

Can the boundary profiles at 26N be used to extract buoyancy-forced Atlantic Meridional Overturning Circulation signals?

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Abstract. The temporal variability of the Atlantic Meridional Overturning Circulation (AMOC) is driven both by direct wind stresses and by the buoyancy-driven formation of North Atlantic Deep Water over the Labrador and Nordic Seas. In many models low frequency density variability down the western boundary of the Atlantic basin is linked to changes in the buoyancy forcing over the Atlantic Sub-Polar Gyre (SPG) region, and this is found to explain part of the geostrophic AMOC variability at 26N. In this study, using different experiments with an OGCM, we develop statistical methods to identify characteristic vertical density profiles at 26N at the western and eastern boundaries which relate to the buoyancy-forced AMOC. We show that density anomalies due to anomalous buoyancy forcing over the SPG propagate equatorward along the western Atlantic boundary, through 26N, and then eastward along the equator, and poleward up the eastern Atlantic boundary. The timing of the density anomalies appearing at the western and eastern boundaries at 26N reveals ~2-3 years lags between boundaries along deeper levels (2600-3000m). Record lengths of more than 26 years are required at the WB to allow the buoyancy forced signals to appear in as the dominant EOF mode. Results suggest that the depth structure of the signals, and the lagged covariances between the boundaries at 26N, may both provide useful information for detecting propagating signals of high latitude origin in more complex models, and potentially in the observational RAPID array. However, time filtering may be needed, together with the continuation of the RAPID program in order to extend the time period.

1 Introduction

The Atlantic Meridional Overturning Circulation (AMOC) plays a key role in controlling the Earth's energy budget. It transports warm water to the north, overlaying a return flow southward of colder and denser water (Cunningham et al., 2007). Due to its large net heat transport, low frequency variability in this circulation can have an important impact on Atlantic sea surface temperatures and, therefore, on the wider climate (Knight et al., 2005; Sutton and Dong, 2012). Decadal prediction systems have shown that the upper ocean temperatures over the subpolar gyre can be predicted due to the leading role of the ocean heat transport (Robson et al 2012; Hermanson et al 2014; Robson et al., 2017). In order to make these decadal predictions it is essential that we ensure the best ocean initial conditions are available with a well reproduced AMOC.

30 The RAPID program has been monitoring AMOC and boundary densities at 26N since 2004 (Cunningham et al., 2007; McCarthy et al., 2015). The observational record so far has revealed a large range of AMOC variability on different time-scales; from high frequency (Balan Sarojini et al., 2011), to large anomalies persisting at inter-annual time-scales (Blake et al., 2015; Roberts et al., 2013), or decadal trends (Smeed et al., 2014; Jackson et al 2016). Using different methods, several studies have investigated the observed weakening (since 2005) of the AMOC and have related the trends to earlier high latitude density changes (Jackson et al, 2016; Robson et al 2014, 2016).

The AMOC circulation is driven both by direct wind stresses and by the buoyancy-driven formation of North Atlantic Deep Water (NADW) over the Labrador and Nordic Seas. Theories of the response of the AMOC and the ocean gyres to wind stress or buoyancy input rely on energy being transmitted through the ocean by planetary Rossby waves, or along the ocean margins by boundary waves (Johnson and Marshall, 2002; Hirschi et al., 2007; Hodson and Sutton, 2012; Jackson et al., 2016). In particular, changes in the NADW may produce a chain of events in the North Atlantic on a range of time-scales from months to decades. The adjustment has been studied in an extensive literature. Some model studies (Kawase, 1987; Huang et al., 2000; Johnson and Marshall, 2002; Getzlaff et al 2005; Marshall and Johnson, 2013) suggest that AMOC anomalies propagate with boundary Kelvin wave speeds resulting in a very short lead time (of order a few months) between subpolar and subtropical AMOC changes. Roussenov et al (2008) suggested that this boundary propagation may also involve higher mode Kelvin and topographic Rossby waves, leading to longer propagation times (of order years). The advection of the NADW outflow also moves down the western boundary more slowly in the Deep Western Boundary Current (DWBC), although lagrangian float observations show that a large fraction of this NADW moves away from the boundary and enters the ocean interior near the Flemish Cap and the Grand Banks (Bower et al., 2009). Using a coupled climate model Zhang (2010) and Zhang et al. (2011) showed AMOC variations associated with NADW formation propagating more in line with the advection speed, with much longer lead times (several years) between subpolar and subtropical AMOC variations. Getzlaff et al. (2005) have shown that the high latitude adjustment to AMOC anomalies can result from a superposition of a fast wave response and a slower advective signal in ocean model experiments with different resolutions. Interestingly, the speed of propagation along boundaries of the density/velocity anomalies related to AMOC changes is found to be model-resolution dependent (Getzlaff et al 2005; Hodson and Sutton, 2012). These propagating density anomalies will also affect the geostrophic AMOC variability at 26N.

Model simulations clearly show a large range of mechanisms leading to AMOC variability at latitude 26N (Hirschi et al., 2007; Biastoch et al., 2008; Cabanes et al., 2008; Ducez et al., 2011; Polo et al., 2014; Pillar et al., 2016). Buoyancy forcing generally operates from inter-annual to decadal time-scales, while the wind forcing mostly acts from intra-seasonal to inter-annual time-scales (Cabanes, et al., 2008; Kanzow et al., 2010; Ducez et al., 2011; Polo et al., 2014; Pillar et al., 2016). Using an adjoint OGCM, Pillar et al (2016) have found that inter-annual to inter-decadal AMOC variability of ~5 Sv amplitude can be excited by heat fluxes in the subpolar North Atlantic, with freshwater fluxes playing a more minor role. Originating from higher latitudes, the Western Boundary (WB) density anomalies explain most of the variance in the zonal density gradients, and hence

geostrophic transports, at 26N, especially at decadal time-scales (Hirst et al., 2007; Polo et al., 2014). The Eastern Boundary
65 (EB) explains only a small part of the inter-annual variability in zonal density gradients in the upper 1500m, and this is mostly
due to local wind forcing (Polo et al., 2014).

Despite the many model studies showing boundary wave connections between the Labrador Sea and lower latitudes, and their
importance for the AMOC, less work has been done on the vertical structure of these anomalies, and yet it is the vertical
70 density structure at 26N that is primarily measured by the RAPID array. We now benefit from more than 10 years of boundary
density records at 26N, and therefore can consider how best to use the vertical structure in these data to study the lower
frequency variability. If low frequency signals can be identified from the vertical structure this would help us to assimilate the
most important signals that need to be reproduced in climate forecast models. This poses the question we address in this paper;
can we extract the buoyancy forced signals from vertical density profiles, such as those sampled at 26N?

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An earlier attempt to extract buoyancy signals at 26N was made by Polo et al (2014, hereafter PA14) using a NEMO 1° OGCM
forced with full ECMWF reanalysis meteorology from 1958-2010, denoted as the CTRL experiment. This CTRL experiment
was compared to runs using only inter-annual wind or buoyancy forcing, allowing separation of buoyancy from wind forced
variations in the AMOC. PA14 found that the buoyancy-forced AMOC anomalies at 26N could be related to changes in deep
80 water formation over the Labrador Sea some years before. Although they showed a coherent WB vertical signal at 26N, they
were not successful in isolating the buoyancy forced signal in the CTRL experiment, due to the confounding influence of the
wind forced variability. They did not look in detail at the propagation or how the vertical structure associated with the buoyancy
forced anomalies develops. In the present work we extend the work of PA14 by (i) developing statistical means of isolating
the buoyancy forced AMOC variability from the full variability in the CTRL using the density profiles at 26N; (ii) analysing
85 the propagation of the buoyancy forced signals from the Labrador Sea down to the subtropics; and (iii) developing statistical
covariance relationships linking the AMOC to the Labrador Sea that might potentially be used in a data assimilation context
to modify the low frequency AMOC variability. The diagnostics developed are also tested on RAPID observations and on
output from the state-of-the-art coupled model HadGEM3-GC2 (Williams et al 2015), which has a ¼° NEMO global ocean.

90 In this paper Section 2 presents the methodology used to analyse AMOC variability and its sources in several runs of the
NEMO model (1°x1° horizontal resolution) driven by different components of ECMWF atmospheric forcing from 1960-2012,
as in PA14, but now including some validation of the boundary density variability in the CTRL run against the RAPID
observations. Sections 3 and 4 describe the modes of the 26N density profile variability in the model and the associated
propagation occurring upstream and downstream, respectively. Section 5 describes a statistical analysis of boundary densities
95 in the RAPID observations and compares these modes with the NEMO experiments. Section 6 discusses density variability in

the coupled experiments GC2 in comparison to the lower resolution results. Section 7 discusses the results and limitations of the interpretations. Finally, Section 8 summarizes the main conclusions.

2 Methodology

100 This section describes the model experiments and statistical methods used to understand the boundary density variability and its relation to the AMOC variability.

2.1 Forced experiments

105 The forced ocean-only model (hereafter NEMO1) is based on NEMO V3.0; it uses the tripolar ORCA grid in a global configuration with $1^\circ \times 1^\circ$ horizontal resolution and a tropical meridional refinement to $1/3^\circ$. The model has 42 vertical levels with thicknesses ranging from 10m at the surface to 250m at the ocean bottom. Initial conditions are taken from the second iteration of a 50-yr cyclic model spin up, each cycle spanning the period 1958–2009 (Balmaseda et al., 2013). The model is forced with daily atmospheric fluxes as boundary conditions taken from the ECMWF Re-Analysis ERA-40; (Uppala et al. 2005) from 1958 to 1978, and the Interim ECMWF Re-Analysis (ERA-Interim; Dee et al. 2011) from 1979 to 2009.

110 The control experiment (CTRL) is forced with time-varying daily surface heat, freshwater and momentum fluxes for the period 1958–2009. The sea surface temperature (SST) is weakly relaxed to daily values with a relaxation time scale of ~ 1 month, while the sea surface salinity (SSS) is restored to climatological SSS with a time scale of 1 year. There is no ice model; instead, wherever the sea ice concentration in the observations exceeds 55%, the model SSTs are nudged more strongly (1-day time scale) to the freezing point (-1.88°C). The restoration to SSS and SST is stronger under sea ice (30 days and 1 day, respectively).

115 Following the work of PA14 we also have a set of simulations where the momentum and buoyancy forcing is decoupled from one-another. In the experiment referred to as BUOY, the momentum flux is taken from the ERA-Interim 1989–2009 seasonal climatology, while the buoyancy forcing (heat, freshwater flux, and SST) is still inter-annually varying. In the experiment referred to as WIND, the momentum flux is fully varying, but the buoyancy forcing is from the same seasonal climatology. These experiments allow us to identify and distinguish the AMOC signals and processes associated with buoyancy and wind forcing, to the extent that they are independent. We use the BUOY experiment as reference for the buoyancy-forced only signals, and the CTRL experiment as the “truth” which includes both buoyancy and wind forcings, as well as the interaction between them. Where appropriate, we also include the WIND experiment and SUM (as the sum of anomalies from BUOY and WIND). Results are discussed in Sections 3, 4.

125 2.2 Coupled experiment

We also analyse 120 years of monthly-mean data from a control run of the high-resolution coupled ocean-atmosphere model HadGEM3-GC2 (hereafter GC2, Williams et al 2015). The ocean component is NEMO v3.4 with the ORCA025 tripolar grid configuration, using Met Office parameters for “Global Ocean 5.0” (GO5.0, Megann et al 2014), with the CICE sea-ice model. The atmosphere component is GA 6.0 of the Met Office Unified Model (UM; Walters et al 2011) at a resolution of N216
130 (~60Km in mid-latitudes) and 85 levels. This model was used in the Met Office seasonal and decadal prediction systems (GloSea5 and DePreSys3 respectively). The model has been used to study the North Atlantic variability and its predictability (Menary et al 2015; Williams et al 2015; Ortega et al 2017; Robson et al., 2016). Results are shown in Section 6.

2.3 Model evaluation

We use the RAPID array (McCarthy et al., 2015; Smeed et al 2017) to evaluate the boundary densities in the model. We use
135 the merged profiles at the Western Boundary (26.52N, 76.74W) and Eastern Boundary (26.99N, 16.23W) for the period April 2004 to December 2009.

Figure 1 shows the profiles of temperature, practical salinity and density at the boundaries for RAPID and CTRL experiment. The main differences occur at the EB in the upper 1000m, especially in temperature and salinity (Fig. 1a, b), probably due to
140 a bias in the location of the Mediterranean waters. Beneath 2000m, the variables in CTRL and RAPID are very similar.

