes in the basin-scale wind forcing during 1958–2007 (Duan et al., 2013). A recent study using altimeter observations shows a clear strengthening of the Southern Hemisphere supergyre in all three oceans since 1993 as indicated in the large trends of sea surface height and their contrast. Argo observations and ECCO assimilations suggest that the strengthening extends to deeper than 2000 m (Qu et al., 2019). The spin-up of the Southern Hemisphere supergyre is attributed to the poleward shift and strengthening of westerly winds that are linked to an increasingly positive southern annular mode (Qu et al., 2019).

5.4 Roles of salinity in inter-ocean exchange

Ocean salinity is one of the basic variables that determines the oceanic stratification, sea level change and climate change (e.g., Llovel and Lee, 2015; Kido and Tozuka, 2017; Sprintall et al., 2019). However, the role of salinity in ocean circulation has been largely underestimated until the recent decade when *in situ* observations of subsurface and surface salinity from Argo and satellite salinity missions became available. These new observations have revolutionized our understanding of the influence of salinity on ocean circulation and dynamics (Vinogradova et al. 2019, and references therein).

Four major processes control the salinity in the Indian Ocean: net air-sea fluxes (evaporation minus precipitation), freshwater inflow from large rivers in the Bay of Bengal, inflow of relatively fresh waters from the Pacific Ocean via the Indonesian Throughflow, and inflow of saltier waters from the Red Sea and the Persian and Arabian Gulfs. These different drivers combine to give the Indian Ocean salinity its unique flavour: a strong east-west gradient in the North Indian Ocean (salty in the Arabian Sea and fresh in the Bay of Bengal) and strong north-south gradients in the South Indian Ocean (fresh in the tropics, and salty in the subtropics) (Fig. 4).

Salinity is a crucial variable to understand Indian Ocean dynamics. For instance, salinity has strong ties with the Indian Ocean Dipole (e.g., Du and Zhang, 2015; Durand et al., 2013; Grunseich et al., 2011; Kido and Tozuka, 2017; Nyadjiro and Subrahmanyam, 2014; Zhang et al. 2016; Section 6.2), the EGC (Menezes et al., 2013; Section 4.2.3), LC transport,

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2669 Ningaloo Niño and marine heatwaves off western Australia (e.g., Feng et al., 2015a), and the El Niño/La Niña climate
 2670 mode (e.g., Hu and Sprintall, 2016; Zhang et al., 2016). Salinity plays an essential role in the dynamics of the seasonal
 2671 Wyrski Jets in the equatorial zone (e.g., Masson et al., 2003), extra-equatorial Rossby waves (Heffner et al., 2008; Menezes
 2672 et al., 2014b; Vargas-Hernandez et al., 2015; Banks et al., 2016), Madden-Julian and Intraseasonal Oscillations (e.g.,
 2673 Grunseinch et al., 2013; Guan et al., 2014; Subrahmanyam et al., 2018), barrier-layer dynamics (e.g., Drushka et al., 2014;
 2674 Felton et al., 2014), and the North Indian Ocean (e.g., D’Addezio et al., 2015, Fournier et al., 2017; Mahadevan et al.,
 2675 2016; Nyadjro et al., 2011, 2012, 2013; Wilson and Riser, 2016; Spiro Jaeger and Mahadevan, 2018).
 2676
 2677 Salinity variability within the Indonesian Seas has been shown to control the transport of the ITF. Andersson and
 2678 Stigebrandt (2005) proposed that a downstream buoyancy pool in the outflowing ITF region acts to regulate the ITF
 2679 transport. Gordon et al. (2003, 2012) pointed out that low salinity surface water from the South China Sea is drawn into
 2680 the Java Sea. Combined with the monsoonal precipitation over the Maritime Continent and seasonal monsoon winds, this
 2681 freshwater plug contributes to the seasonal fluctuation of the Makassar Strait Throughflow transport and inhibits the inflow
 2682 of tropical Pacific surface water from the Mindanao Current (e.g., Gordon et al., 2012; Lee et al., 2019). Recently, Hu and
 2683 Sprintall (2016) found that about 36% of the interannual ITF transport is attributable to the salinity effect associated
 2684 with freshwater input anomalies due to the ENSO cycle. Jyoti et al. (2019) further examined this salinity effect and found
 2685 that the unprecedented sea-level rise in the southern Indian Ocean since the beginning of the 21st Century is attributed to
 2686 the accelerated heat and freshwater intrusion by the ITF. A significant strengthening of the ITF transport in the 2000s has
 2687 given rise to a subsequent warming and freshening of the eastern Indian Ocean (e.g., Hu and Sprintall, 2017a, 2017b,
 2688 Section 6.1). The southeast Indian Ocean is one of the few places in the global ocean where the halosteric component of
 2689 sea level rise is as large as the thermosteric component (Llovel and Lee, 2015).
 2690

2691 6 Modes of Interannual Climate Variability in the Indian Ocean

2692 6.1 ENSO teleconnection and the Indian Ocean Basin mode

2693 ENSO influences the Indian Ocean circulation through the Pacific-to-Indian Ocean oceanic waveguide and atmospheric
 2694 teleconnections. Through the atmospheric bridge, El Niño conditions in the Pacific induce an anticyclonic wind anomaly
 2695 pattern in the southeast Indian Ocean (Xie et al., 2002), whereas La Niña induces a cyclonic wind anomaly pattern (Feng
 2696 et al., 2013). The ENSO teleconnection also drives SST variability over the western Indian Ocean during ENSO
 2697 development. The tropical Indian Ocean experiences prolonged warming (cooling) that peaks in the following boreal
 2698 spring and persists into boreal summer, after the decay of El Niño (La Niña) events, the so-called Indian Ocean Basin
 2699 (IOB) mode (Yang et al., 2007). The westward propagating Rossby waves induced by ENSO may also help sustain the

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Deleted: (2019) further examined the salinity effect in the ITF transport following Hu and Sprintall (2016) and found that the unprecedented sea-level rise in the southern Indian Ocean since the beginning of the 21st Century is attributed to the accelerated heat and freshwater intrusion by the ITF.

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2712 warming (cooling) of the tropical Indian Ocean (Xie et al., 2002), fueled by regional air-sea coupling (Du et al., 2009).
2713 The IOB warming has a capacitor effect for El Niño to influence boreal summer climate, such as for the Indian monsoon
2714 (Zhou et al., 2019), and remote impacts in the northwest Pacific (Xie et al., 2009, 2016), including China and Japan (Hu
2715 et al., 2019). Details of the Indo-Western North Pacific capacitor effect are summarized in Xie et al. (2016) and Kosaka et
2716 al. (2021). The relationship between ENSO and IOB evolves on a decadal time scale, and the persistent IOB warming
2717 after El Niño has been evident since the 1970s (Xie et al., 2010). Based on the CMIP5 multi-model experiments, the IOB
2718 warming tends to persist longer after the El Niño events under global warming scenarios (Zheng et al., 2013).

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2719 The ITF variability lags ENSO by 8-9-months, found in ocean model results (England and Huang, 2005) and derived from
2720 the geostrophic transport across an Australia-Indonesia XBT section (Liu et al., 2015). The variability of the ITF transport
2721 drives sea level and upper ocean heat content anomalies in the southeast Indian Ocean. Through the waveguide, ENSO
2722 has a direct influence on the strength of the Leeuwin Current (Section 4.2.4), with a stronger poleward volume and heat
2723 transport during a La Niña event (Feng et al., 2008). A stronger Leeuwin Current during La Niña events leads to greater
2724 baroclinic instability of the current and enhanced generation of eddies that leads to interannual variability of the eddy
2725 kinetic energy in the southeast Indian Ocean (Feng et al., 2005; Zheng et al., 2018). The increase of the ITF transport and
2726 enhancement of rainfall in the Indonesian Seas during strong La Niña events can drive up to 0.2-0.3 psu freshening
2727 anomalies in the upper southeast Indian Ocean (Phillips et al., 2005; Feng et al., 2015a; Hu and Sprintall, 2017a; Section
2728 5.1.2), which may have a compound effect in accelerating the Leeuwin Current (Feng et al., 2015a). Both ENSO and the
2729 IOD (see Section 6.2) influence the ITF and thus the exchange of heat from the Pacific into the Indian Ocean, but in
2730 concurrent IOD and ENSO events it appears that the influence from the IOD dominates (Sprintall and Revelard, 2014).

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2731 Due to the opposing effects of the winds and dissipation, ENSO induced sea level and upper ocean heat content anomalies
2732 in the southeast Indian Ocean do not propagate far into the western Indian Ocean; instead, wind anomalies generate sea
2733 level and heat content anomalies of opposite signs in the western Indian Ocean through Rossby wave propagations
2734 (Masumoto and Meyers, 1998; Xie et al., 2002; Zhuang et al., 2013; Ma et al., 2019; Volkov et al., 2020; Nagura and
2735 McPhaden, 2021). Thus, the joint forcing of the oceanic waveguide and atmospheric teleconnection results in variations
2736 of meridional overturning circulation and heat transport in the Indian Ocean on a multi-year time scale, in phase with the
2737 ITF variability (Ma et al., 2019).

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2738 6.2 The Indian Ocean Dipole

2739 There is increasing evidence that positive IOD events are more frequent and intense during the 20th century (e.g., Abram
2740 et al., 2008; Cai et al., 2013; Abram et al., 2020a,b; and references therein). A rare occurrence of three consecutive positive
2741 IOD events took place in 2006-2008 (Cai et al., 2009b). The skewness towards more positive and fewer negative IOD
2742 events (Cai et al., 2009a) is due potentially to an anthropogenically-driven shoaling thermocline in the eastern Indian

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2748 Ocean (Cai et al., 2008). The three consecutive positive IOD events rarely occurred in Coupled Model Intercomparison
2749 phase 5 (CMIP5) models and the more recent frequent occurrence was consistent with regional Indo-Pacific Walker
2750 circulation trends (Cai et al., 2009c,d). An anthropogenic contribution was proposed since positive IOD events became
2751 more frequent over the period 1950–1999 in the CMIP5 models. Projected mean-state changes in the Indian Ocean with
2752 stronger easterly winds and a shoaling thermocline in the southeast Indian Ocean during austral spring favour positive
2753 IOD development, with a reduction in skewness between positive and negative IOD events likely (Cai et al., 2013; Figure
2754 17), and a three-fold increase in frequency of extreme positive IOD events by 2100 compared to the previous century (Cai
2755 et al., 2014a). However, model biases in Indian Ocean mean-state and IOD variability challenge these projected changes:
2756 models with excessive IOD amplitude bias tend to project a strong IOD-like warming pattern and increase in extreme
2757 pIOD occurrences, consistent with an enhanced Bjerknes feedback, and hence the projected IOD changes could represent
2758 spurious artefacts of model biases (Li et al., 2016). Yet, paleoclimate evidence supports trends observed in recent decades:
2759 based on a millennial IOD reconstruction from corals, extreme positive IOD events, as were observed in 1997 and 2019,
2760 were historically rare (Abram et al., 2020b). In the reconstruction, only ten extreme positive IOD events occurred and yet
2761 four events occurred in the last 60 years (Abram et al., 2020b). The increase in event frequency and intensity highlights
2762 the need to improve preparedness in regions affected by IOD events to minimize future climate risks posed by them.