NEMO1 and GC2 are both able to capture important aspects of the observed boundary density profiles such as the mean vertical density gradients (N^2 , Fig. S1a). On the WB the profiles are similar between 1500m to 4500m but the model stratification is stronger between 300-700m. The EB profiles are similar at all levels below 500m (Fig. S1a). However, the
145 NEMO1 model underestimates the density variance at all levels, especially at the WB, while GC2 has a more realistic variance on the WB at depth (Fig. S1b).

The AMOC at 26N in NEMO1 has a time mean (12Sv) and maximum (18 Sv) at a depth of 1000m in the CTRL experiment. The mean AMOC is higher in the BUOY experiment by ~2 Sv. The AMOC measured at 1000m has a prominent trend in the
150 BUOY experiment (+3.2 Sv in 52 years), but the trend is not significant for the CTRL (-0.2 Sv in 52 years). The AMOC seasonal cycle (not shown) in the CTRL presents a maximum in boreal winter with a secondary peak in boreal summer, which is also reproduced in the BUOY experiment. The annual cycle defined as the difference between the maximum (in boreal winter) and the minimum (in boreal spring) is 3.9 and 4.9 Sv for the CTRL and BUOY experiments. After removing the linear trend the standard deviation of both experiments is similar: 2.28 and 1.92 Sv for CTRL and BUOY respectively (see also Fig.
155 S2 for the AMOC distributions). The AMOC at 26N in the RAPID observations presents a mean and maximum of 17 and 31 Sv respectively from 2004-2014, with monthly standard deviation of 4.35 Sv, which double the standard deviation for the

CTRL experiment. Trends have been reported for RAPID data of -0.6Sv/year (Smeed et al 2014; 2018) which could be part of a longer variation cycle (Smeed et al 2018, Jackson et al 2016). Results from the modes of variability at the western boundary density profile are shown in section 5.

160 **2.4 Statistical analysis**

Model experiments are first sampled at the western and eastern Atlantic boundaries at 26N to simulate the monthly-mean density profiles from the RAPID array. Empirical Orthogonal Function (EOF) analysis of these density profiles is used to obtain vertical modes of density variability and related timeseries of Principal Components (PCs) which together represent the largest fractions of the total variance (Bretherton et al., 1992).

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Before calculating the EOFs, the data are processed to remove the seasonal cycle and linear trends. Unlike in PA14, density anomalies are weighted by the thickness of each layer to ensure that all points are appropriately represented for the total density variability. The EOF analysis is computed for the individual boundaries; Western Boundary (WB) and Eastern Boundary (EB), and also for the combined anomalies at both boundaries. Finally, we have explored the combined EOFs by time lagging the eastern boundary variability in order to understand related signals from both boundaries.

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The PC timeseries is normalised (the units are standardised anomalies) and the associated EOF is represented as the regression map of the density anomalies (in kg/m^3) on to the PC timeseries. Therefore, the EOF pattern has density units (density anomalies by 1 standard deviation of the PC). The PC is calculated from monthly data of the density profiles. The PC from 800m (referred to as PC-f800) is calculated using only the information from 800m downwards. However, the resulting EOF pattern, by computing regression map, can have a density signal shallower than 800m.

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Regression analysis of the PC timeseries on other fields (e.g. 3D density) allows us to detect spatial patterns and depth structures of the propagating modes associated with the EOFs at 26N. We show regression and the correlation coefficients where they are statistically significant at the 95% confidence level, according to a Student's t-test for the effective number of degrees of freedom (Metz, 1991).

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2.5 Spectral analysis

In order to remove the high frequencies in the timeseries in section 3, we have used a one-year running mean filter. This filtered timeseries is obtained by taking the average of a data subset (13 months) which is centered in a monthly time step (von Storch and Zwiers, 1999).

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Spectral analysis is used to decompose timeseries to show signals that lie within different frequency bands. The analysis is performed in order to identify the frequencies involved in the propagation of density anomalies at different depths. Power

spectra of the timeseries are obtained using the multi-taper method, which provides more degrees of freedom and therefore
190 more significance (Thomson 1982). The power spectra are tested against the hypothesis that the signals are generated by a
first-order autoregressive process AR(1) with the same time-scale as the original, yielding a red noise spectrum, and the 95%
confidence limit for the rejection of the red noise hypothesis is applied. Additionally, when we have a large internal variability,
we use the decomposition in order to filter some of the timeseries using a Lanczos (1956) filter. This is done in particular for
the control GC2 run simulation.

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We have used the Radon Transform (RT) function (Dean 1983) in order to estimate the phase speed of propagation of density
anomalies. The angle of the maximum RT standard deviation determines the propagation phase speed. We calculate the RT
every 0.1 degrees. The phase speed averaged for the Hovmöller diagrams has been estimated at 3000m level for CTRL and
BUOY experiments.

200 **3 Modes of vertical density variability at the 26N boundaries**

3.1 Linearity of AMOC and the boundary density signals in the NEMO1 model

The forced ocean model experiments enable us to isolate the boundary density variability associated with buoyancy forcing
i.e. from the BUOY experiment, as noted by PA14. We focus here on developing the vertical density fingerprint of this signal
and using it to identify the buoyancy forced AMOC signal as it appears in the CTRL run.

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Figure 2a shows the monthly AMOC (the total AMOC minus Ekman component) variability, defined as the integral of the
meridional transport at 26N down to 1028m, for both the CTRL and BUOY experiments (as in Fig. 1a in PA14). There is a
prominent decadal signal in both CTRL and BUOY with peaks in 1975, 1985 and 1995, although CTRL also shows additional
monthly and inter-annual variability. Similar decadal signals of AMOC variability reproduced by reanalysis products (Wang
210 et al 2010; Karspeck et al., 2015). The monthly-mean timeseries correlate at 0.43, but the correlation rises to 0.62 when using
a 1 year running mean filter (Fig. 2b). Wind forced inter-annual variability explains most of the remaining differences; when
the 1-year smoothed AMOC anomalies from BUOY and WIND are summed the correlation with CTRL rises to 0.86 (SUM
in Fig 2b).

215 Most of the boundary density variability is also recreated in BUOY and WIND. The correlations between boundary density
anomalies in the CTRL and SUM are shown as a function of depth in Fig. 2c-d. For the WB, most of the CTRL variability is
reproduced by SUM from 1800m to 4000m (Fig. 2c and Fig. S3). For the EB SUM explains most of the variability seen in the
CTRL experiment at all depths (Fig. 2d and Fig. S3). A 1-year low pass filtering does not influence the correlations for the
WB, although for the EB filtering reduces the correlation at some depths. We now relate this density variability with the
220 AMOC signals in Fig 2a, b.

3.2 EOFs of boundary density profiles

Figure 3 shows the Principal Component timeseries of the first EOF computed using monthly density profiles on the western and eastern boundary at 26N for the CTRL (blue) and BUOY (red) experiments. As density profiles near the surface contain significant noise, we calculate the PCs for both full depth profiles (f0m, Fig. 3a-b) and only from 800m downwards (f800m, Fig. 3c-d). The full depth PCs in Fig. 3a both show substantial high frequency noise, and the CTRL timeseries does not correlate well with AMOC variability from Fig 2a. However, when only retaining densities below 800m, in Fig. 3c, the BUOY and CTRL PCs closely match each other ($r=0.80$) and both now correlate with the AMOC timeseries in Fig. 2a ($r=0.33$ and $r=0.88$ for the CTRL and BUOY respectively). These f800m EOFs also explain more of the deeper density variance. Note that, adding a 1-yr filter to the PC-CTRL timeseries increases the correlations with the CTRL AMOC variability to 0.47 (Fig. 3c). However, for the BUOY experiment the temporal filter has little impact on the correlation between PC and AMOC timeseries (0.89).

Figures 3b and d show the corresponding eastern boundary density EOF timeseries for full depth and below 800m variability. The full depth PC is rather different to that on the western boundary, and also to the AMOC variability itself. However, when only the deep variability is retained, the inter-annual variability in BUOY is more similar to the western boundary buoyancy driven variability. In CTRL there is still considerable high frequency wind generated variability in the deep PC, however, when a 1-yr filter is introduced the buoyancy forced AMOC related signal becomes clearly visible on the eastern boundary, and the correlation with the AMOC rises to 0.42 (Fig. 3d). All correlations are summarised in Table 1.

In order to understand the usefulness of truncating the density profiles for the EOFs, Fig. 4a-b summarises correlations between leading PCs at the boundaries and AMOC timeseries in the CTRL experiment by increasing the truncation level for the density profiles. For the WB, maximum correlation is found simultaneously at all truncations especially for deeper levels (Fig. 4a). In contrast for the EB simultaneous correlations are always low even with a 1-year filter, but correlations increase greatly when a lag to the AMOC is applied (Fig. 4b). Correlations still require a 1-year filter to remove noise but now peak at 0.7 with a lag around 2-3 years, with the more deeply truncated signals also showing the longer lags. Similar lag increased EB correlations are found for the BUOY experiment (not shown). The nature of this lag in the EB-AMOC correlation is related to signal propagation between boundaries and will be discussed in the next section.

Figure 4c-d shows the vertical structure of density anomalies associated with the leading EOF modes for both boundaries using the 800m depth truncation and monthly-mean data for the three experiments CTRL, BUOY and WIND. The leading EOFs are very similar between CTRL and BUOY on the WB, with maxima between 1200m and 4000m. Note that in PA14 Fig. 3b, their PC-WB in the CTRL experiment with only 500m depth truncation was substantially different and was mainly wind-driven. The WIND experiment has much less variability at greater depths, and the PC is uncorrelated with the PC-WB from CTRL

($r=0.09$, Fig. 4c). Therefore, on the WB, Fig. 3c and 4c show that the EOF analysis successfully extracts the buoyancy-forced signal related to the AMOC in the CTRL.

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On the EB the leading EOFs, even below 800m, show more correspondence between CTRL and WIND (Fig. 4d, with PC correlations $r=0.87$). This explains why further filtering and the application of a lag to the PC timeseries in Fig. 3d, 4b is needed to extract the weaker buoyancy-forced AMOC-related signal on the EB. The relationship between these buoyancy driven density variations at both boundaries is now explored further.

260 **3.3 Relationship between boundaries**

Figure 5a shows lead-lagged correlations between the WB leading PCs (computed from 800m depth) and the EB at different depth truncations, in the CTRL experiment. Dashed (solid) lines indicate PCs without (with) the application of a 1-year running mean filter. For truncations deeper than 1600m the highest correlation is found when the EB is lagged up to 30 months (Fig. 5a), revealing longer links between boundaries for deeper levels. However, this lag is much less clear when shallower depths are retained, and substantial WB-EB density correlations start to be seen with lag0, unlike in Fig. 4b. Even in the BUOY experiment, only when the EB EOFs are truncated to below 1600m is a strong lag clearly seen between the boundaries, again reaching up to 30 months for the deepest signals (Fig. 5b).

In order to reduce upper level noise at the EB and to bring out the deep density signal connecting the boundaries more clearly, we compute the combined EOF while truncating the EB to below 1600m, which shows a maximum in the WB-EB correlations for CTRL (Fig. 5a). The new combined EOF shows similar WB structure for all lags (Fig. 5c) with a deeper signal around 2000m on the EB (Fig. 5d). Figure 5e shows the timeseries, which is very similar to the PC1-WB in Fig. 3c. The new combined EOF explains more density variance in CTRL (43% compared to 40% explained by PC1-WB) and a slightly longer lag between WB and EB (18 months in CTRL and 25 months in BUOY, not shown).

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We conclude that the deep densities on the EB contain a very clear signal of the buoyancy forced AMOC variability but that this signal plays no detectable role in the direct (lag 0) control of the AMOC. The EB lag signal can also be seen in relation to the PC-WB density, although the signal is less clearly lagged, probably reflecting the noise still present in the upper layer densities on the WB.

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The WB clearly contains the core density information on the buoyancy-forced AMOC changes at low frequencies. Using PC1-WB, we will now identify the propagating density signal connecting the boundaries (section 4) and search for similar signals in the RAPID data (section 5), and in the higher resolution GC2 model (section 6).

4 Propagation of the buoyancy-driven signals

285 Motivated by the lagged signal at the EB, we analyse the i) spatial coherence of the anomalies at deeper levels (~3000m), and
ii) the propagation fingerprints of the connecting signals.

4.1 Spatial regression patterns

Figure 6 shows the spatial density anomalies averaged 2700-3000m, regressed onto the PC1-WB for the CTRL and BUOY
experiments respectively, at different lags. At these deep levels the regression patterns in CTRL (Fig. 6a-c) and BUOY (Fig.
290 6d-f) are very similar to each other. This agreement suggests that, although the magnitude of the regressions is stronger in
BUOY, we are identifying the same signal in both experiments. The density signal in the Labrador Sea appears from lag -30
(Fig. 6a, d) and intensifies and propagates down the WB to the Equator (Fig. 6b, e), and then across the equator and poleward
at the EB (Fig. 6c, f).

295 Density regressions at shallower levels (900-1300m) against PC1-WB in the CTRL experiment (suggested by the maximum
in the EOF profiles in Fig. 4) also finds anomalies beginning near the Labrador Sea, leading PC1-WB by 30 months, and a
pattern of equatorial Kelvin and Rossby-waves as in Johnson and Marshall (2002) from lag -30 to lag ~0 (Fig. S5). However,
the absence of such tropical signals in BUOY at long lags (lag -30) shows that the CTRL signals could be due to wind driven
Ekman pumping (see supplementary material in Fig. S6). Therefore, we concentrate on the deeper signal which is clearly
300 related to buoyancy.

4.2 Wave path and phase speed

Figure 7a shows the wave-track defined following the signals in Fig. 6 and using the topography at the 3269 m model level.
The wave-track starts in the Labrador Sea (60N) and proceeds southwards to the equator. We plot it along the equator and then
North along the eastern boundary to 55N. The path avoids entering into the Gulf of Mexico, but at these depths this is not
305 expected.

The Hovmöller diagram along this path (Fig. 7b) shows the propagation related to peaks in the deep water formation at high
latitudes for the CTRL experiment (BUOY is very similar, not shown). Density anomalies propagate continuously along the
track from the Labrador Sea around to the British Isles. The propagation shows density maxima in 1975, 1985 and 1995, also
310 seen as peak-AMOC years in Fig. 1-2. Additionally, the propagation speed from the Radon transform is similar for both
experiments (0.41 and 0.31m/s for the BUOY and CTRL respectively). This phase speed is consistent with the lags found
between boundaries (i.e. density anomalies in CTRL will take ~25 months to travel between WB and EB following the defined
track). Background currents at 2700-3000m level in the CTRL are shown in supplementary material (Fig. S4).

315 Figure 8 shows the regression of density anomalies from Fig. 7b onto PC1-WB for the CTRL and BUOY experiments. The density anomalies in both experiments show a continuously propagating pattern from the Labrador Sea right around to 40N on the eastern boundary. The signal is stronger in BUOY, as noted in Fig. 6, but otherwise the regression patterns are very similar.

Importantly these propagating buoyancy-related signals are clearly seen in the CTRL experiment where wind and buoyancy
320 forcing are both applied, again suggesting that the analysis is extracting the same buoyancy forced processes. Therefore, the diagnostic methods developed should allow identification of similar signals in other models, and in the observations. In the next section we look for similar buoyancy-forced signals in the RAPID observations.

5 Modes of WB density variability in RAPID data

The RAPID timeseries dataset is considerably shorter than the datasets analysed for buoyancy signals in the NEMO model.
325 Nevertheless Fig. 9 uses the same EOF analysis on the WB density profile for the RAPID array and for the NEMO experiments for the common period 2004-2009. The first two leading EOFs and their PCs indicate inter-annual variability in RAPID data (black lines in Fig. 9a, c). EOF1 has a maximum density anomaly at ~1000m (with a maximum of 0.04 kg/m³ in Fig. 9b), while the second mode describes variations at deeper levels (~3000m, Fig. 9d).

330 The density EOFs of the WB in both the CTRL (blue lines) and WIND (green lines) experiments for the common period 2004-2009 look quite similar to the EOFs from the RAPID array, in particular with EOF1 peaking at around 1000m and EOF2 peaking much deeper at ~3000m (Fig. 9b, d). However, the density anomalies in the RAPID modes are larger than in the model modes (larger amplitude of density variance in RAPID compared with CTRL is also seen at all depths in Fig. S1).

335 However in BUOY (Fig. 9, red lines), although the deep 3000m peak shows up as EOF1, it has very small amplitude (0.001 kg/m³) showing very little variability over the short 2004-2009 period. EOF1 is associated with a decrease up to 2008 and an increase thereafter, up to 2010 (Fig. 9a, similar to the signal at the end of the period in Fig. 3c). BUOY EOF2 shows very little density signal at depth, and mainly represents interannual variability (Fig. 9c-d).

340 The PC1 timeseries in CTRL and WIND also look similar to PC1 from RAPID, reaching a peak in 2007 and declining to 2010 (Fig. 9a). Therefore, for these common 6 years of simulation 2004-2009, the wind-forced inter-annual density variability on the WB is remarkably well captured by the model.

Figure 10 shows the statistics of the correlation between the leading WB mode for CTRL and both, WIND (Fig. 10a) and
345 BUOY (Fig. 10b) experiments using different numbers of years for each sub-period (x-axis). Figure 10a (first box) shows the test of sampling 6 years periods (as we have in 6 years in the common period in Fig. 9a). For all 6-year sub-periods selected

between 1958 and 2009 we found that EOF1 and PC1 in CTRL and WIND agreed well (Fig. 10a, first box), even in periods when the buoyancy signal was known to be changing rapidly. This dominance of the wind forcing over short time periods is not surprising and was noted previously (PA14).

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We find that typically 16 years of data are needed to find a significant (above 0.35) correlation between CTRL and BUOY (Fig. 10b), however the leading mode in CTRL may still be a mix of wind and buoyancy signals (as seen from Fig. 10a). The best extraction of buoyancy forcing signals occurs when we have periods longer than 35 years, when the wind-forced signal nearly disappears (Fig. 10a). Figure 10c shows the correlation between the PC1-WB and the AMOC timeseries for the CTRL over different periods. This shows how much AMOC variance can be explained by the PC1-WB. Only for periods of >25 years is PC1-WB in the CTRL experiment able to explain more than 25% of the AMOC variance ($r > 0.5$) and is therefore able to extract an AMOC-related signal at 26N (Fig. 10c). We note that as the sub-periods get longer, with more years (x-axis), we have fewer independent sub-periods and therefore less dispersion in the sample distribution (i.e. the box is smaller). However Fig. 10 still suggests that for long periods the correlation between PC1-WB CTRL and PC1-WB BUOY is high, correlation between PC1-WB CTRL and PC1-WB WIND is low, and thus the WB density profile associated with PC1-CTRL is mainly buoyancy-forced signal.

360

Although EOF2 from CTRL and RAPID both represent deeper density variability (Fig. 9d) they are still dominated by wind forcing over this short time period, and we note that EOF2 from WIND shows the same deep density peak. The PC2s from CTRL and WIND show very similar timeseries and even the RAPID PC2 shows a considerable level of agreement with CTRL and WIND (Fig. 9c). However, PC1 from BUOY, which has a similar EOF but represents only the buoyancy forced component, has only lower-frequency changes with no relationship to the other timeseries (Fig. 9a). The BUOY PC2 timeseries (Fig. 9c) also shows no comparable variability.

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Therefore, the short record of the RAPID array would not allow us to follow buoyancy-forced signals from the NEMO1 model. Hence, it appears likely that the variability seen here in the RAPID record, both at shallow and deeper depths, is mainly related to wind-forcing (both PC1 and PC2 in CTRL experiment and RAPID show agreement in Fig. 9). Note that the same EOF analysis using the longer record now available for RAPID data (2004-2018), but not for these model results, still gives similar density profiles to Fig. 9 (not shown).

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The NEMO1 model results suggest that the lower frequency buoyancy forced signals from higher latitudes may start to dominate over the wind forced signals after ~25 years of RAPID data have been collected, when their leading density variability should show up at deeper depths ~3000m. In the next section we look at the ability of the analysis to extract buoyancy forced signals, and their propagation, in a higher resolution HadGEM3-GC2 coupled model run, which is the current UK operational coupled model.

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6 Boundary density in a high resolution coupled model

6.1 Density propagation in GC2

Figure 11a shows a wave track for the GC2 model (section 2.2) bathymetry using the 3138m level boundary, and figure 11b shows the Hovmöller diagram of the density anomalies for 120 years along this track, after applying a 2 year low pass filter as GC2 is much noisier at higher frequencies (see methods). The high frequency density anomalies (<2years period) in a similar high-pass Hovmöller (not shown) are very noisy and do not propagate, therefore we suggest that the high frequency signal is dominated by local Ekman pumping.

Anomalies propagate down the WB from ~40N (track point 200) to the equator and across to the EB. These signals can be traced back across the Gulf Stream to subpolar latitudes (points 100-200), but only appear as lower frequency decadal variations in the subpolar gyre and into the Labrador Sea (points 0-100). The Radon Transform phase speed is ~2m/s, which is faster than the phase speed calculated in NEMO1 at the same depth (Fig. 7), and closer to the theoretical and observed Kelvin wave propagation speeds (Illig et al 2004; Polo et al 2008). Other authors, using models of different complexity, have found a similar phase speed of density anomalies from AMOC variations propagating along the western boundary (Getzlaff et al 2005; Zhang, 2010; Marshall and Johnson, 2013). Therefore, as the density anomalies propagate down the deep western boundary, we would expect to find this deep density variability signal using the EOF analysis.

6.2 EOFs at the WB in GC2

Figure 12a shows the vertical density profiles associated with the first two EOFs from 800m downwards at the 26N WB in the GC2 experiment. The profile location is the closest grid-point to the western wall at the Bahamas, which reaches the bottom at 3200m, and the first 2 EOFs explain more than 70% of the total density variance. EOF1 shows an equivalent barotropic vertical structure peaking near the bottom ~3000m (blue line) while EOF2 changes sign between 900m and 3000m (red line).

The PC1 timeseries associated with EOF1 is plotted in Fig. 12b (blue line). Unlike in NEMO1, this PC1 shows high frequency variability which is nevertheless still correlated with the AMOC-Ekman (i.e. the peak AMOC stream function at 900m after the variability due to Ekman has been removed) without filtering at $r=0.45$, rising to $r=0.49$ with high-pass (<2 years) filtering. However, PC1 becomes less correlated with the AMOC after 2 year low-pass filtering, $r=0.25$ (Table 2). In contrast, for PC2 the unfiltered correlation with the AMOC-Ekman is low ($r=0.13$) but this increases with 2 years low pass filtering ($r=0.32$, Fig. 12c, Table 2). As density anomalies at deep levels are able to be excited by wind alone (already seen in both NEMO1 and the RAPID observations in Fig. 9), the two EOFs in GC2 each capture some wind and some buoyancy forcing.

Although the PC1 and PC2 are orthogonal by construction, and thus correlation between timeseries is zero, after time filtering the PCs the modes are correlated ($r=0.41$) and the lead-lag correlations with the AMOC present a cycle between vertical profiles modes, with 16 months between the peaks (Fig. 12d). This indicates the limitations of extracting low frequency AMOC-related signals in more complex environments using linear methods.

The high frequency PC1-WB (<2years) represents high frequency density signal at 26N, which could be wind-forced. It is correlated with the AMOC and is independent of the number of years used to identify it (Fig. 12e). PC2-WB correlates better with a lower-frequency AMOC signal (>2years) and is also independent of the number of years used to identify it (Fig. 12f).

If we filter the density anomalies prior to performing the EOF, then the leading mode corresponds to PC2-WB seen here. This confirms that we cannot isolate the low frequencies by identifying a deep density signature, as works well in NEMO1, therefore time-filtering is needed to identify the inter-annual buoyancy signal in GC2 as a leading mode. Nevertheless, the buoyancy signal is still traceable emerging from the Labrador Sea, and is well captured in the PCs, representing relevant information on AMOC variability.