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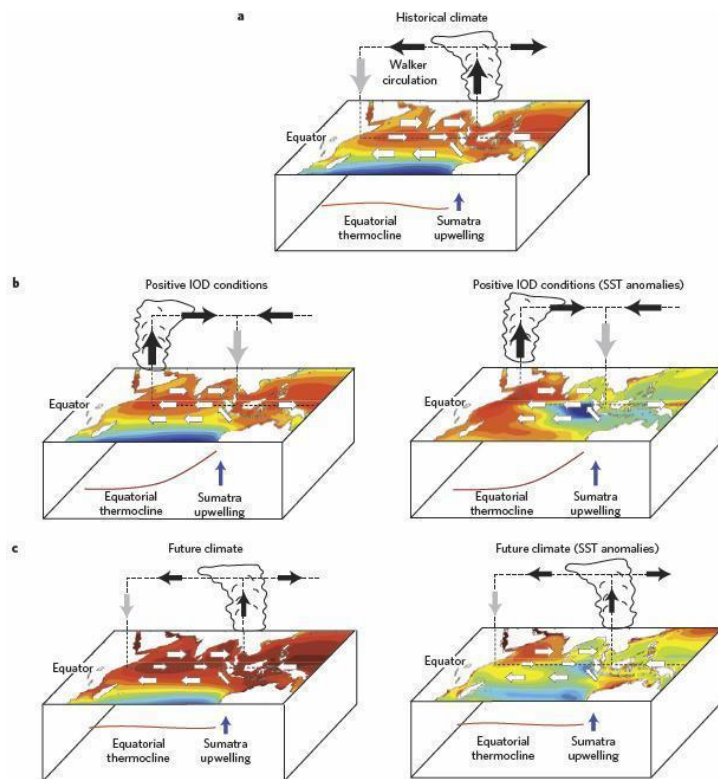
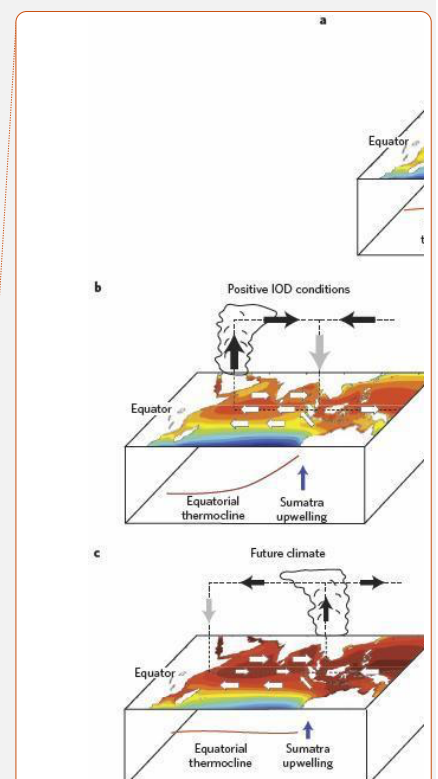


Figure 17: Historical austral spring mean climate and positive IOD conditions for the twentieth century, and future austral spring mean climate. a, Historical mean climate, indicating SSTs, surface winds, the associated atmospheric Walker circulation, the mean position of convection and the thermocline. In the western Indian Ocean, the descending branch is broad and not well-defined, as indicated by a grey arrow. b, Typical conditions during a positive IOD event. c, Projected future mean climate based on a CMIP5 multi-model ensemble average. Diagrams with total SST fields are shown on the left; diagrams with SST anomalies referenced to the 1961–1999 mean for b,



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2773 and referenced to the basin mean for c, are shown on the right. Reprinted from Cai et al. (2013) with permission
2774 from Springer Nature.

2775 While model simulations and paleo proxy records suggest changes in the frequency and magnitude of IOD events in a
2776 warming climate, there is less observational evidence from other sources. Given the short observational record in the
2777 Indian Ocean, the role of decadal to multi-decadal variability across the broader Indo-Pacific region has recently emerged
2778 as a compounding factor: the number and frequency of IOD events have been observed to vary on decadal timescales.
2779 Decadal variations in SST featuring an IOD-like out-of-phase pattern between the western and eastern tropical Indian
2780 Ocean have been linked to the PDO (Krishnamurthy and Krishnamurthy, 2016) or IPO (Dong et al., 2016). A combination
2781 of processes transmits the signal from the Pacific to the Indian Ocean through both the atmospheric and oceanic bridges,
2782 leading to variations in the subsurface temperature structure in the Indian Ocean (Zhou et al., 2017; Jin et al., 2018a).
2783 Decadal modulations of the background state of the eastern Indian Ocean thermocline depth can thus pre-condition the
2784 Indian Ocean to more or less IOD events (Annamalai et al., 2005). Consequently, positive IOD events were unusually
2785 common in the 1960s and 1990s with a relatively shallow eastern Indian Ocean thermocline, while the deeper thermocline
2786 in the 1970s and 1980s was associated with frequent negative IOD and rare positive IOD events (Ummenhofer et al.,
2787 2017). The Indian Ocean stands out as a region with high skill in decadal predictions (Guemas et al., 2013) and improved
2788 understanding of decadal modulation of IOD events can aid in decadal prediction efforts for the Indian Ocean region.

2789 The relationship between ENSO and the IOD has been subject to ongoing debate. Recent research has shown that around
2790 two-thirds of IOD variability arises as a remote response to ENSO (Stuecker et al., 2017; Yang et al., 2015), with the
2791 remaining variability being independent of ENSO. Stuecker et al. (2017) argue that the ENSO-driven IOD can be seen as
2792 a combination of remotely driven wind and heat flux anomalies modulated by seasonally-varying Bjerknes feedback in
2793 the Indian Ocean. Further, they suggest that the ENSO-independent IOD events arise out of white noise atmospheric
2794 forcing coupled to these feedbacks (Stuecker et al., 2017). Variability internal to the Indian Ocean basin and unrelated to
2795 ENSO, arising from ocean-atmosphere feedback processes, does however modulate the evolution of IOD events and can
2796 lead to early termination of IOD events; as a result, including internal variability improves the predictability of the IOD
2797 (Yang et al., 2015). IOD variability internal to the Indian Ocean resembles recharge oscillator dynamics for ENSO, but
2798 equatorial heat content is less effective as a precursor for the IOD than for ENSO because of the strong impact of remote
2799 forcing from the Pacific on the IOD. Internal Indian Ocean dynamics however may contribute to the biennial nature of
2800 the IOD through the cycling of Kelvin/Rossby wave energy across the basin (McPhaden and Nagura, 2014). The
2801 relationship between ENSO and the IOD is not only one-way: IOD events have also been shown to influence the
2802 development of ENSO in the following year (Izumo et al., 2010; Wang et al., 2019; Cai et al., 2019; and references therein).

2803 Different types of IOD events have been described, each with distinct evolution and regional impacts (Du et al., 2013;
2804 Endo and Tozuka, 2016). Du et al. (2013) distinguished three types of IOD events according to the timing of their peak

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2805 amplitude and overall duration: ‘unseasonable’ events that develop and mature mostly within June-August (JJA), ‘normal’
2806 events that develop and mature mostly within September-November (SON), and ‘prolonged’ events that develop in JJA
2807 and mature in SON, with the latter two described as the canonical IOD events (Du et al., 2013). The unseasonable IOD
2808 events have only been observed since the mid-1970s and have been suggested to be a response to the rapidly warming
2809 Indian Ocean SST and a weakened Walker circulation during austral winter (Du et al., 2013). The seasonal evolution and
2810 type of ENSO also seems to play a role in determining the IOD evolution and type, with atmospheric influences transmitted
2811 through variations in the Walker Circulation and oceanic ones through anomalous oceanic Rossby waves affecting timing
2812 and evolution of IOD events, especially during their developing phase (Guo et al., 2015; Zhang et al. 2015; Fan et al.,
2813 2017). However, Sun et al. (2015) suggested more IOD events independent of ENSO since the 1980s, along with higher
2814 correlations between the IOD and Indian summer monsoon activity, likely due to mean-state change in the tropical Indian
2815 Ocean due to weaker equatorial westerlies. The relationship between ENSO and the IOD has weakened in recent decades,
2816 linked to changes in the ENSO-induced rainfall anomalies over the Maritime Continent (Han et al., 2017).

2817 Recent advances in understanding variability and change in IOD characteristics have implications for the relationships
2818 between SST and regional rainfall patterns in Indian Ocean rim countries. For example, different types of IOD events
2819 exhibit distinct regional impacts, with only the canonical events associated with enhanced rainfall over East Africa due to
2820 the low-level moisture convergence over the region (Endo and Tozuka, 2016). The effect of Indian Ocean SST on East
2821 African rainfall is most pronounced during the short rains (September-November), though Williams and Funk (2011)
2822 argued that warming Indian Ocean SST in recent decades ~~was~~ also associated with reduced long rains for the March-June
2823 season in Ethiopia and Kenya. Changes in the tropical atmospheric circulation across the Indo-Pacific on multi-decadal
2824 timescales (Vecchi and Soden, 2007; L’Heureux et al., 2013) have further implications for the relationship between Indian
2825 Ocean SST and regional rainfall: When the Pacific Walker cell weakened and the Indian Ocean one strengthened post-
2826 1961, the East African short rains became more variable and wetter (Nicholson, 2015). Similarly, Manatsa and Behera
2827 (2013) described an epochal strengthening in the relationship between the IOD and East African rainfall post-1961, with
2828 73% of short rain variability in East Africa explained by the IOD, up from 50% in previous decades. After 1997, this
2829 increased further to 82%, explaining spatially coherent events across the region and frequent rainfall extremes (Manatsa
2830 and Behera, 2013). Recent observed and projected changes in frequency and intensity of IOD events highlight the
2831 increasing need for preparedness in vulnerable regions affected by these events. One such event is the recent 2019 positive
2832 IOD, the largest Indian Ocean Dipole on record since the 1960s (Du et al. 2020), which caused extreme rainfall and floods
2833 over Japan and China (Takaya et al. 2020; Zhou et al., 2021). The 2019 IOD was unique in that it developed independently
2834 from any El Nino events and resulted from westward propagating Rossby waves in the southwest tropical Indian Ocean
2835 (Du et al., 2020).

2836 6.2.1 Biogeochemical Variability

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2838 IOD events are associated with distinct changes in primary productivity, as measured by chlorophyll. During positive IOD
2839 events, increased chlorophyll indicative of phytoplankton blooms is apparent in the normally oligotrophic eastern Indian
2840 Ocean in fall (Wiggert et al., 2009; Currie et al., 2013). Positive chlorophyll anomalies occur in the southeastern Bay of
2841 Bengal in boreal winter, while negative anomalies are observed over much of the Arabian Sea and southern tip of India.
2842 In a case study of the 2006 positive IOD event, Iskandar et al. (2010) using an eddy-resolving biophysical model found
2843 the offshore chlorophyll signal in the southeastern Indian Ocean to be associated with regions of high eddy kinetic energy
2844 implying that cyclonic eddies injected nutrient-rich water into the upper layer enabling the bloom. Currie et al. (2013)
2845 emphasize the importance of assessing the relative contributions of IOD events and remote impacts from ENSO on primary
2846 productivity in the Indian Ocean through their respective influence on upper-ocean properties for improved understanding
2847 and ultimately predictions of productivity, ecosystems, and fisheries within the basin. Little attention has been paid so far
2848 to resultant effects of these blooms on biogeochemical cycling (Wiggert et al., 2009).

2849 6.3 The subtropical Indian Ocean Dipole

2850 The subtropical Indian Ocean Dipole (SIOD) is a climate mode in the southern Indian Ocean, which tends to arise and
2851 peak in the austral summer (Behera and Yamagata, 2001). During the SIOD's positive phase, the climate mode has positive
2852 SST anomalies in the southwestern Indian Ocean and negative SST anomalies in the northeastern region (Behera and
2853 Yamagata, 2001; Suzuki et al., 2004; Hermes and Reason, 2005). During the positive phase, enhanced precipitation occurs
2854 over southern Africa (Behera and Yamagata 2001; Reason 2001, 2002). Recent studies have shown that the SIOD affects
2855 the Indian summer monsoon rainfall (Terray et al., 2003), rainfall over southwestern Australia (England et al., 2006) and
2856 tropical cyclone trajectories in the southern Indian Ocean (Ash and Matyas, 2012).

2857 Initially, SST anomalies associated with the SIOD were considered to be generated directly by latent heat flux anomalies
2858 (Behera and Yamagata, 2001). However, recent studies (Morioka et al. 2010, 2012) based on a mixed layer heat budget
2859 analysis revealed the importance of mixed layer depth anomalies generated by latent heat flux anomalies. Wind anomalies
2860 associated with the anomalous Mascarene High suppress latent heat loss and shoal the mixed layer in the southwestern
2861 part, while latent heat release is enhanced and the mixed layer deepens anomalously in the northeastern part (Morioka et
2862 al. 2010, 2012). With these changes in the upper ocean heat capacity, warming of the surface mixed layer by the
2863 climatological shortwave radiation is enhanced in the southwestern part and becomes less effective in the northeastern
2864 part. As a result, the dipole SST anomalies appear in the southern Indian Ocean.

2865 Because the above mechanism operates more effectively as the thickness of the mixed layer becomes thinner, the return
2866 period of the SIOD is becoming shorter associated with the shoaling trend of the mixed layer (Yamagami and Tozuka,
2867 2015). Whether this mechanism is associated with decadal-to-interdecadal variations and/or global warming awaits further
2868 study. Many coupled models are relatively successful in simulating the SIOD with some biases in the location and structure

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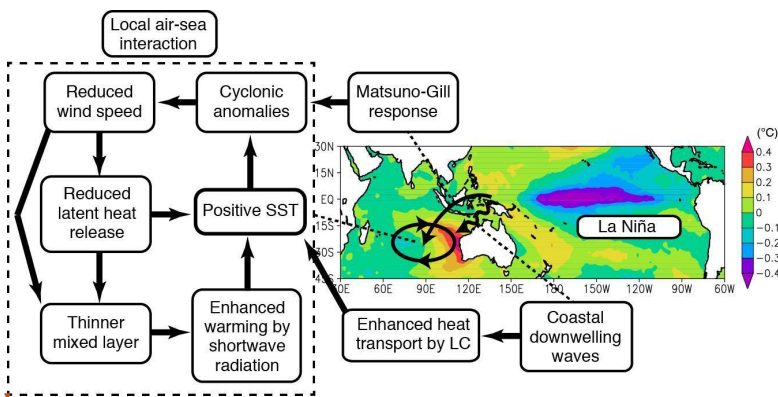
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2870 of the SST anomaly (Kataoka et al., 2012). However, no study has examined if the SIOD is modulated by climate modes
2871 of variability with decadal-to-interdecadal timescales or changes with global ocean warming.

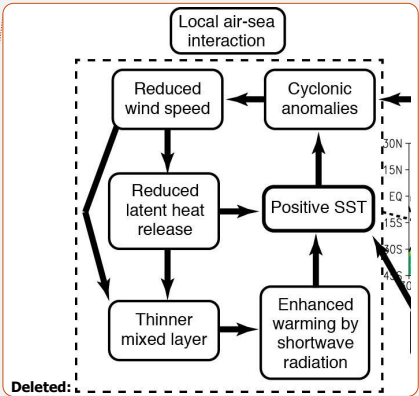
2872 **6.4 Ningaloo Niño and marine heatwaves in the Indian Ocean**

2873 The Ningaloo Niño (Niña) phenomenon is an interannual climate mode associated with anomalously warm (cold) water
2874 in the eastern Indian Ocean (Feng et al., 2013; see Figure 18). This mode is seasonally phase-locked, with a peak during
2875 austral summer (Kataoka et al., 2014). The mode exerts significant impacts on rainfall over Australia (Kataoka et al., 2014)
2876 and affects marine ecosystems and fisheries (e.g. Pearce et al. 2011). The phenomenon can alter biological productivity,
2877 with negative chlorophyll anomalies during Ningaloo Niño (Narayanasetti et al., 2016). Ningaloo Niños can develop in
2878 response to remote ENSO forcing from the western Pacific transmitted as a coastally trapped wave (Kataoka et al., 2014).
2879 During the La Niña events, high sea level anomalies propagate poleward along the west coast of Australia, intensifying
2880 the Leeuwin Current and causing poleward advection of heat and anomalously warm waters (e.g. Benthuisen et al., 2014;
2881 Section 4.2.4). Poleward transport of tropical, low salinity waters can further enhance the total geostrophic transport of
2882 the Leeuwin Current (Feng et al., 2015a).



2884 **Figure 18: Schematic diagram illustrating generation mechanisms (i.e. local air-sea interaction, atmospheric**
2885 **teleconnection, and oceanic wave propagation) of the Ningaloo Niño. SST anomalies are regressed against the**
2886 **Ningaloo Niño Index to illustrate typical SST anomalies associated with the phenomenon.**

2887 Atmospheric teleconnection can further enhance the development of Ningaloo Niño. A reduction in southerly winds over
2888 the shelf, which would strengthen the Leeuwin Current, can arise through a Gill-type response with low sea level pressure



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anomalies in the southeast Indian Ocean owing to the Niño3.4 SST anomalies (Feng et al., 2013; Tozuka et al., 2014). Ningaloo Niños can arise from local air-sea interactions off western Australia, through the wind-evaporation-SST feedback during its initial stage (Marshall et al., 2015) and coastal SST-wind-Leeuwin Current (Bjerknes) feedback (Kataoka et al. 2014). In the coastal feedback mechanism, positive SST anomalies lead to northerly alongshore wind anomalies and coastal downwelling anomalies, causing enhancement of the positive SST anomalies (Kataoka et al., 2014). During the Ningaloo Nino's development phase, estimates of air-sea heat flux contributions have been found to be dependent on products and their resolution and bulk flux algorithms (Feng and Shinoda, 2019). Since the late 1990s, Ningaloo Niño events have occurred more frequently (Feng et al., 2015b). This decadal increase is corroborated by coral proxy records of Leeuwin Current strength, with the most extreme SST anomalies associated with Ningaloo Niños occurring since 1980 (Zinke et al., 2014).

More generally, marine heatwaves refer to prolonged, extremely warm water events. Over the past decade, most studies on marine heatwaves in the Indian Ocean have focused on the eastern sector of the Indian Ocean. Major events in the Indian Ocean have been associated with phases of ENSO. Along the west coast of Australia, marine heatwaves have occurred predominantly at subtropical reefs during La Niña events due to increased heat transport (Zhang et al., 2017). The term "marine heatwave" was first coined owing to a +5°C warm water event in 2011 off Western Australia during a strong La Niña (Pearce et al., 2011). The 2011 event was associated with the strongest recorded Leeuwin Current transport anomaly, bringing warm tropical waters south, and was partly due to air-sea heat fluxes (Feng et al., 2013; Benthuyssen et al., 2014).

Across Australia's northwestern shelf, marine heatwaves have been found to occur at tropical coral reefs from El Niño due to solar radiation and a weakened monsoon (Zhang et al., 2017). During the strong El Niño of 2015-2016, the southeast tropical Indian Ocean experienced the warmest and longest marine heatwave on record, with weakened monsoon activity and anomalously high air-sea heat flux into the ocean (Benthuyssen et al., 2018). The anomalously warm water conditions persisted into winter, during one of the strongest negative IOD events (Benthuyssen et al., 2018). The 2016 marine heatwave was associated with coral bleaching spanning Australia's inshore Kimberley region to remote coral reef atolls (Gilmour et al., 2019). More broadly across the Indian Ocean during 2016, marine heatwaves have been studied in terms of their ecological impacts, such as coral bleaching in the western Indian Ocean (e.g. Gudka et al., 2018), the Maldives (e.g. Ibrahim et al., 2017) and consequences for fishes in the Chagos Archipelago (Taylor et al., 2019).

Trends in marine heatwave metrics indicate widespread regions across the Indian Ocean where events have increased in frequency, based on SST from 1982-2016, especially in the central and southwestern sectors (Oliver et al., 2018). Over the same time period, the duration and intensity of marine heatwaves have increased [in the Indian Ocean and globally](#) (Oliver et al. 2018, [Marin et al. 2021](#)). Primary climate modes of variability correlated with an increased occurrence of marine heatwaves include the following: (1) the positive phase of the Dipole Mode Index for the northwestern sector, the

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2922 tropical sector, and south to the Seychelles Islands, (2) the positive phase of the Niño3.4 index for the south-central sector,
2923 and (3) the negative phase of the El Niño Modoki index, which measures the strength of the Central Pacific ENSO, for the
2924 eastern Indian Ocean (Holbrook et al., 2019). While the marine heatwaves in the eastern Indian Ocean have been well
2925 documented, there have been fewer studies into the physical mechanisms causing marine heatwaves across the basin and
2926 other regions and less confidence, for example in the Bay of Bengal, in the local processes causing reported events on a
2927 range of time scales (Holbrook et al. 2019). There are indications that increased extremes in El Niño (Cai et al., 2014b)
2928 and La Niña events (Cai et al., 2015) due to mean ocean warming trends increase the likelihood of marine heatwave
2929 occurrence in the southeast Indian Ocean (Zhang et al., 2017).