Here we summarise the comparison between density profile modes in different environments as follows:

i) In the NEMO1 CTRL simulation with 52 years of data, we can find a vertical density pattern at 26N that is buoyancy-forced (i.e. similar to the mode in the BUOY experiment). This has maximum density anomalies at 1500-3000m and corresponds to low-frequency variability of the AMOC. The source is the density changes over the Labrador Sea. In comparison, the leading density mode for the WIND experiment has a maximum at 1000m and represents inter-annual variability of the AMOC.

ii) However, if the same methodology is applied over shorter periods, the wind-forced variability dominates both the first 2 EOFs, with density anomalies at 1000m and 3000m. The limit for the time-period is about 25 years to extract buoyancy-forced signals that can be related to the AMOC at lower frequencies.

iii) When CTRL is compared with the RAPID array PCs, we find similar vertical profiles in the CTRL and observed PCs. Suggesting the short period of RAPID does not allow to extract relevant buoyancy-signals. Wind-forced signals are predominant showing density anomalies that are also relevant for the geostrophic part of the AMOC at 1000m (but without any lag and at inter-annual timescales).

iv) The same methodology applied to a more complex GC2 earth system model results in a leading PC that shows positive anomalies between 1000 and 3000m. This PC1 is related to the AMOC at short inter-annual timescales (predominantly wind-forced). The PC2 mode shows reversing density anomalies between 1000 and 3000m and is related to the AMOC at lower

frequencies (period >2 years). The analysis shows that these PCs1 are correlated to AMOC independently of the number of years used in the calculation, although correlation with PC1 have less spread for sub-periods longer than ~26 years (Fig. 12e-f).

450 The EOF below 800m method seems to be appropriated to detect buoyancy forced signals from density profiles if we have more than ~25 years of data in NEMO1. However, a Lanczos time filtering for periods > 2years prior to the EOF analysis is recommended in more complex environments, as would be required for RAPID observations, and for coupled models. From similarities between NEMO, RAPID and GC2, we conclude that the density profile mode that is likely to be buoyancy-forced, corresponds to density anomalies at deep levels (3000m) that covary negatively with density anomalies at upper levels
455 (1000m). Timeseries should be filtered with periods >2years and the PC should then correlate with the AMOC at 26N 1000m for the same frequencies.

7 Discussion

In this work we have used model output and statistical methods to identify vertical density profiles along the boundaries that are consistent with the buoyancy-forced variability. We have shown that the most relevant profile at 26N is found at the WB using EOF analysis after truncating the density profile from 800m (PC1-WB). This truncation is very effective in emphasising low frequency (decadal timescale) signals, and in NEMO1 negates the need for temporal filtering, which can also add spurious signals or lead to excess smoothing. Caveats that warrant further discussion include the differences between the EOFs in NEMO1 and GC2, and the role of the EB.

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We note that vertical profiles of the WB EOFs in GC2 and NEMO1 are different, especially in the top 1500m, and the frequencies of the dominant variability in boundary signals in GC2 are higher than the decadal signature seen in NEMO1. We find that shallow (<1500m) density signals related to PC1-WB in NEMO1, have different timing in CTRL and BUOY experiments in the tropics (Fig. S5-6), suggesting that at 26N wind forcing is modulating the buoyancy-forced density signal in CTRL. This is also an argument to suggest that in GC2 the shallow signal is probably wind-forced signal.

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Although PCs in GC2 contain the low-frequency AMOC-related variance it is perhaps not surprising that the details of the WB EOFs are different to those in NEMO1, given the range of AMOC variability in models (Biastoch et al., 2008; Cabanes et al., 2008; PA14; Ortega et al., 2017). The 1° horizontal resolution, and even the ¼° model, may still be too coarse to correctly capture propagating boundary signals (Johnson and Marshall, 2002; Getzlaff et al., 2005; Hodson and Sutton, 2012) from the Labrador Sea. Therefore, we may not expect the exact details of the boundary density EOFs, or the phase speeds or phase lags identified from the boundary and Labrador Sea signals, to be very realistic. Nevertheless, the methods reveal, in two very different modelling environments, boundary signals consistently related to the geostrophic AMOC at 26N.

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Regarding the differences in the phase speed propagation in NEMO1 and GC2, there are factors that influence propagation (and hence the phase speed) of boundary waves in numerical models such as i) model resolution (Hsieh et al., 1983; Hodson and Sutton, 2012). ii) Lateral viscosity (Davey et al., 1983). iii) Orientation of the coastal boundary relative to the ocean grid (Schwab and Beletsky, 1998). iv) ocean stratification along the propagation trajectory. v) Additional complexity in a coupled model in comparison with an ocean-only model, potentially influencing the propagation. We think that both resolution and complexity are explaining the differences in the phase speed along boundaries shown in the Hovmöller plots (Figures 7 and 11), however, specific sensitivity experiments should be designed in order to properly understand these factors.

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It is also worth noticing that, due to noise from high frequency wind-forcing, time filtering is needed to see a more clear buoyancy related signal in GC2. Wind forcing can be projected onto density anomalies at deep levels (as is seen in the RAPID data and in GC2), and similar time-filtering may be needed in the observations in order to extract the buoyancy-forced signals.

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Understanding the differences between these boundary density EOFs between models, would be useful for interpreting observations as the RAPID record becomes longer, and it should be a focus of further work.

495 In Section 3.3, we tested the value of using WB and EB together in a single EOF. The fact that combining the WB and EB at different lags does not improve the explained variance shows that the WB alone captures most of the variability at decadal time-scales. This is in agreement with previous results by PA14 and also with the propagating signal from the Labrador Sea that can reach the subtropics along the WB (as in shown by Jackson et al., 2016), but the propagation up the EB is less clear. This is not to say that the EB observations do not make up an important component of the RAPID array observations. Indeed, it has already been shown, in observations (Kanzow et al, 2010; Ducheze et al 2011), and in models (PA14), that the EB is
500 important for understanding the wind-forced variability in the observed AMOC at 26N from sub-annual to inter-annual time-scales. Here we were able to show the density propagation at 3000m depth on the EB which has a clear AMOC-related signature, although with a lag. Therefore, the EB observations are still important in understanding the role of decadal-buoyancy forced variability.

505 We want to comment that the correspondence of the PCs and boundary wave modes is not direct. The boundary EOFs only capture variability at single locations (26N). Figure 4 shows that the teleconnection between density signals at the WB and EB occurs very quickly at 800m-1100m (around 8 months as a fast response) however, at deeper levels the EB signal is delayed. The different lags in Fig. 5 reflect different propagation speeds at different depths, mainly along the equator (Hovmöller diagrams at different depths suggest that the phase speeds change at the equator). The propagation along either the WB or the
510 EB represent boundary waves with more coherent EOF vertical modes, albeit different modes on the WB and EB.

The temporal variability in RAPID WB density profile suggests that wind forcing is still dominating variability in the short record. More years of RAPID data may be needed to allow the buoyancy forced variance to appear in the leading EOFs. Although it may be helpful to perform time filtering which could allow buoyancy forced signal to be found in shorter periods
515 of data, the continuation of the RAPID array would be crucial in order to understand the wind-forced inter-annual variability and to detect the density signals linking the subpolar North Atlantic with the AMOC at 26N.

8 Conclusions

In this work we have used NEMO1 OGCM experiments which separate buoyancy and wind forced signals (BUOY and WIND experiments; Polo et al., 2014), together with statistical techniques, to develop methods to extract the Atlantic boundary density
520 profile signatures at 26N most associated with the buoyancy-forced AMOC from an experiment with both buoyancy and wind forced variability (CTRL). After finding the “best” vertical profile on the western boundary, we describe the temporal-spatial structures related to this signal. The main findings are summarized as follows:

- 525
- Using EOF analysis and outputs from OGCM experiments we find that the vertical density structure at both the western (WB) and eastern boundaries (EB) at 26N show characteristics that can be unambiguously linked to buoyancy-forcing in the Labrador Sea.
- 530
- The vertical structure associated with the leading EOF mode of density variability on the WB (EOF1-WB) shows positive anomalies from 1500 to 3000m depth that can be related to earlier changes in the North Atlantic deep water formation, and to density anomalies over the Labrador Sea, which are seen to lead PC1-WB by ~30 months. The PC1-WB is found to be very robust in both the CTRL and BUOY experiments, driving buoyancy-forced AMOC variability on decadal timescales. PC1-WB explains 40% and 70% of the 26N deep density variance for the CTRL and BUOY experiments respectively.
- 535
- The PC1-WB is found to lead density anomalies at 1000-1500m on the EB (associated with PC1-EB) by ~7months. The result of combining both boundaries into a single density EOF allows to extract the correlated variance and the optimal lag between the boundaries. The combined EOF variance shows maxima when lagging the EB between 7 months and 3 years. The longer period lagged relationship is consistent with density propagation at 2700-3000m.
- 540
- In the CTRL experiment, density anomalies at 2000-3000m propagate southwards along the WB and eastward along the Equator and then up to the African coast impacting the vertical structure on both boundaries at 26N. The propagation is continuous from the Labrador Sea around the basin and up to the British Isles. This density signal propagates at an average speed ~0.3m/s, consistent with the propagation speeds in the BUOY experiment.
- 545
- The same method is applied to the RAPID array data for the common period with the NEMO1 simulations. The two leading EOFs for the WB have anomalous densities at 1000m and 3000m respectively and are well simulated by the CTRL experiment (Fig. 9). This inter-annual variability is unequivocally wind-forced. The observational record must be longer in order to identify the buoyancy-forced vertical anomalies, which have lower frequencies.
- 550
- The same method was able to extract boundary signals from the higher resolution model HadGEM3-GC2. Despite the greater complexity in GC2, the vertical density profiles on the WB at 26N can be clearly related to the geostrophic AMOC, although some time-filtering is needed in order to separate the timescales.

- 555
- After filtering (periods below and above 2 years), PC1-WB is found to be more related to AMOC at 1000m at high frequencies, with an EOF profile with strong positive density anomalies at 1000m and 3000m. PC2-WB is related to the AMOC at 1000m at low frequencies, and shows positive (negative) density anomalies at deep 1000m (3000m) levels.
- 560
- We also show clear density propagation from the Labrador Sea around the basin to the British Isles along a wave track at 3000m depth (defined by 3138m bathymetry), which explains part of the AMOC variability. However, temporal filtering is needed to make this stand out above the noise.
- 565
- We conclude that the buoyancy-forced density signals at 26N will be distinguishable in the observations (as well as in coupled models) if the available record is long enough (>26 years), selecting density profiles with opposite anomalies at 1000m and 3000-3500m, with time filtering >2 years to help eliminate high frequency wind-driven signals.

Author contribution

570 IP has made the calculations with the data from numerical simulations and observations. IP, JR and KH have analysed the results and IP prepared the manuscript and all authors have contributed to the writing and editing the manuscript.

Competing interests

The authors declare that they have no conflict of interest.

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Code and data availability

The code used for the analysis is done in Matlab and it could be available by request to the first author. The data for the three forced experiments belongs to ECMWF and it could be available by request. GC2 data belongs to MetOffice and it could be available by request. RAPID data is fully available online at: <http://www.rapid.ac.uk>.

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Table 1. Correlations between PCs f800m and AMOC timeseries in NEMO1. Correlations between the geostrophic AMOC and the PCs at the boundaries for the CTRL and BUOY experiments. Significant correlations are highlighted with a star based on a t-test at 95% confidence level and after calculating the number of effective degrees of freedom.