2930 6.5 Monsoon variability and links to the Indian Ocean

2931 Several monsoon systems surround the Indian Ocean, notably the South Asian monsoon, the East Asian monsoon and
2932 the Australian monsoon. These monsoon systems are remotely influenced by global coupled modes of variability such as
2933 ENSO, which is often associated with dry conditions in the South Asian monsoon (e.g., Rasmusson and Carpenter,
2934 1983; Ropelewski and Halpert, 1987) and Australian monsoon (e.g., Risbey et al., 2009; Jourdain et al., 2013), although
2935 the relationship with the Indian monsoon has recently weakened (e.g., Kumar et al., 1999). In the Indian Ocean, the IOD
2936 has a strong influence on the Asian monsoon systems, but is weak during the Australian monsoon period. The IOD tends
2937 to oppose the ENSO teleconnection to the South Asian monsoon by enhancing monsoon rainfall (e.g., Ashok et al.,
2938 2004; Chowdary et al., 2015; Krishnaswamy et al., 2015; Pokhrel et al., 2012). However, the exact combination of SST
2939 patterns between the Indian Ocean and the Pacific is crucial for determining the rainfall response in the Asian monsoons
2940 (e.g., Lau and Wu, 2001; Ratna et al., 2020; Yuan and Yang, 2012), and the relative strengths of the teleconnections
2941 have varied over time (Krishnaswamy et al., 2015). Furthermore, there is evidence that the Indian Ocean forcing of the
2942 South Asian monsoon may be primarily driven by ENSO, with pure IOD events only weakly influencing monsoon
2943 rainfall (Cretat et al., 2017).

2944
2945 The monsoon systems around the Indian Ocean tend to vary in phase and are also linked to the western North Pacific
2946 Monsoon (e.g., Gu et al., 2010). There is a biennial oscillation in the strength of the monsoon systems, with a strong
2947 Asian monsoon preceding a negative IOD and coinciding with cold eastern Pacific SSTs, followed by a strong
2948 Australian monsoon and subsequently by a reversal in the SST patterns (Loschnigg et al., 2003; Meehl & Arblaster,
2949 2011). Thus, each monsoon system interacts with the ocean dynamics and thermodynamics and with the other monsoon
2950 systems through a complex set of teleconnections.

2951
2952 At a regional scale, upwelling in the Arabian Sea reduces rainfall along the western Ghats of India during the monsoon
2953 due to a reduction in evaporation and water vapour transport (Izumo et al., 2008). Moisture fluxes across the Arabian

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Sea are crucial to accurate simulation of the Indian Monsoon, yet many models fail to accurately capture these (Levine and Turner, 2012). In the Bay of Bengal, the shallow surface mixed layer, supported by the vertical salinity gradient, leads to rapid variations in SST (e.g., Sengupta and Ravichandran, 2001; Vecchi and Harrison, 2002) that interact with intraseasonal oscillations (Gao et al., 2019) in the atmosphere and thus with the active/break cycles on the monsoon (e.g., Lucas et al., 2014). This strong and rapid variability in upper ocean conditions in the Bay of Bengal, and the potential feedbacks on the monsoon, motivated multiple observational research programmes with field campaigns in the Bay of Bengal, as discussed in the next section.

7. Multiscale upper ocean processes in the Bay of Bengal

Reflective of its name, the Bay of Bengal is in many ways analogous to a large-scale estuary with seasonally reversing winds and boundary currents that facilitate the transport, stirring, and mixing of water masses. To the north, the Ganga-Brahmaputra-Meghna watershed delivers on average 1300 km³ in annual runoff of freshwater with a seasonal peak in discharge from July to September (Sengupta et al., 2006). During the southwest monsoon (boreal summer), the Summer Monsoon Current (Fig. 10) flows eastward advecting high salinity waters from the Arabian Sea into the southern Bay of Bengal, balancing the Bay's net outflow of freshwater. Instabilities and eddies result in mesoscale stirring of these different water types and create a strongly filamented and complex near-surface thermohaline structure. Lateral and vertical gradients in stratification are further modified by submesoscale processes, instabilities, and mixing. The resultant shallow stratification allows for rapid coupling with the atmosphere. Collectively, these conditions present a natural laboratory to study multi-scale mixing processes and their link to air-sea interaction. This section discusses new understanding of physical processes in the Bay from the large-scale to sub-mesoscale and finally at the smallest mixing scales.

Recent focus on the Bay of Bengal's upper ocean structure has been prompted by the need to understand atmosphere and ocean coupling with the aim of ultimately informing monsoon forecasting efforts at the intraseasonal timescale and shorter. Several bi-lateral international collaborations (Lucas et al., 2014; Wijesekera et al., 2016; Mahadevan et al., 2016; Vinayachandran et al., 2018; Gordon et al., 2019, 2020) have collectively supported multiple field campaigns, beginning in 2013 and concluding in 2019, using a combination of shipboard, moored, and autonomous platforms. These atmospheric and oceanic measurements have provided new insights into the BoB's structure and the processes that regulate that structure, particularly at fine lateral scales (<5 km).

Results from these combined efforts span from large-scales, e.g., the quantification of coastal transport along the Sri Lankan coast (Lee et al., 2016) and the mesoscale stirring of freshwater (Sree Lekha et al., 2018), to intermediate scales, e.g., high-resolution (order 100 m) frontal surveys that hint at the roles of submesoscale (Ramachandran et al., 2018) and non-hydrostatic processes in setting stratification (Sarkar et al., 2016), to small-scales with direct measurements of

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2985 microstructure yielding new insights into the BoB's mixing regimes (Jinadasa et al., 2016; Thakur et al., 2019; Cherian et
2986 al. 2020).

2987 7.1 The Bay's Forcing and Upper Ocean Structure

2988 At the largest scales, the Bay is forced by air-sea fluxes of buoyancy and momentum, which are strongly modulated by
2989 the monsoon and vice versa. Precipitation and multiple river systems, including the Ganga-Brahmaputra-Meghna system,
2990 contribute to freshwater input that creates a barrier layer in the surface Bay of Bengal, which is strongest in the northern
2991 Bay weakening toward the south. The Bay's stratification, in particular its barrier layer, is unique in how it impacts the
2992 evolution of seasonal SST, in turn setting the lower boundary condition for the development of the monsoon (Li et al.,
2993 2017). For this reason, recent emphasis has been placed on understanding processes that determine the Bay's upper ocean
2994 salinity and temperature structure.

2995 The monsoon cycle of surface forcing plays a first-order role in controlling the Bay's upper ocean temperature structure.
2996 Direct flux measurements are a critical component in our ability to accurately capture/represent and predict the magnitude
2997 and variability of monsoon air-sea coupling. Recent studies have shown that of the air-sea heat flux terms, shortwave
2998 radiation and latent heat flux are the largest drivers of variability to the total heat tendency. These variables are also those
2999 which reanalysis products struggle most to accurately represent, showing biases up to 75 W/m² (Sanchez-Franks et al.,
3000 2018). High-quality air-sea surface flux measurements over the BoB historically have been limited to the few sites
3001 maintained by the RAMA array (McPhaden et al., 2009). However, regional measurement efforts have expanded and
3002 baseline surface measurements are now collected and sustained through India's National Institute of Ocean Technology's
3003 met-ocean buoy program (Venkatesan et al., 2018), as well as the recent transition of an 18°N air-sea flux buoy from
3004 Woods Hole Oceanographic Institution to Indian National Centre for Ocean Information Services (Weller et al., 2016).

3005 Precipitation and riverine discharge along the Bay's margins respectively contribute roughly 60% and 40% of the 0.14 Sv
3006 net freshwater delivered to the Bay (Sengupta et al. 2006; Wilson and Riser, 2016). Precipitation peaks in early summer
3007 (June) with a value near 0.4 m month⁻¹, while discharge peaks slightly later in summer (August) with a value near 0.3 m
3008 month⁻¹. Evaporative loss (included in the net 0.14 Sv) is relatively steady throughout the year at 0.1 m month⁻¹ (Wilson
3009 and Riser, 2016). Estimates of river discharge from gauged sources are known to have uncertainties (underestimates)
3010 related to unmonitored tributaries and streams. For large deltas, altimeter-based elevations offer a means of extrapolating
3011 gauge data over space and time. Papa et al. (2010, 2012) applied such an approach to the Ganga-Brahmaputra River
3012 system for the period 1998-2011. This time series allows for assessment of interannual variability over time ranges not
3013 spanned by gauged efforts. Papa et al. (2012) note a 12,500 m³/s standard deviation in interannual variability in the Ganga-
3014 Brahmaputra discharge. Importantly, such data sets are also easily accessible by the general public, facilitating progress
3015 and understanding by the scientific community.

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3020 The Bay's upper ocean temperature and salinity structure is an integrated representation of the above summarized
3021 sources/sinks of heat and freshwater, combined with the physical processes that redistribute these quantities. The
3022 thermohaline structure of the Bay is remarkable in several regards—for shallow mixed layer depths (< 5 m, Sengupta and
3023 Ravichandran, 1998), for inversions of temperature (Shroyer et al., 2016, 2019; Thadathil et al. 2016), for large-scale
3024 coherent layering that spans 100 kms (Shroyer et al., 2019), an active mesoscale field and the strong influence of river
3025 discharge over the interior basin. The Bay's salinity stratification is a critical, if not dominant, contributor to the upper
3026 ocean density stratification. It supports the formation of barrier layers that are frequently observed to be warmer than the
3027 mixed layer thereby providing a substantial subsurface heat reservoir with the potential to modify air-sea interaction
3028 (Girishkumar et al., 2011; Shroyer et al., 2016). For example, in conditions supportive of formation of a diurnal warm
3029 layer (low winds, strong insolation), subsurface turbulent fluxes can act to modulate the diurnal SST cycle by transporting
3030 (typically) warm barrier layer waters into the mixed layer at night while still cooling the base of the diurnal warm layer
3031 (DWL) during the day (Shroyer et al., 2016). A similar phenomenon, albeit on a much different scale, results with passage
3032 of cyclones, which often show a salty wake even in the absence of a cool wake which is common for cyclones elsewhere
3033 (for e.g. Chaudhuri et al. 2019, Qiu et al. 2019). Below, we review recent progress on understanding of processes that
3034 determine the Bay's upper ocean thermohaline structure.

3035 7.2 Lateral Processes

3036 7.2.1 Stirring from the Margins

3037 The Bay of Bengal has an active mesoscale eddy field that stirs diverse source waters into the interior of the Bay of Bengal.
3038 The origins of these source waters are the Arabian Sea waters to the west, the Ganga-Brahmaputra-Meghna at the northern
3039 tip, Andaman Sea waters to the east, and Equatorial waters to the south. This stirring effectively contributes to a quasi-
3040 stationary balance of the fresher waters from the north and the high salinity waters from the west and south over time.
3041 Lateral advection is a fundamental contributor to the formation of the barrier layer (George et al., 2019) and the freshwater
3042 budget of the Bay (e.g. Sree Lekha et al., 2018). In the northern Bay, the dispersal of water from the periphery into the
3043 interior depends critically on mesoscale stirring and the time varying Ekman transport, as indicated from mooring (Sree
3044 Lekha et al., 2018) and ship-based surveys (Shroyer et al., 2019), and constrained by modelling results (Sree Lekha et al.,
3045 2018). Here, the advection of freshwater by the mesoscale stirring also plays an important role in determining SST over
3046 the northern BoB (Buckley et al. 2020), as these waters are typically associated with relatively shallow mixed layers. In
3047 the southern Bay, measurements have suggested the competing influences of mixing and advection of salty Arabian Sea
3048 water in the erosion and reformation of the barrier layer during the southwest monsoon (George et al., 2019;
3049 Vinayachandran et al., 2018). In particular, George et al. (2019) show that maintenance of the barrier layer and the
3050 associated maximum depth of mixing was critically dependent on horizontal advection through its impact on stratification.

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Surface freshwater input also has an impact on barrier layer evolution; several freshening events were captured at various stages of their seasonal evolution in the southern Bay of Bengal in recent observations (Vinayachandran et al., 2018). These events play a significant role in the formation of a thick barrier layer, showing that during the southwest monsoon the shoaling of the mixed layer in the southern BoB has a similar magnitude and behaviour to that in the northern BoB (Vinayachandran et al., 2018).

7.2.2 Inter-basin exchange

Inter-basin exchange is critical to the Bay's salinity budget; since the Bay receives net freshwater input, this freshwater must be balanced by salty water imported from either the Arabian Sea or the western equatorial Indian Ocean (Jensen et al., 2001; Sanchez-Franks et al., 2019), and turbulent transport of salt into the fresh water layer is necessary to maintain the BoB's long-term salinity balance. Observations show that intrusion of high salinity water from the Arabian Sea enters the BoB between 80°-90°E during the southwest monsoon, (e.g. Murty et al, 1992; Vinayachandran et al., 2013) and has been found in several models (e.g. Vinayachandran et al., 1999; Han and McCreary, 2001 and Jensen, 2001). More recent observational and modeling studies show that both lateral and vertical transfer of heat and salt occur at multiple space-time scales. Seasonal currents play an important role in transporting heat and salt in and out of the BoB, but the role of mesoscale eddies on lateral transports is not well known.

Using unique year-long mixing measurements detailed in Section 7.3, Cherian et al. (2020) tentatively estimated a turbulent salt flux of $1.5\text{e-}6 \text{ psu ms}^{-1}$ out of Arabian Sea water averaged between 85°E and 88.5°E at 8°N through the 34.75 psu isohaline between August and January. Over those 6 months, this flux would increase the salinity of a 75m layer of water by 0.3 psu, though much of this would be cancelled out by surface fluxes. The magnitude and timing of this salt flux roughly match that necessary to restore the Bay's near-surface salinity after the large freshwater input in August as estimated by a few modelling studies (Akhil et al., 2014; Benshila et al., 2014; Wilson and Riser, 2016). This is the first direct measurement of turbulence that supports the hypothesis of intrusion of high salinity water from the Arabian Sea during the southwest monsoon (Vinayachandran et al., 2013).

7.2.2.1 Andaman Sea Exchange

The Irrawady river drains into the Andaman Sea, a marginal sea at the eastern edge of the Bay. Export from the Andaman is then another source of freshwater for the Bay, particularly at intermediate densities ($22\text{-}25 \text{ kg m}^{-3}$). A striking example of the interaction between strong surface forcing and an anticyclonic eddy can be found in the fortuitous crossing of an intrathermocline eddy (ITE) in 2013, as reported by Gordon et al. (2017). The water mass characteristics clearly identify ITE waters from the Andaman Sea; and, analysis of ancillary Argo data suggest a similar water type often penetrates westward into the Bay extending from the three passages connecting the two basins. While at the time of transit the observed ITE had a very weak surface expression, a week prior to encountering the ITE a clear sea surface high ($>10 \text{ cm}$)

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3085 is evident in AVISO SSHA. Tropical cyclone Lehar passed near the location of this sea surface high in the interim, and
3086 the working conjecture is that the winds associated with Lehar were sufficient to modify a typical mode-1 anticyclone into
3087 the observed ITE.

3088 7.2.2.2 Arabian Sea Exchange

3089 Near-surface exchange from the Arabian Sea into the Bay of Bengal is influenced by the Sri Lanka Dome (SLD), an
3090 upwelling thermal dome that recurs seasonally within the SMC in the wind shadow of Sri Lanka (Vinayachandran and
3091 Yamagata 1998, de Vos et al. 2014, Burns et al. 2017). The SLD has long been recognized as a prominent circulation
3092 feature in the southwestern bay during the summer monsoon; and it has been noted as a region of enhanced productivity
3093 (Vinayachandran et al., 2004, de Vos et al. 2014), cool SST (Burns et al. 2017), and consequently depressed convection
3094 (Figure 15). The SLD displays pronounced interannual variability (Cullen and Shroyer 2019). In some years the SLD has
3095 a strong surface manifestation (amplitude of the low ~30 cm) that persists well beyond the southwest monsoon; in other
3096 years the SLD has a weak expression that is intermittent and short-lived (~1-2 months). The SLD is not fixed in location
3097 despite its strong association with the wind stress curl. Its position varies from year-to-year as well as over the course of
3098 one season. Variations in its location and strength may influence the properties of waters entrained and upwelled within
3099 the SLD.

3100 At intermediate depths (<~200 m), the signature of the neighboring Arabian Sea is notable across much of the basin
3101 (Gordon et al., 2016). During summer, Arabian Sea High Salinity Water (ASHSW; density near 22-24 kg m⁻³) is
3102 carried/advectioned into the Bay of Bengal as a 'high salinity core' via the Southwest Monsoon Current (SMC, Webber et
3103 al., 2018; Sanchez-Franks et al., 2019) and then spread north along the bay's central spine (Hormann et al., 2019). During
3104 this journey, salt is mixed upward into the near-surface fresh layer (Cherian et al 2020; Section 7.3). A nearly two-year
3105 long moored current record in the southern BoB captured seasonally varying large eddies generated by the SMC and
3106 Northeast Monsoon Current. These eddies included a cyclonic eddy, the SLD, and an anticyclonic eddy south of the SLD
3107 (Wijesekera et al. 2016c). These observations revealed that the average transport over a nearly two year period into the
3108 BoB was about 2 Sv (1 Sv = 10⁶ m³ s⁻¹) but likely exceeded 15 Sv during summer of 2014, which is consistent with the
3109 transport associated with the SMC (e.g., Schott et al. 2009; Webber et al. 2018). The observations further indicate that the
3110 water exchange away from coastal boundaries, in the interior of the BoB, may be largely influenced by the location and
3111 strength of the two eddies that modify the path of the SMC. The strength and location of the SMC itself is dependent on
3112 a combination of local and remote forcing (Webber et al., 2018).

3113 As discussed above several hypotheses have been suggested for cyclonic eddy (SLD) and anticyclonic eddy formation in
3114 the southern BoB. It has been suggested that the cyclonic wind stress-curl over southwestern BoB generates the SLD
3115 (McCreary et al., 1996; Vinayachandran and Yamagata 1998; Schott et al., 2001; Cullen and Shroyer 2019). Based on

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numerical simulations, de Vos et al., (2014) argued that the separation of SMC from the (southern) boundary of Sri Lanka may lead to SLD, where a cyclonic vorticity is generated by lateral frictional effects. A mechanism for the anticyclonic eddy formation has been proposed by Vinayachandran and Yamagata (1998), where the interaction of the SMC with Rossby waves arriving from the eastern boundary leads to the anticyclonic eddy. Pirro et al (2020a) proposed a new hypothesis wherein the anticyclonic eddy is generated by a topographically trapped Rossby wave response of the SMC to perturbations by the Sri Lankan coast. They reported that observations of the size, location and origins of the SLD were broadly consistent with their hypothesis, based on a laboratory experiment designed to mimic natural flow in the BoB by creating an eastward jet (SMC) on a simulated β plane.

High-resolution sampling of the interior BoB has provided a more detailed look at the lateral extent of typical ‘patches’ of Arabian Sea water, which tend to remain well-defined over scales of 10-50 km, suggesting the importance of eddy activity in exchange (Shroyer et al., 2019). While many studies have traced origins of ASHSW from the eastern Arabian Sea, entering the Bay of Bengal directly via the southwest monsoon current (e.g., Jensen et al, 2016); a recent study suggests an equatorial pathway may also be relevant (Sanchez-Franks et al., 2019; [Section 7.2.3](#)). Highly salty and highly oxygenated waters from the Persian Gulf and the Red Sea have also been noted in the southern regions of the Bay of Bengal (Jain et al., 2017). These waters are injected into the Bay of Bengal via current systems (equatorial and the southwest monsoon current) with important repercussions for the oxygen concentrations of the Bay of Bengal oxygen minimum zone (Sheehan et al., 2020).

Velocity and hydrographic profiles from a shipboard survey in December 2013 combined with drifter observations, satellite altimetry, global ocean nowcast/forecast products, and coupled model simulations were used to examine the circulation in the southern Bay of Bengal during the Northeast monsoon (Wijesekera et al. 2015). The observations captured the southward flowing East India Coastal Current (EICC, e.g., Shetye et al. 1994) off southeast India and east of Sri Lanka. The EICC was approximately 100 km wide, with speeds exceeding 1 m s^{-1} in the upper 75 m. East of the EICC, a subsurface-intensified 300-km-wide, northward current was observed, with maximum speeds as high as 1 m s^{-1} between 50 m and 75 m. The EICC transported low-salinity water out of the bay and the subsurface northward flow carried high-salinity water into the bay during typical northeast monsoon conditions (Wijesekera et al. 2015; Jensen et al. 2016).

7.2.3 Equatorial Connections

The Equatorial undercurrent (EUC) in the Indian ocean is seasonally variable. The summer–fall EUC tends to occur in the western basin in most years but exhibits evident interannual variability in the eastern basin (Chen et al. 2015), with different processes dominating its generation in the western and eastern basins. In the eastern basin reflected Rossby waves from the eastern boundary play a crucial role in the EUC, whereas directly forced Kelvin and Rossby waves control the EUC in the western basin.

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Equatorial Kelvin waves, commonly interpreted as Wyrtki (1973) jets, propagate eastward along the equator during April/May and September/October. Upon reflection from the IO eastern boundaries, energy of Wyrtki jets is reflected back in part as long Rossby waves that disperse slowly during the following two months and reach the central-eastern BoB during July-August (Han et al., 1999, 2001; Han, 2005; Nagura and McPhaden, 2010a). The remaining energy is partitioned into two coastally-trapped Kelvin waves traveling poleward (Moore, 1968), which excite long Rossby waves propagating westward. Therefore it is suggested that planetary waves driven by remote forcing from the interior IO contribute significantly to the formation, strength and intensity of the BoB circulation (Vinayachandran et al. 1998; Nagura and McPhaden, 2010b; Chen, 2015). A subset of these planetary waves are the mainstay of intraseasonal oscillations (ISOs), a sub-seasonal phenomenon of period less than 120 days. The genesis of oceanic ISOs has been attributed to multiple mechanisms: external forcing (e.g., atmospheric ISOs and Ekman pumping, e.g. Duncan and Han 2012) and internal processes (upper ocean processes and instabilities e.g. Zhang et al. 2018).