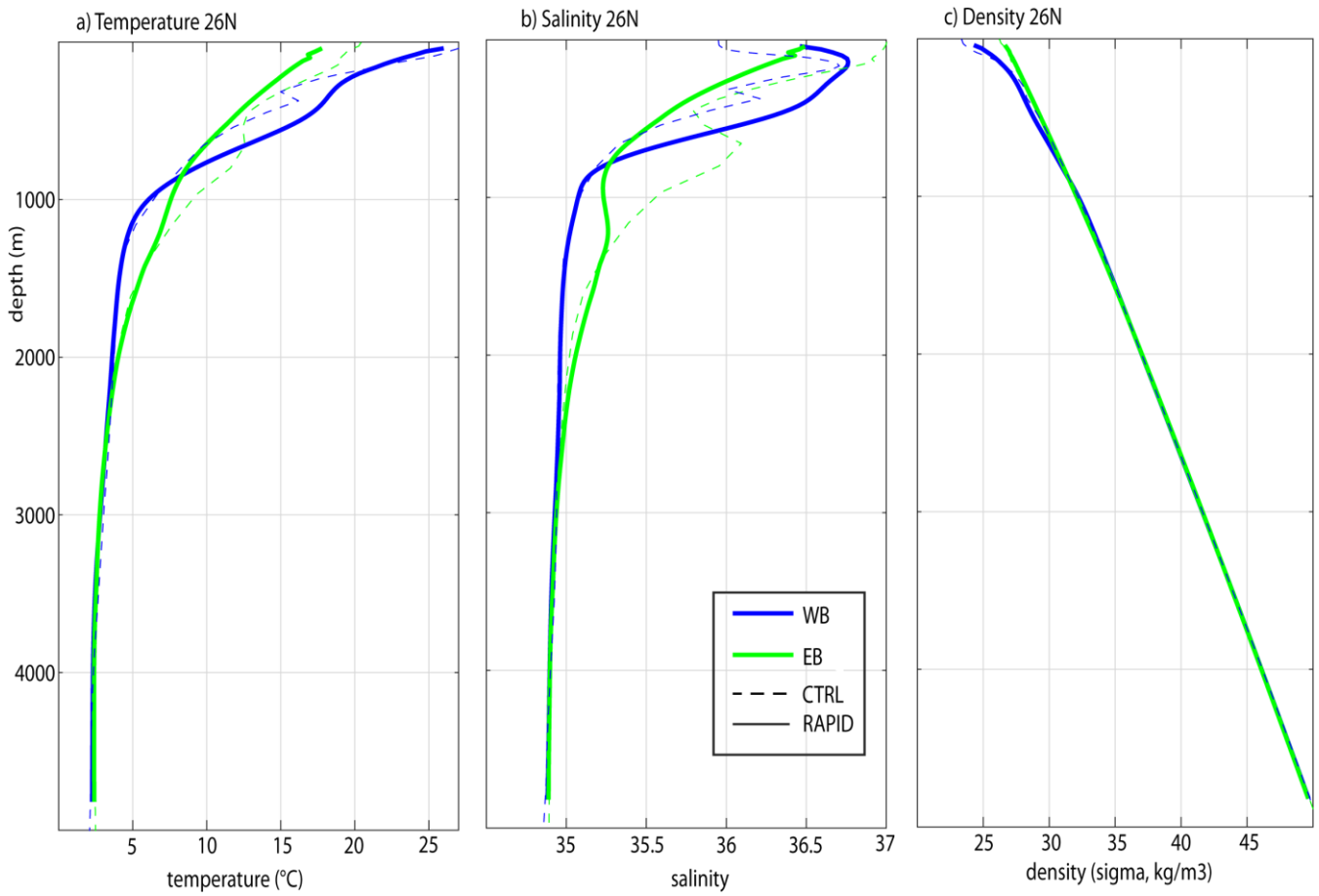
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Correlation Nf / 1yrF	PC1-WB AMOC	PC1-EB AMOC	PC1-EB- PC1-WB
BUOY	0.88*/0.89*	0.80*/0.81*	0.83*/0.84*
CTRL	0.33*/0.47*	-0.04/0.42*	0.36*/0.67*

Table 2. Correlations between PCs f800m WB and AMOC timeseries in GC2. Correlations between the geostrophic AMOC and the PCs at the WB for the GC2 experiment. Significant correlations are highlighted with a star based on a t-test at 95% confidence level and after calculating the number of effective degrees of freedom.

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Correlation	Nf	< 2 years	>2 years
PC1-WB, AMOC	0.45*	0.49*	0.25
PC2-WB, AMOC	0.13	0.06	0.32*
PC1-WB, PC2-WB	0	0.11	0.32*



795 **Figure 1. Temperature, salinity and density profiles at 26N.** a) Temperature profiles (in °C) at 26N for Western Boundary (WB, blue line) and Eastern Boundary (EB, green line) for the CTRL experiment (dashed line) and the RAPID data (solid line) for the period 2004-2009. b) same as a) but for the salinity profiles. c) same as a) but for the density profiles (in kg/m³).

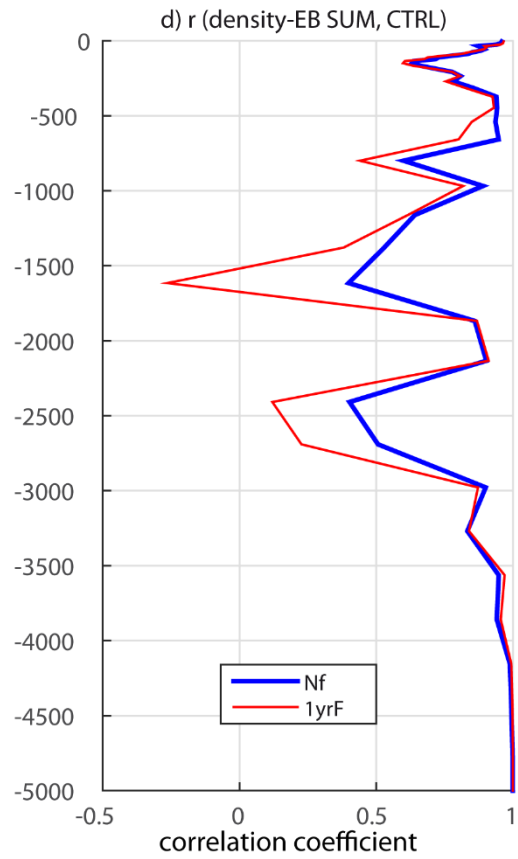
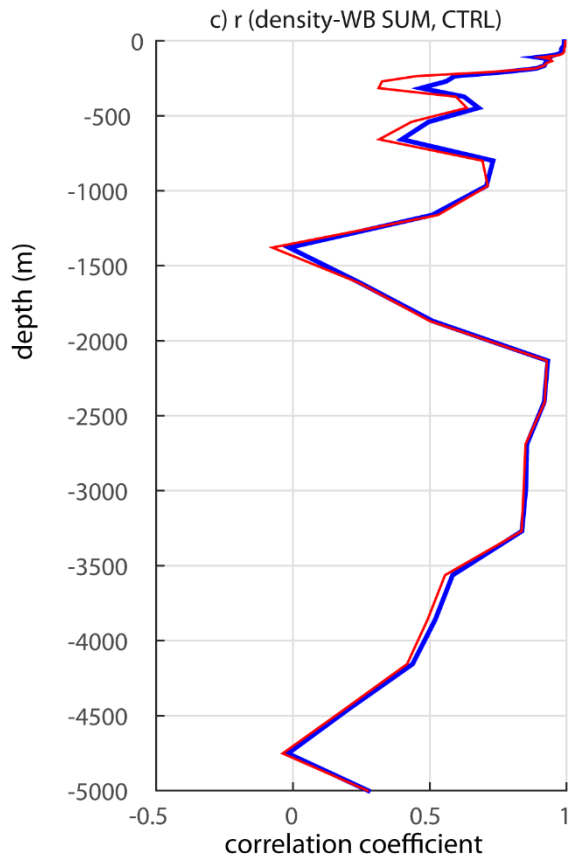
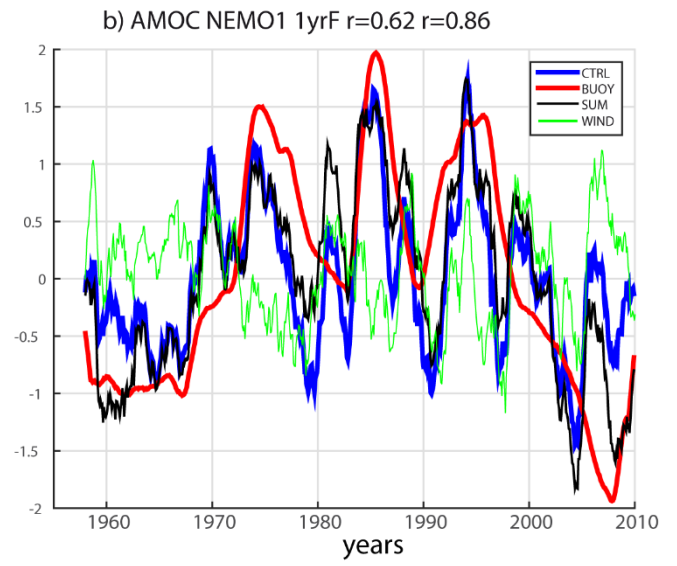
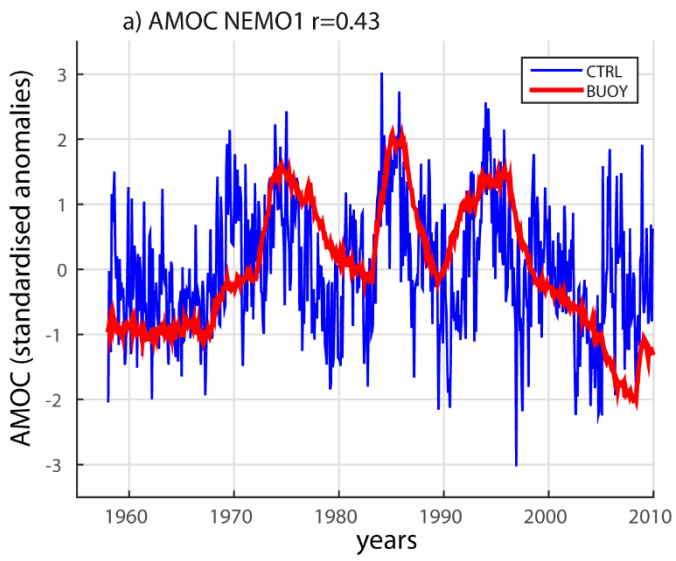
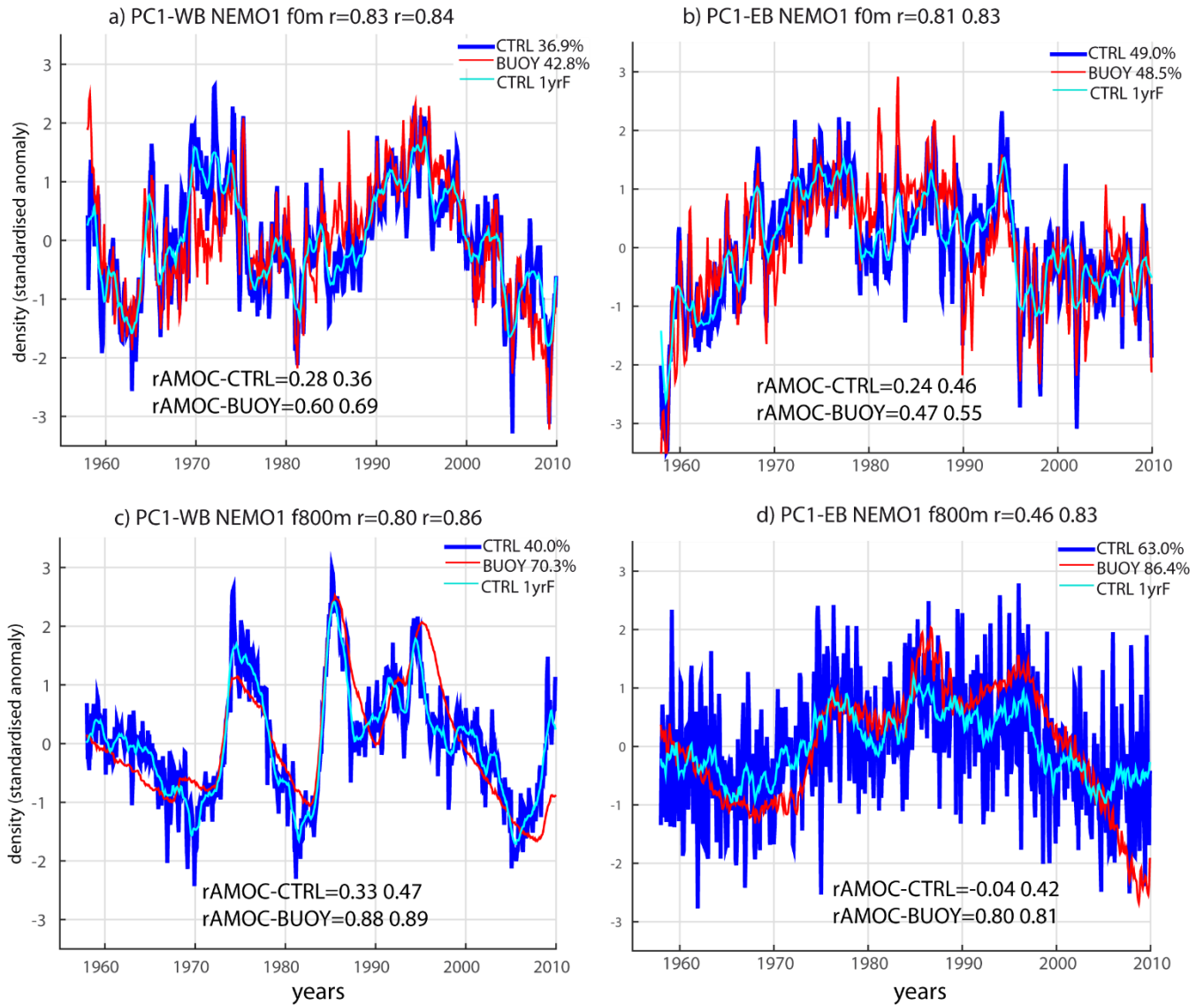


Figure 2. AMOC and linearity of the forcings for the density profiles at the boundaries. a) Geostrophic AMOC defined at 26N and 1100m for the CTRL (blue line) and BUOY (red line) experiments. In the title, correlation between both timeseries is detailed. b) Same as a) but the timeseries have been smoothed with one-year running mean filter. Black line refers to the AMOC for the SUM (BUOY+WIND). Green line is the filtered AMOC timeseries for the WIND experiment. In the title, correlations between blue-red and blue-black lines are detailed. c) Correlation coefficient between profiles of density anomalies at the WB for the CTRL experiment and the SUM. Blue line refers to anomalies without time filter, red line refers to anomalies that have been smoothed with one year running mean filter. d) Same as c) but for the EB.