Observations in the IO have captured a range of variabilities in the 30 – 120 days frequency band (e.g., Girishkumar et al., 2013), and past research has identified roughly three distinct ISO bands in the context of the thermocline: 30-60 days, 60-90 days, and 120 days (Han et al., 2001; Girishkumar et al., 2013). Pirro et al. (2020b) discussed interaction between 30-60 day ISOs and the SMC in the southern BoB using long-term moored observations. They estimated that the background mean flow acceleration resulting from the meridional divergence of wave momentum flux in the thermocline was about 10^{-8} m s^{-2} . As a result, within a wave period, ISOs can enhance the eastward flow in the thermocline by about 25%. The negative shear production computed for the same period is consistent with this finding suggesting that the mean flow gained kinetic energy at the expense of the ISO band. The meridional heat-flux divergence was $-10^{-7} \text{ }^{\circ}\text{C s}^{-1}$ and has a tendency for cooling the thermocline by about 0.5°C when ISOs are active (Pirro et al., 2020b). Observations have also captured energetic and consequential 5-20 day convectively coupled Kelvin waves in the atmosphere (Baranowski et al. 2016) that generate oceanic Kelvin waves, affect surface heat fluxes and generate upper ocean turbulence (Pujiana and McPhaden, 2018).

High salinity waters from the western Arabian Sea and the western Equatorial Indian Ocean can route to the Bay of Bengal via the Somali Current and the Indian Ocean EUC (Sanchez-Franks et al., 2019). Changes in strength of the Bay of Bengal high salinity core are linked to the convergence of the East Africa Coastal Current and the wintertime southward-flowing Somali Current, with anomalously strong equatorial Undercurrent (Fig. 19). Because of the seasonal reversal of currents, two junctions form naturally, one in the western equatorial Indian Ocean (Somali Current) and another south of India (monsoon currents), which effectively act as ‘railroad switches’ rerouting water masses to different basins in the Indian Ocean depending on the season (Fig. 19, Sanchez-Franks et al., 2019).

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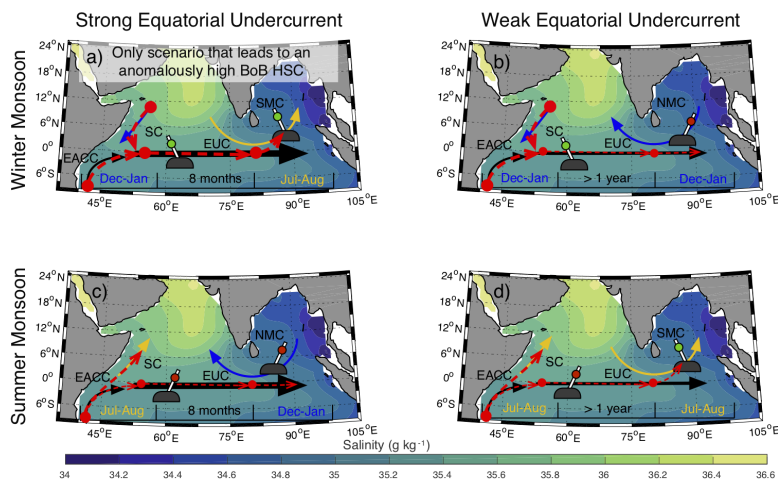
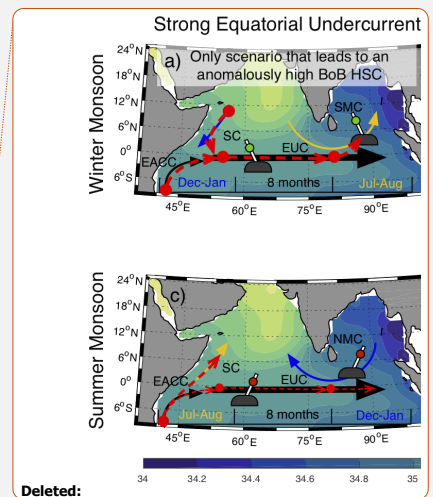


Figure 19: Seasonal circulation pathways in the northern Indian Ocean, or Railroad Switch schematic, on subsurface (90 m) salinity climatology (psu; shaded) from the Argo optimally interpolated product for the four Equatorial Undercurrent scenarios: (a, b) winter monsoon and strong (weak) Equatorial Undercurrent and (c, d) summer monsoon and strong (weak) Equatorial Undercurrent. Red dashed arrows indicate high-salinity advection. BoB = Bay of Bengal; HSC = high-salinity core; SMC = Southwest Monsoon Current; SC = Somali Current; EUC = Equatorial Undercurrent; EACC = East African Coastal Current. From Sanchez-Franks et al. (2019).

7.3 Vertical Mixing

Strong stratification in the Bay of Bengal plays a critical role in setting the upper ocean turbulence, notably leading to relatively weak mixing compared to other regions (e.g. Gregg et al., 2006). However, large-scale inferences suggest that mixing must play a key role in at least two regards. First, the net surface flux during the southwest monsoon on average is warming but yet the SST cools (Shenoi et al, 2002). Second, the large-scale salt balance must be closed through upward mixing of high-salinity water carried into the Bay via the Summer Monsoon Current (Vinayachandran et al., 2013).

Recent year-long direct measurements of mixing in the Bay have helped link the seasonal cycle in mixing to the seasonal cycle of winds, currents and freshwater. These year-long measurements were recorded by mixing meters called χ pods. χ pods consist of two temperature microstructure sensors and a suite of ancillary sensors necessary to infer the rate of



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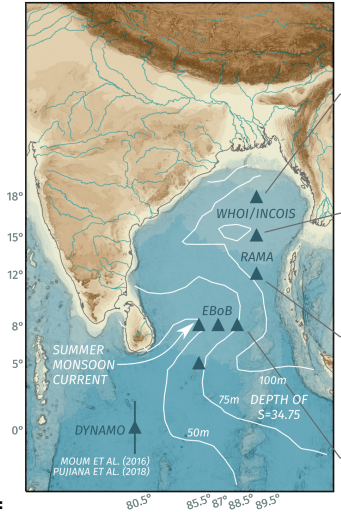
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3207 dissipation of temperature variance at 1Hz frequency for up to a year (Moum & Nash, 2009). χ pods have been deployed
3208 on moorings in three different regions of the Bay (Figure 20): the air-sea buoy at 18°N, top 65m (Thakur et al., 2019),
3209 RAMA moorings along 90°E (mixing measurements at 15m, 30m and 45m; Warner et al. 2016), and the EBoB array in
3210 the south-central Bay (mixing measurements spanning between 30m and 100m at sites in the region 85°E-88°E, 5°N-8°N,
3211 Cherian et al., 2020). Across the basin, turbulence within and near the base of the mixed layer shows strong seasonality
3212 that parallels the monsoon cycle in winds (Thakur et al., 2019, Warner et al., 2016). In the thermocline of the south-central
3213 Bay (EBoB array), mixing is correlated with packets of downward propagating near-inertial waves implicating wind
3214 forcing. As depicted in Figure 20, both near-surface and thermocline mixing are relatively high during the NE and SW
3215 monsoons (Dec-Feb, May-Sep) and relatively low during the transition (Mar, Apr). Cyclones during the post-monsoon
3216 months of October and November can drive a hundredfold increase in near-surface mixing both locally and throughout
3217 the Bay (Warner et al. 2016). **Turbulence** profiles collected by a fast thermistor on a CTD rosette during a basin-wide
3218 survey before and after the passage of cyclone Madi (6-12 Dec, 2013) show a basin-wide increase in diffusivity linked to
3219 near-inertial waves forced by the cyclone (Wijesekera et al., 2016b).

3220 Indirect estimates of turbulent diffusivity and turbulent heat fluxes at the base of the mixed layer can be found as the
3221 residual of a mixed layer heat budget whose terms are estimated using a combination of mooring and satellite
3222 measurements. Girishkumar et al. (2020) use this approach to indirectly estimate seasonal median turbulent diffusivities
3223 using decade-long RAMA mooring records at 90°E. They find a robust seasonal cycle of mixing at 8°N, 12°N, and 15°N;
3224 and strong latitudinal variability in turbulence, with larger diffusivities inferred at 8°N relative to 12°N and 15°N in all
3225 seasons. When comparisons are possible, the indirect estimates compare well against the more direct but time-limited
3226 estimates of Warner et al (2016) at 90°E, 12°N.

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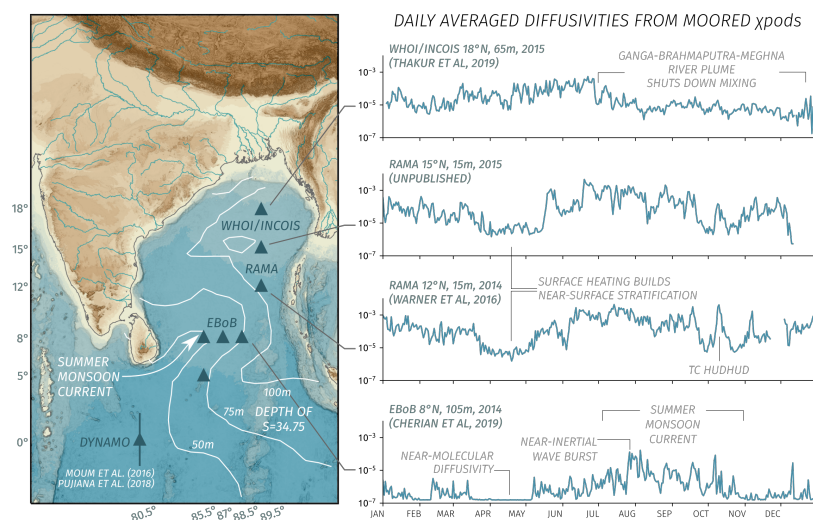


Figure 20: Annual cycle of daily averaged temperature diffusivities derived from xpod measurements. The data are from two different years, 2014 and 2015, depending on location. Note the similar wind-forced seasonal cycle at 12°N, 15m and 15°N, 15m and the dramatically different seasonal cycle at 8°N, 105m (reflecting near-inertial wave activity) and at 18°N, 65m reflecting freshwater influence.

The influence of freshwater is a critical caveat to the above generalizations: the arrival in August of the Ganga-Brahmaputra-Meghna freshwater plume at 18°N has been observed to suppress turbulence (diffusivity $K_T < 10^{-5} \text{ m}^2 \text{ s}^{-1}$) for multiple months (Aug-Nov) at depths of approximately 50-65 m (Figure 20). This buoyant lens limited the vertical extent of the influence of Tropical Cyclone Komen as compared to a previous (weaker) storm (Chaudhuri et al 2019, Thakur et al 2019). Similar observations of extremely weak turbulence below strong, salinity-stratified surface layers have been reported throughout the Bay using data from a variety of platforms: ship-based microstructure (Jinadasa et al, 2016) profiling floats with a temperature microstructure sensor (Shroyer et al, 2016) and glider-based microstructure measurements (St. Laurent and Merrifield, 2017). Lucas et al (2016) find that near-inertial shear was elevated at the base of the mixed layer but not elevated at the base of the barrier layer — direct evidence that salinity stratification can insulate deeper depths from the effects of near-surface forcing (downward propagating near-inertial waves in this case). Li et al. (2017) use a combination of observations and modelling results to demonstrate that barrier layers in the Bay of Bengal

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influence the amplitude of intraseasonal oscillations in SST and precipitation. However, a recent coarse resolution coupled modelling study suggests that freshwater has little influence on SST or rainfall, since the SST tendency caused by a reduction in mixing is offset by changes in surface heat fluxes (Krishnamohan et al., 2019)

Surface freshwater advection can create subsurface reservoirs of heat and salt that can be accessed when the winds are strong enough, such as during cyclones that regularly form in the Bay during October and November. In one dramatic example Qiu et al (2019) report up to 5 psu increases in SSS and only a smaller 0.5°C decrease in SST following the passage of Cyclone Phailin (2013). In this case, mooring records indicate that mixing was limited to the isothermal layer (Chaudhuri et al. 2019). Subsurface warm layers (i.e. temperature inversions stabilized by strong salinity stratification) are also observed, representing a reservoir of heat that can be accessed if a storm excites enough turbulence, as appears to have happened during the passage of Cyclone Hudhud (Warner et al, 2016). The influence of stratification in limiting the extent of vertical mixing and creating subsurface warm layers mean that cyclone-induced cooling is generally either weak or negligible in the Bay, unlike in other ocean basins (Sengupta et al, 2008). Subsurface warm layers influence SST on longer timescales too: Girishkumar et al (2013) find that the wintertime SST at 8°N, 90°E is quite sensitive to the thickness of the barrier layer, and to the presence of temperature inversions (subsurface warm layers) in the barrier layer on intraseasonal and interannual timescales.

Long periods of near-molecular diffusivities (weeks to a month) were also inferred at multiple χ pods along 8°N between 50 m and 100 m during transition months of March and April. Here freshwater insulation does not appear to be the major factor. Instead the period of weak turbulence may be linked to low levels of near-inertial energy (a consequence of weak wind forcing in March and April) and the absence of strong mean oceanic flows during these transition months (Cherian et al 2020). Relatively weak diffusivities are also present in the LADCP fine structure estimate of depth-integrated (thermocline to bottom) turbulent kinetic energy dissipation ϵ (Kunze et al, 2006) and the Argo fine structure-based 250-500 m diffusivity estimates of Whalen et al. (2012). The extended presence of such weak turbulence suggests that the Bay's internal wave field is weaker than might be expected from the Garrett-Munk internal wave spectrum at least during some months of the year. Another (related) question is the issue of representation of such weak background mixing in climate models and whether that matters to known biases in such models.

Published efforts so far have been directed towards understanding the modulation of turbulence by larger-scale variations in the wind, currents and freshwater. Questions remain as to the impact of small-scale mixing on the large-scale long-term T-S structure in the Bay as well as the influence of subsurface mixing and the ensuing modification of SST on coupled ocean-atmosphere phenomena such as the MJO and the MISO (Section 3.2)

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7.4 Where vertical and lateral processes meet: The Role of Submesoscale

Freshwater inflow from the Ganga-Brahmaputra-Meghna (GBM) and the Irrawaddy river in the Bay of Bengal is stirred by the mesoscale eddies into sharp frontal gradients (in salinity and in density) at $O(1-10\text{km})$ scales with shallow vertical extent. These fronts are acted upon by winds seasonally, setting up complex sub-mesoscale structures with salinity differences $O(1\text{psu})$ over $1-10\text{km}$, developing bore-like features with $O(0.5\text{psu})$ difference over a few meters horizontally (Nash et al 2016; Figure 21). Wavenumber spectra of temperature at $O(1-10\text{km})$ scale show a -2 slope in many regions of the Bay (Mackinnon et al 2016), a signature of frontogenesis in the Bay at these scales. The BoB is thus replete with fronts which evidently slump at sub-mesoscales due to both symmetric and baroclinic instabilities (Ramachandran et al. 2018), and show higher stratification near fronts (Sree Lekha 2019).

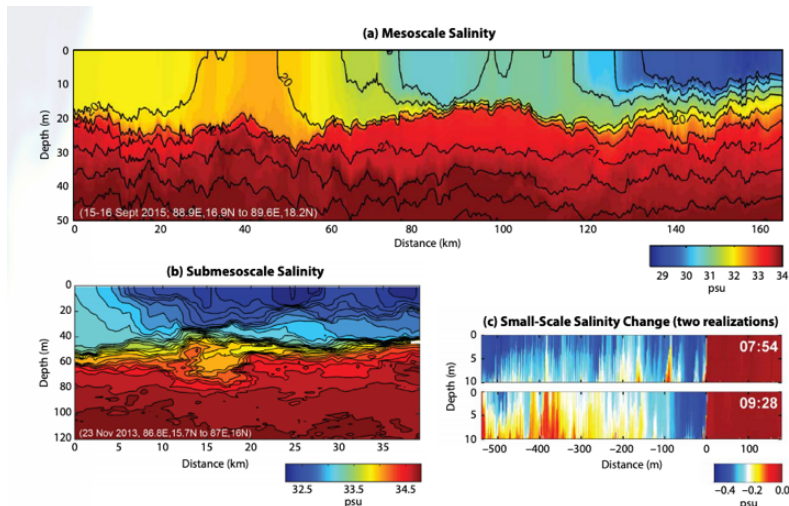


Figure 21: Observed salinity gradients at mesoscale, sub-mesoscales and small horizontal scales from in the Bay of Bengal (Nash et al. 2016).

The fronts and filaments at $O(1-10\text{km})$, which are dominated by salinity gradients and weakly compensated, have strong implications for setting up the density stratification in the top $50-100\text{m}$ in the BoB (Section 4.4.1). The stratification in this depth range often has multi-layered structure with stratification varying at $O(1-10\text{km})$ scales (Lucas et al 2016), showing evidence that the stratification in the Bay cannot be explained simply in terms of vertical processes, and horizontal

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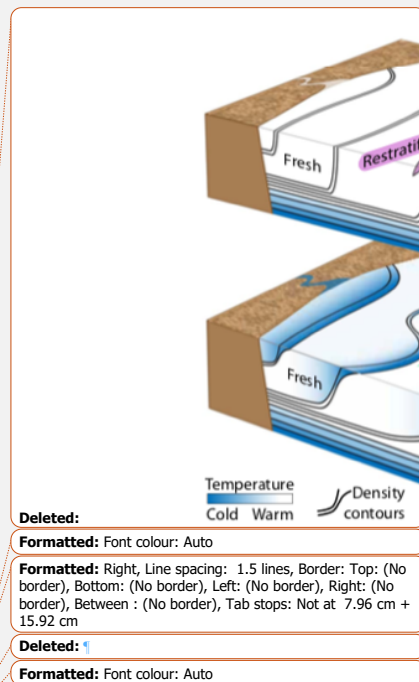
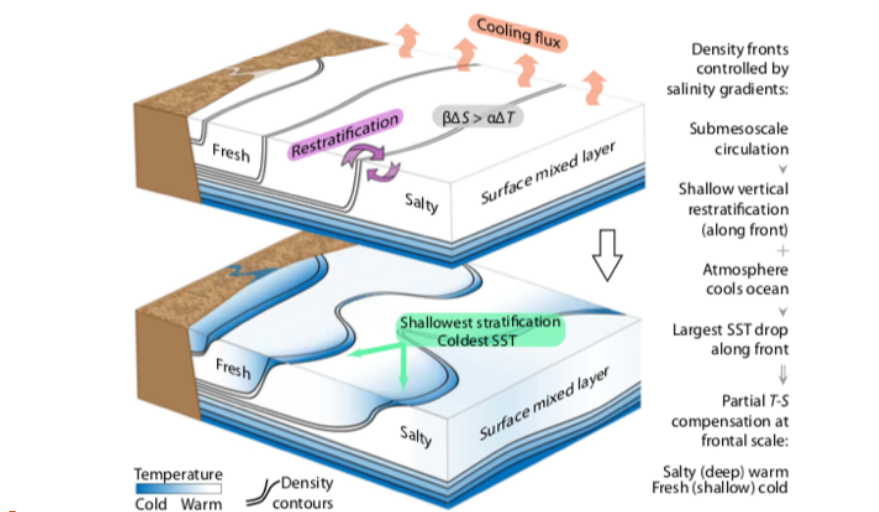
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submesoscale processes are intimately coupled with the vertical processes at these scales. Ramachandran et al. (2018) show that a mesoscale strained region with strong fronts ($O(1\text{kg/m}^3 \text{ over } 40\text{km})$) and weak down front wind shows multiple dynamical signatures of sub-mesoscale instabilities. Ageostrophic secondary circulations arising near the fronts and the accompanied sheared advection plays an important role in setting the stratification (Pham and Sarkar 2019). Both observations and process modeling show $O(1\text{-}10\text{km})$ patches of low potential vorticity consisting of subducted warm water patches due to a combination of baroclinic and forced symmetric instabilities, creating barrier layers whose thickness varies laterally at sub-mesoscales (Ramachandran and Tandon, 2020 JGR-in review).

During winter, the temperature gradients in the horizontal compensate for the salinity gradients to reduce the density gradient, and the sub-mesoscale processes in BoB lead to a unique situation. Jaeger & Mahadevan (2018) show that surface cooling fluxes combined with submesoscale instabilities of the haline fronts during wintertime leads to shallower mixed layers on the less saline (cooler) side. Therefore, cold SSTs in wintertime in the Bay mark surface trapped waters (Fig. 22), whereas in other regions of the world ocean, cold filaments mark upwelling of nutrient-rich waters. Further, since the shallow fresher mixed layers lead to larger drops in temperature, this develops the correlation between SST and SSS at $O(1\text{-}10\text{km})$ scales.