810 **Figure 3 AMOC at 26N and individual EOF of density profiles at the 26N boundaries: sensitivity to depth truncation.** a) Timeseries
associated with the leading mode of density variability of the water column at the WB for the CTRL (blue line) and BUOY experiment
(red line), using the full water column. Cyan line corresponds to the timeseries of the CTRL smoothed with one-year running mean filter.
In the title, correlations between blue-red and blue-cyan lines are detailed. b) Same as a) but for the EB. c) -d) Same as a)-b) but
815 considering the EOF analysis with profiles from 800m downwards. Timeseries are dimensionless and the percentage of explained variance
is seen inset. Also on bottom, the correlations between PCs and the AMOC are detailed, for timeseries without filter and smoothed with
one year running mean filter.

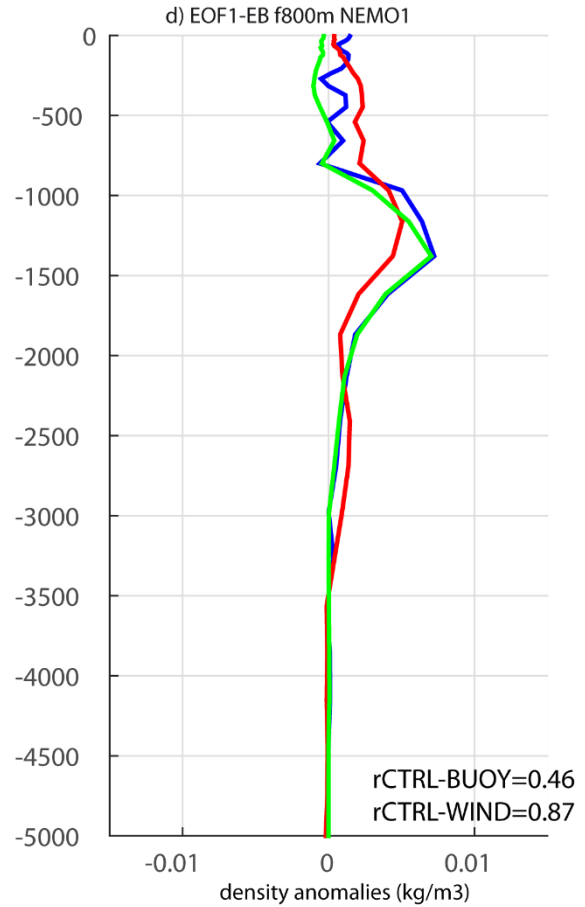
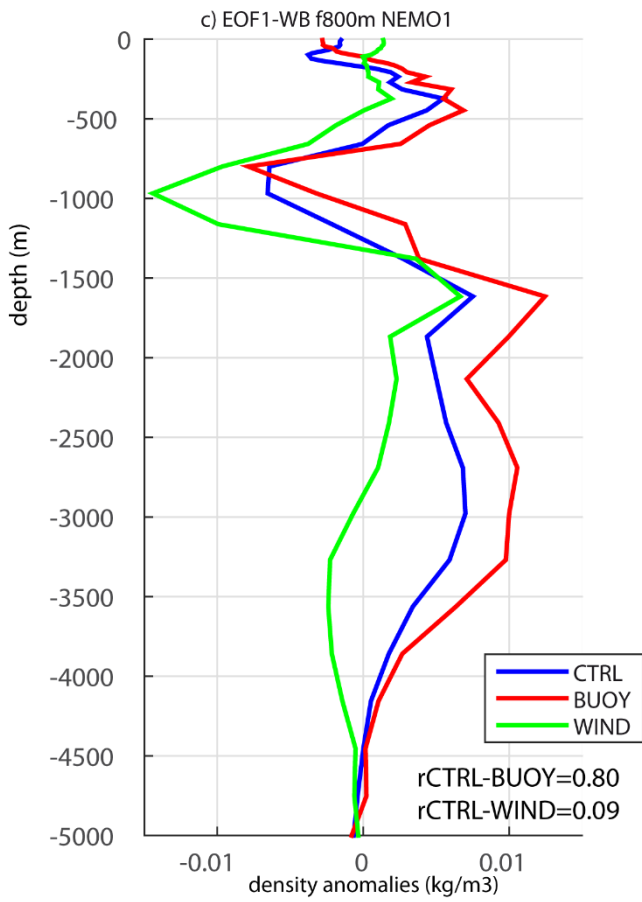
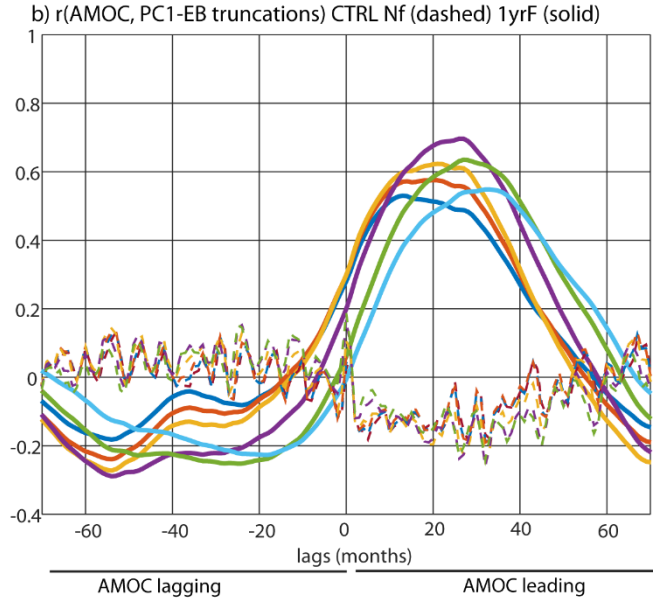
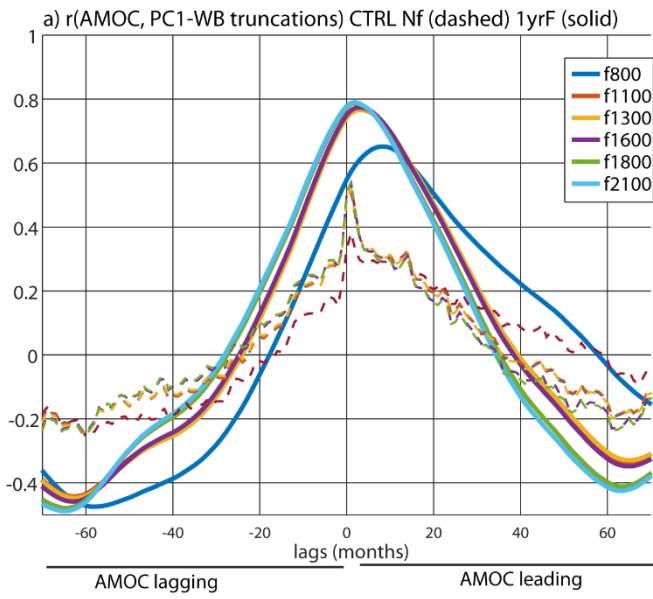


Figure 4. Individual EOF of density profiles at the 26N boundaries from 800m. a) Lead-lag correlations between AMOC and PC1-WB at different truncations (represented with different colours). Positive (negative) values over x-axis represents the AMOC leading (lagging) the PC1-WB. 1-year running mean filter (solid line) and no filter (dashed line) is applied to the timeseries for the CTRL experiment. b) Same as a) but for the correlations between AMOC and PC1-EB. c) Profile of density anomalies (in kg/m³) associated with the leading mode of density variability of the water column from 800m downwards at WB for the CTRL (blue line) and BUOY (red line) and WIND (green line) experiment. Associated timeseries are in Fig. 3c except for the WIND experiments. d) Same as c) but for the EB, associated timeseries are in Fig. 3d except for the WIND experiment. Correlation coefficient between timeseries of the PCs for CTRL with BUOY and WIND are detailed.

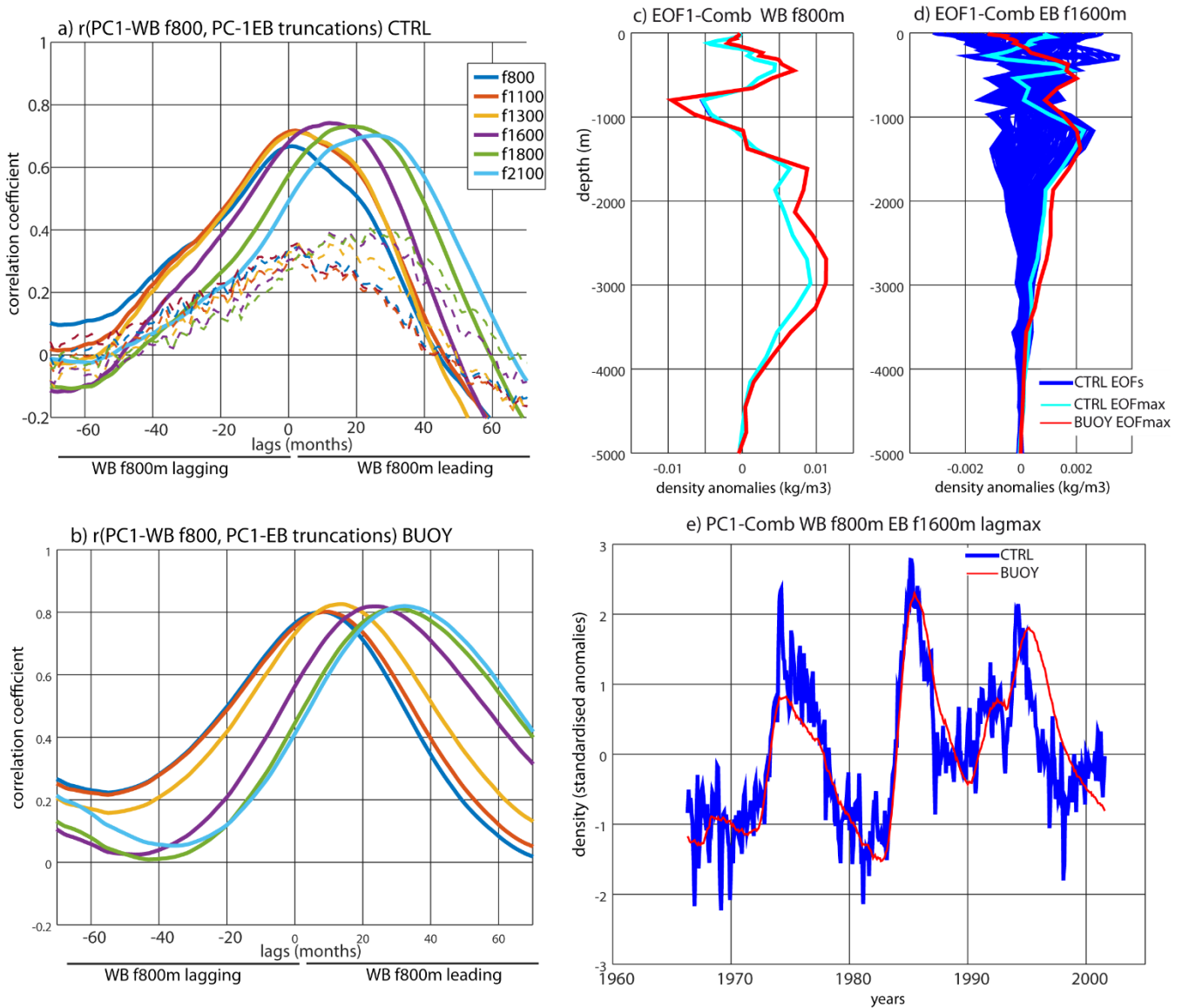
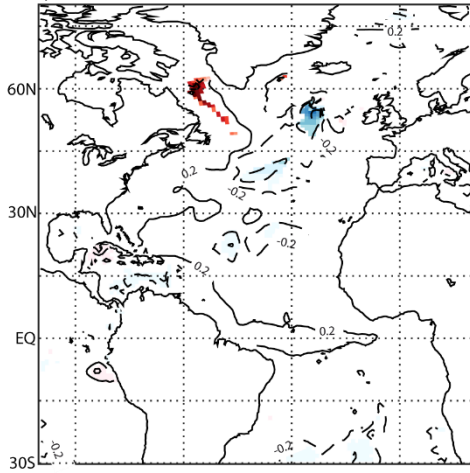
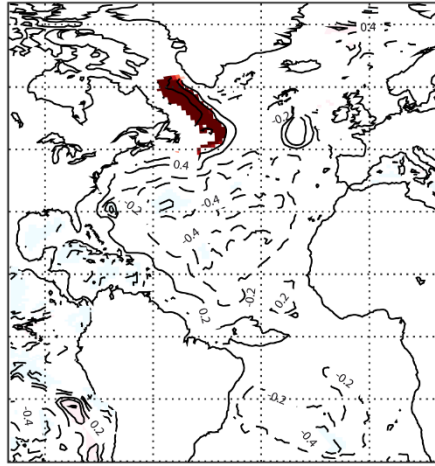


Figure 5. Relationship between Boundaries and the combined EOF. a) Lead-lag correlations between PC1-WB from 800m and PC1-EB at different truncations (represented with different colours). Positive (negative) values over x-axis represents the PC1-WB leading (lagging) the PC1-EB. 1-year running mean filter (solid line) and no filter (dashed line) is applied to the timeseries for the CTRL experiment. b) Same as a) but for the BUOY experiment and only 1-year running mean filter is plotted. c) - d) are the EOFs as the result of combining the density profile at the WB (from 800m) at lag 0 and the EB (from 1600m) at different lags. The CTRL experiment is plotted in blue and the cyan line corresponds to the EOF at the lag in which the explained fraction of variance is maximum (EOFmax). The latter equivalent EOFmax profile for the BUOY experiment is plotted in red. e) The associated timeseries of EOFmax for the CTRL (blue) and BUOY (red) experiment.

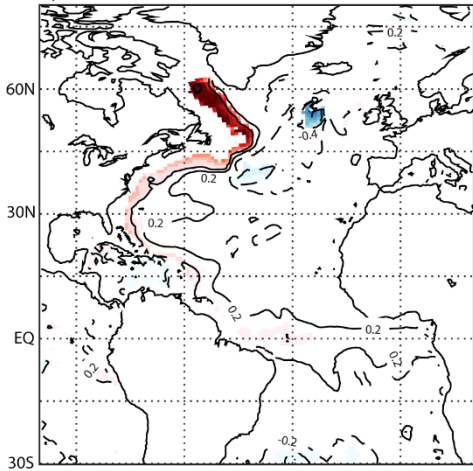
a) PC1-WB x den 2700-3000m CTRL -30 months



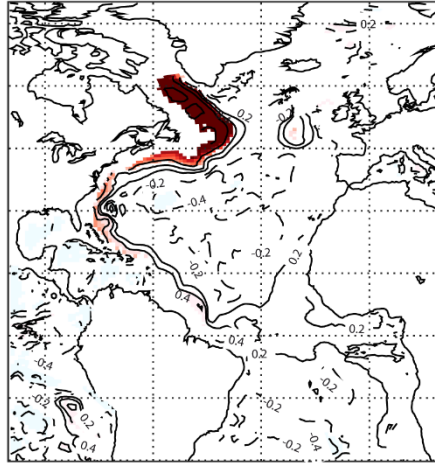
d) PC1-WB x den 2700-3000m BUOY -30 months



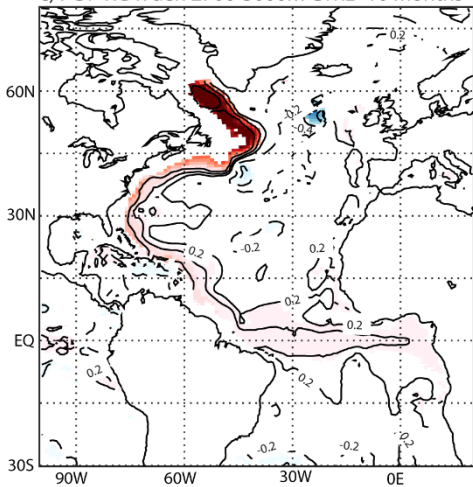
b) PC1-WB x den 2700-3000m CTRL -20 months



e) PC1WBxspa2700-3000m BUOY -20 months



c) PC1-WB x den 2700-3000m CTRL -10 months



f) PC1-WB x den 2700-3000m BUOY -10 months

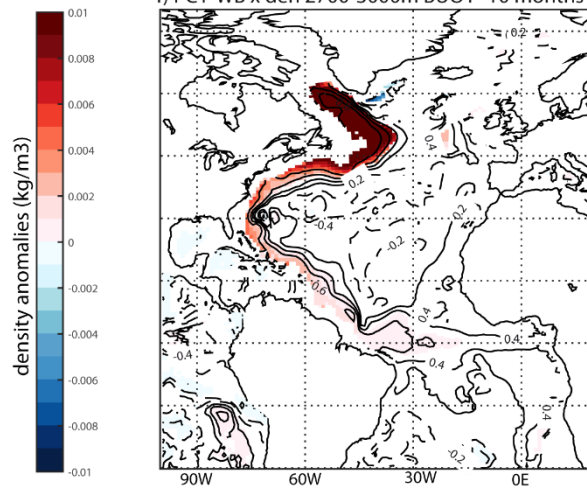
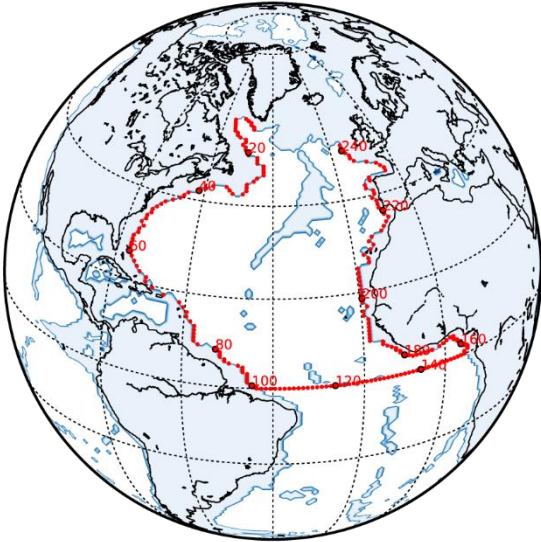
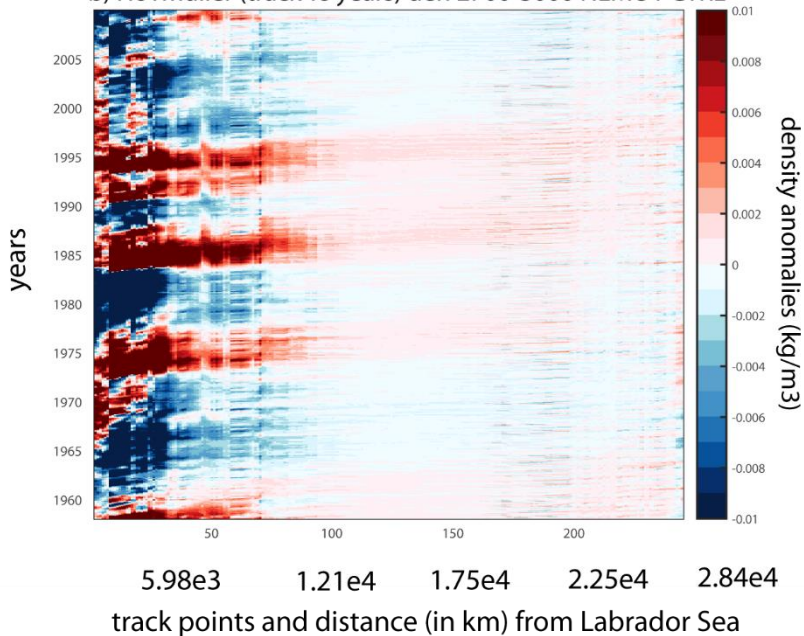


Figure 6. The spatial relationship of density anomalies with PC1-WB in the CTRL and BUOY experiment. a) Density anomalies (in kg/m³) averaged from 2700m to 3000m levels, 30 months in advance projected onto PC1-WB for the CTRL experiment. Black lines correspond to correlation contours every 0.2. Only significant areas are plotted with a Student's t-test at 95% confidence level considering only effective degrees of freedom. b) and c) are same as a) but for lags -20 and -10 months. d)-f) same as a)-c) but for the BUOY experiment.

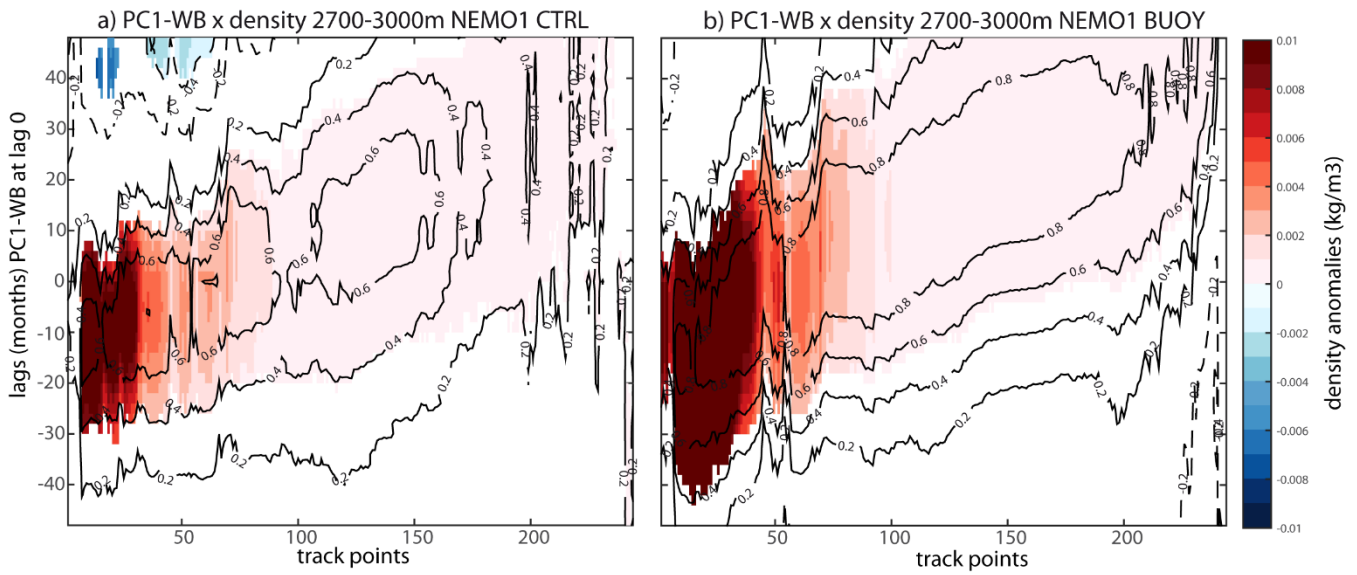
a) Track points



b) Hovmuller (track vs years) den 2700-3000 NEMO1 CTRL



845 **Figure 7. Hovmöller diagrams of density anomalies along the wave-track in the CTRL experiment.** a) Map of the wave-track used in the study. The track-points correspond to the first point before the coast following the bathymetry at 3269 m. b) Density anomalies (in kg/m³) along the track averaged from 2700m to 3000m levels for the CTRL experiment. Distance (in km) from Labrador Sea is detailed together with track-points on the x-axis.