3309 **Figure 22: Interaction of submesoscale salinity gradients with atmospheric cooling leads to shallow cold regions**
3310 **(From Spiro Jaeger and Mahadevan, Science Advances 2018)**

3311 7.5 Putting the Pieces Together

3312 7.5.1 Coupled ocean-atmosphere phenomena

3313 Due to the presence of a barrier layer over much of the Bay of Bengal, entrainment and upwelling of waters from the
3314 thermocline are inhibited, and the evolution of SST is largely driven by net air-sea heat flux variability (Duncan and Han,
3315 2009). However, the dependency of SST on surface fluxes is controlled by subsurface processes such as formation of
3316 barrier layers, entrainment warming and cooling of the mixed layer, penetrative solar radiation and zonal advection
3317 (Thangaprakash et al., 2016). Advection is important in influencing the SST as lateral variations in the mixed layer depth
3318 alone can result in variations in air-sea fluxes of roughly 20 Wm^{-2} over distances of kilometers (Adams et al., 2019). This
3319 magnitude is similar to uncertainty in air-sea flux products (Weller et al. 2016) thus implying that variations in sub-
3320 mesoscales are important for heat balance in the northern BoB. The coupling of the ocean-atmosphere over BoB at large
3321 scales implicates the air-sea interaction and the mixed layer heat budget in [the](#) BoB (Rahaman et al. 2019), although at
3322 oceanic mesoscale and finer scales in the horizontal and at sub-seasonal timescales this coupling is a topic of active
3323 research.

3324 7.5.2 Implications for biogeochemistry in the Bay

3325 Eddies in the central BoB arise not by the baroclinic instability of boundary currents but rather due to planetary wave
3326 dynamics off the equator that triggers coastal Kelvin waves around the Bay. The Kelvin waves then trigger south-westward
3327 propagating Rossby waves, which result in large mesoscale structures in the Bay (Cheng et al. 2018). The Andaman and
3328 Nicobar Islands are also shown to be very important for the generation of these eddies; without these islands the number
3329 of eddies would have reduced to almost half in the western bay of Bengal (Mukherjee et al., 2019). These eddies provide
3330 much of the horizontal stretching and stirring of the tracers, including those relevant to the ecosystems

3331 Eddies have tremendous potential to influence ocean biogeochemistry by providing “new” nutrients to the ocean’s
3332 euphotic layer (Stramma et al., 2013). However, we do not fully understand the spatial distribution of nutrients within the
3333 eddy surface area – e.g., there is a debate whether nutrients upwell at the core and downwell at the edge of the eddy, or
3334 vice versa. Further, such discrepancy also continues in the type of eddies – i.e., whether upwelling occurs in cyclonic and
3335 downwelling occurs in anticyclonic eddies and vice versa (Mahadevan, 2014; Mahadevan et al., 2012; Martin and
3336 Richards, 2001). But there is a consensus that eddies do impact biogeochemistry (McGillicuddy et al., 2007).

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3337 There have been only a handful of studies on the role of eddies in biological productivity in this region (Kumar et al.,
 3338 2007; Singh et al., 2015). Kumar et al. (2007) observed an increase in surface nutrients in the Bay through eddies during
 3339 both fall-2002 and spring-2003 followed by higher biomass. Despite being highly eutrophic, biological activity did not
 3340 increase following cyclonic eddies during the summer-2003 in the northern Bay (Muraleedharan et al., 2007). But primary
 3341 production switched from 'regenerated' to 'new' production during summer-2003. In a ^{15}N based new production estimate
 3342 to assess the role of cyclonic eddies in enhancing primary production, Singh et al. (2015) carried out measurements of ^{13}C
 3343 based primary production at four stations in the Bay of Bengal (around a cyclonic eddy close to 17.8°N , 87.5°E) during
 3344 winter 2007. The measurements sampled one cyclonic eddy during the campaign. The highest surface productivity (2.71
 3345 $\mu\text{M C d}^{-1}$) and chlorophyll a ($0.18 \mu\text{g L}^{-1}$) were observed within the eddy due to intrusion of nutrients from subsurface
 3346 waters. Given new nitrogen input via vertical mixing, river discharge or aerosol deposition, the additional primary
 3347 production due to this new nutrient input and its contribution to the total production increased from 40% to 70%. Eddies
 3348 could be a reason for the otherwise unexplained high new production rates in the Bay of Bengal (Singh and Ramesh,
 3349 2015). Eddies also seem to have a potential for transferring a high fraction of fixed carbon to the deep.

3350 8. Summary and open questions

3351 This paper summarises a suite of new studies in the Indian Ocean that have been made possible through national, bilateral,
 3352 and international programmes, including the IIOE-2. An increase in high quality observations (both increased spatial
 3353 resolution and the acquisition of longer time series) has led to a substantial increase in our understanding of processes and
 3354 interactions. These in-situ observations, in combination with remote sensing, detailed syntheses and modeling have
 3355 increased our knowledge of the surface circulation and its complex implications for biological production, along with an
 3356 increased understanding of air-sea interaction in the Indian ocean.

3357 There are, however, a number of outstanding questions that require prioritised efforts. Compared to the Atlantic and
 3358 Pacific, where the important boundary currents are now being monitored with a suite of gliders with repeated and sustained
 3359 sections (Todd et al. 2019), the boundary currents and their variability in the Indian Ocean remain poorly constrained.
 3360 Given the anomalous warming of the Indian Ocean, the frequency of heatwaves, and the population supported by the
 3361 Indian Ocean and Monsoons, the air-sea fluxes and the coupled atmosphere-ocean exchange in this ocean remain poorly
 3362 understood at many scales. Understanding of the intermediate, deep and abyssal layer circulation and the vertical
 3363 overturning cells that connect these layers in the Indian Ocean is lacking.

3364 There are still many gaps in current understanding of Indian Ocean biogeochemical cycles, which we have presented here
 3365 in the context of the physical processes that affect them. Although the characterization of the temporal and spatial
 3366 variability in chlorophyll concentration and primary production has greatly improved as a result of recent in situ
 3367 measurements and satellite remote sensing, there are still many areas where there is little or no information about how this

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Moved down [53]: These have important feedbacks on the SST which impacts atmospheric convection with a global reach. At longer time scales, the salinity feedbacks to climate at interannual to decadal timescales need to be investigated in further detail. The decadal variability of the Indian Ocean Dipole and its link to the Pacific decadal variability also needs to be better understood

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On the influence of small scale mixing, increased measurements of ocean mixing both along the equator and new long term measurements in the Bay of Bengal, have shown intensively enhanced mixing during the passage of eddies and during cyclones. However, there are still significant uncertainties in subsurface ocean mixing in setting the large scale balance in the Indian ocean. ¶

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3390 relates to changes in planktonic food web structure and particulate organic matter export to the deep ocean. Although
3391 nutrient limitation patterns were not discussed in this review, it should be pointed out that the importance of nitrogen
3392 verses iron and silica limitation in the Arabian Sea and elsewhere in the Indian Ocean is still a subject of debate - more
3393 nutrient and trace metal measurements are needed along with nutrient limitation bioassays throughout the Indian Ocean.

3394 The number of nitrogen fixation rate measurements in the Indian Ocean has increased significantly over the last decade,
3395 but the importance of this process as a source of new nitrogen to the surface ocean has been quantified in only a few
3396 regions (e.g., off northwest Australia) and its contribution to bloom formation (e.g., the Madagascar Bloom) is still
3397 uncertain. From a spatial standpoint, the quantification of biogeochemical variability in the northern Indian Ocean
3398 (Arabian Sea and Bay of Bengal) has benefited, in particular, from numerous shipboard measurements, moorings and
3399 biogeochemical Argo float deployments in the last decade. Many questions still remain, for example, related to the
3400 influence of freshwater inputs on biogeochemical cycles in the Bay of Bengal. Remarkably, the biogeochemical and
3401 ecological impacts of the Indonesian Throughflow have been examined in only a handful of studies. Similarly, there are
3402 very few studies that focus on the biogeochemical and ecological impacts of the Seychelles-Chagos Thermocline Ridge
3403 (SCTR). The ITF and the SCTR are unique features of the Indian Ocean, yet the understanding of their biogeochemical
3404 and ecological impacts is rudimentary at best. Finally, the quantification of biogeochemical variability in the Leeuwin
3405 and Agulhas Currents and adjacent waters has also benefited from recent measurements, though it is important to point
3406 out that the biogeochemical impacts of boundary currents in the Indian Ocean are still poorly understood compared to the
3407 Atlantic and Pacific.

3408 There are still large uncertainties in air-sea fluxes. Even in the regional basin of the Bay of Bengal where there have been
3409 focused international efforts, the river discharge and rain need to be better represented in models, as do the processes that
3410 set the shallow salinity stratification. These have important feedbacks on the SST which impacts atmospheric convection
3411 with a global reach. At longer time scales, the salinity feedbacks to climate at interannual to decadal timescales need to be
3412 investigated in further detail. The decadal variability of the Indian Ocean Dipole and its link to the Pacific decadal
3413 variability also needs to be better understood, particularly given events like the record breaking 2019 positive IOD that
3414 developed independently from ENSO conditions. There are still large gaps in our understanding of the Indian Ocean
3415 dynamics that lead to these extremes, and consequently in our ability to predict the onset, intensity and frequency of
3416 extreme weather such as rainfall and flooding, associated with anomalously strong climatic mode events, that have major
3417 socioeconomic impacts.

3418 Modeling and observational efforts have both pointed to the increased role of air-sea coupling at higher frequencies to
3419 improve the predictions of sub-seasonal Monsoon forecasts. Observations and models indicate that MISOs may be slowing
3420 down because of the warming in the Indian Ocean (e.g. Sabeerali et al. 2013), which needs to be understood better for
3421 providing reliable monsoon predictions and projections in this climate vulnerable region.

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3423 On the influence of small-scale mixing, increased measurements of ocean mixing both along the equator and new long-
3424 term measurements in the Bay of Bengal, have shown intensively enhanced mixing during the passage of eddies and during
3425 cyclones. However, there are still significant uncertainties in subsurface ocean mixing in setting the large-scale balance in
3426 the Indian ocean.

3427 It has been proposed that the hiatus in warming of the surface atmosphere may have ceased as the Pacific Ocean enters an
3428 El Nino like state (Cha et al. 2018). However, the secular trends in the Pacific Ocean trade winds are expected to continue
3429 to affect the Indo-Pacific Ocean heat content through the Indonesian Throughflow (Maher et al. 2018). The Indian Ocean
3430 thus remains a critical component of the Earth's global response to the continued anthropogenic forcing and the ocean's
3431 role as a clearing house for distributing heat to modulate global warming.

3432 **Code Availability**

3433 No original data analyses were undertaken as part of this review paper.

3434 **Data Availability**

3435 No original data analyses were undertaken as part of this review paper. All data presented in this manuscript have been
3436 previously published and are available from sources identified in the original manuscripts.

3437 **Author Contributions**

3438 HEP and AT designed the review, wrote the introductory and concluding parts and sections in their areas of expertise. HP
3439 and AT reviewed the contributions of the authors and made editorial adjustments. RH wrote the sections on
3440 biogeochemical variability in Section 4. All co-authors contributed to the writing of sections relevant to their areas of
3441 expertise and response to reviewer questions. All authors contributed to refining the manuscript for submission. RF, CU,
3442 JB, BW, AS-F, JH and RM contributed editorial advice.

3443 **Competing interests**

3444 The authors declare that they have no conflict of interest.

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