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Figure 8. PC1-WB and density along the wave-track. a) Density anomalies (in kg/m³) averaged from 2700m to 3000m projected onto PC1-WB for the CTRL experiment. Black lines correspond to the correlation every 0.2. Only significant areas are plotted with a Student's t-test at 95% confidence level considering only effective degrees of freedom. b) Same as a) but for the BUOY experiment.

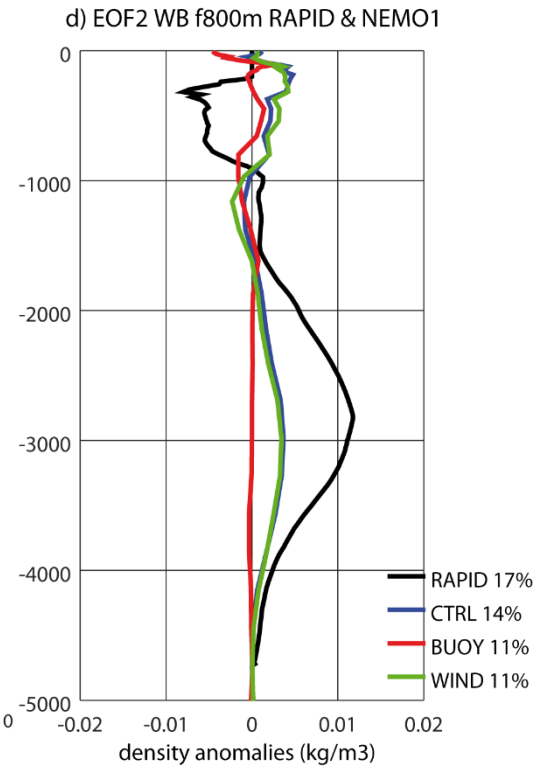
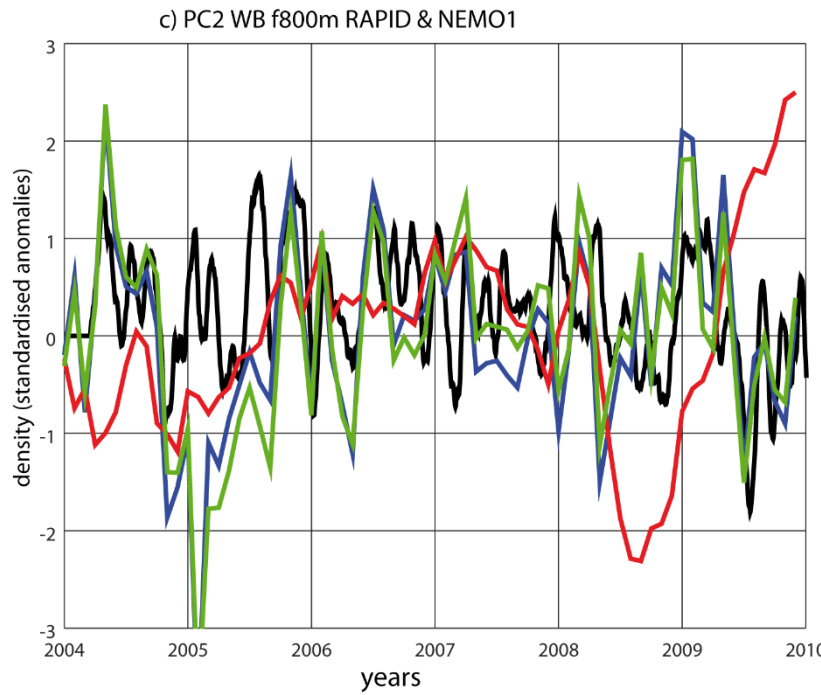
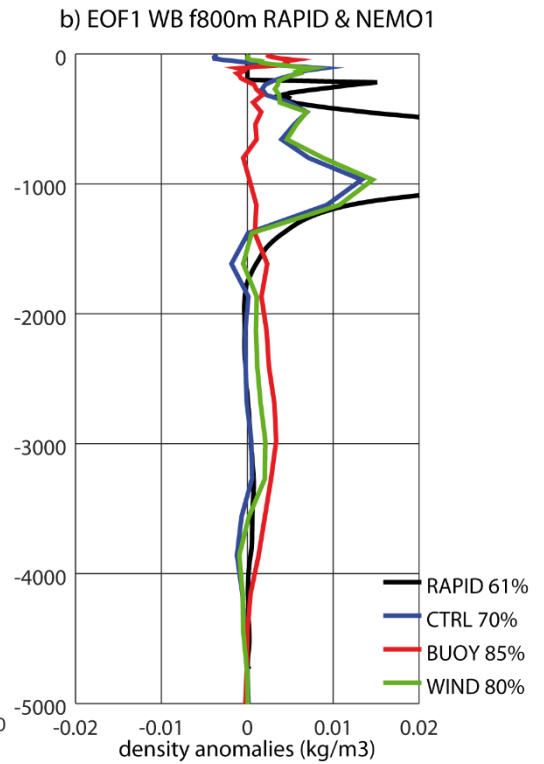
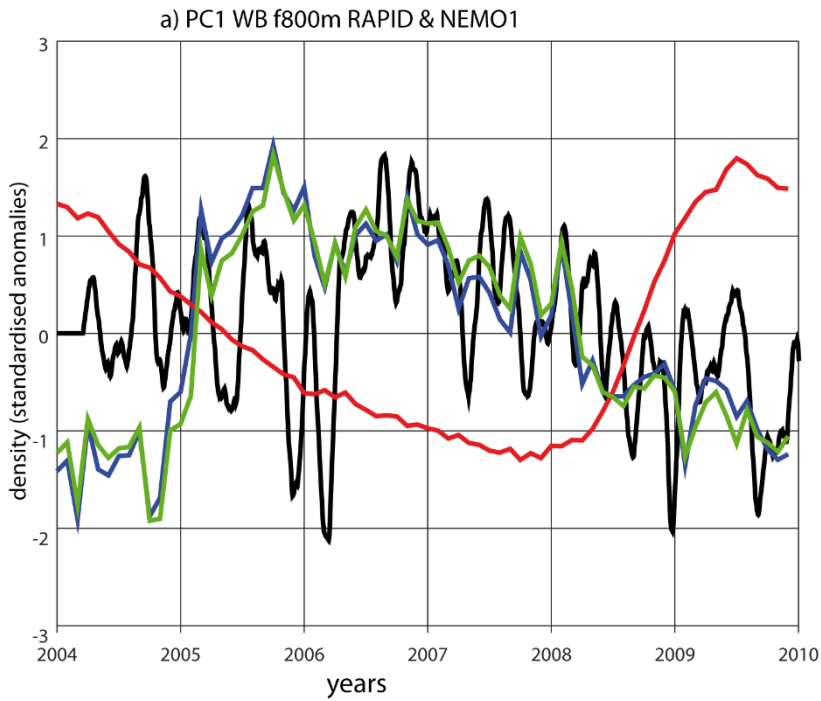
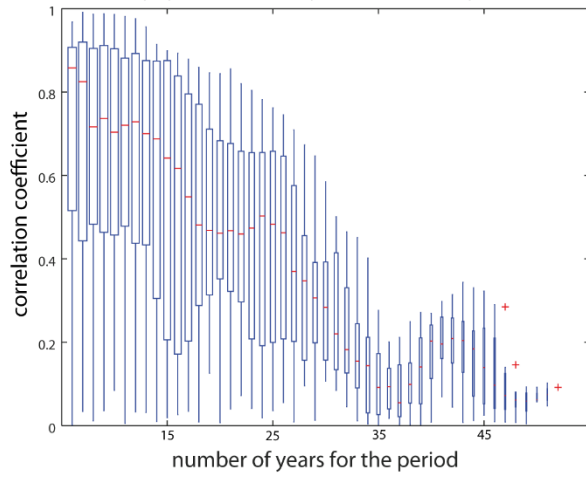


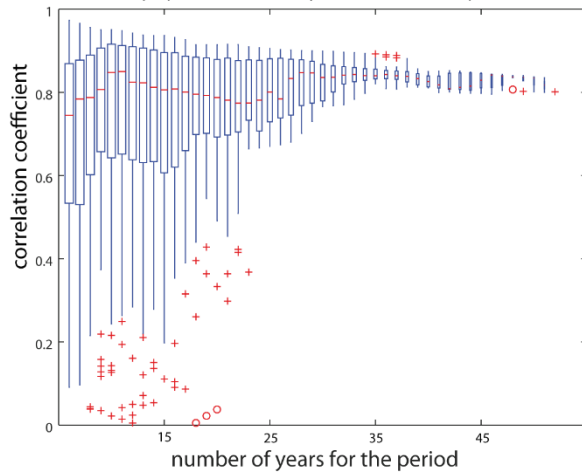
Figure 9. Leading EOFs for the NEMO experiments and the RAPID data. a) PCs timeseries associated with the leading EOF for the WB-26N, for the CTRL (blue line), BUOY (red line) and WIND (green line) experiments, and for RAPID (black line) in the common period 2004-2009. b) Density profiles of the EOF patterns linked to the PCs in a). c) -d) Same as a) -b) but for the second mode. For this time period, the EOFs for the CTRL experiment captures wind-forced only modes.

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a) $r(\text{PC1-WB CTRL}, \text{PC1-WB WIND})$



b) $r(\text{PC1-WB CTRL}, \text{PC1-WB BUOY})$



c) $r(\text{PC1-WB CTRL}, \text{AMOC} > 2\text{yr})$

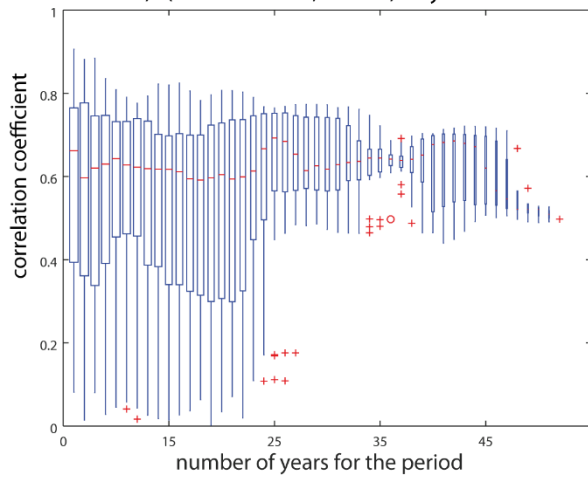
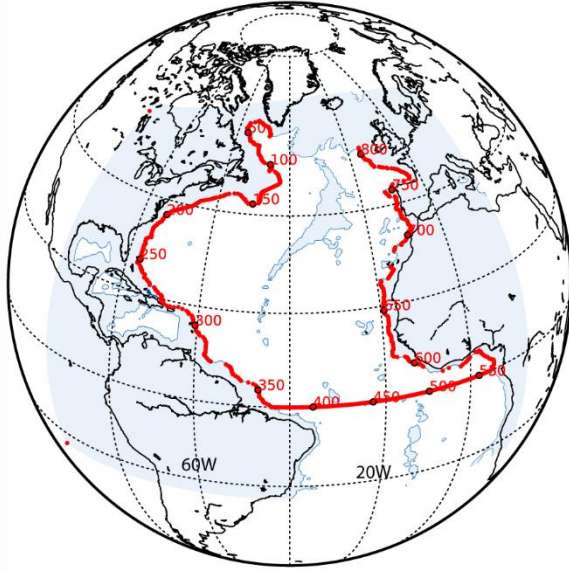


Figure 10. Impact of the number of years used in capturing the buoyancy forced signal in CTRL. a) box-plot of the correlation score (y-axis) between PC1-WB CTRL and the PC1-WB WIND for different number of years (x-axis) in the sub-periods considered. b) Same as a) but for the correlation between PC1-WB CTRL and PC1-WB BUOY. c) Same as a) but for the correlation between PC1-WB CTRL and AMOC filtered with periods >2 years. Red line corresponds to the median, box corresponds to the Inter-Quartile Range (IQR), the whiskers extend to the most extreme data point which is no more than 1.5 times the IQR from the box. Data points outside 1.5 and 3 times the IQR from the box are marked with crosses and circles respectively. The signal extracted by PC1-WB CTRL is always a mix of forcings however, there is a higher probability of extracting the buoyancy-forced signals when periods are longer than ~26 years. If periods are longer than 35 years, the wind-forced signal is negligible. We have to note that, for sub-periods with more years we have a smaller number of independent sub-periods and therefore there is less dispersion in the sample distribution (i.e. the box is smaller).

a) Track points GC2



b) Hovmuller (track vs years) density 3000m F>2yr GC2

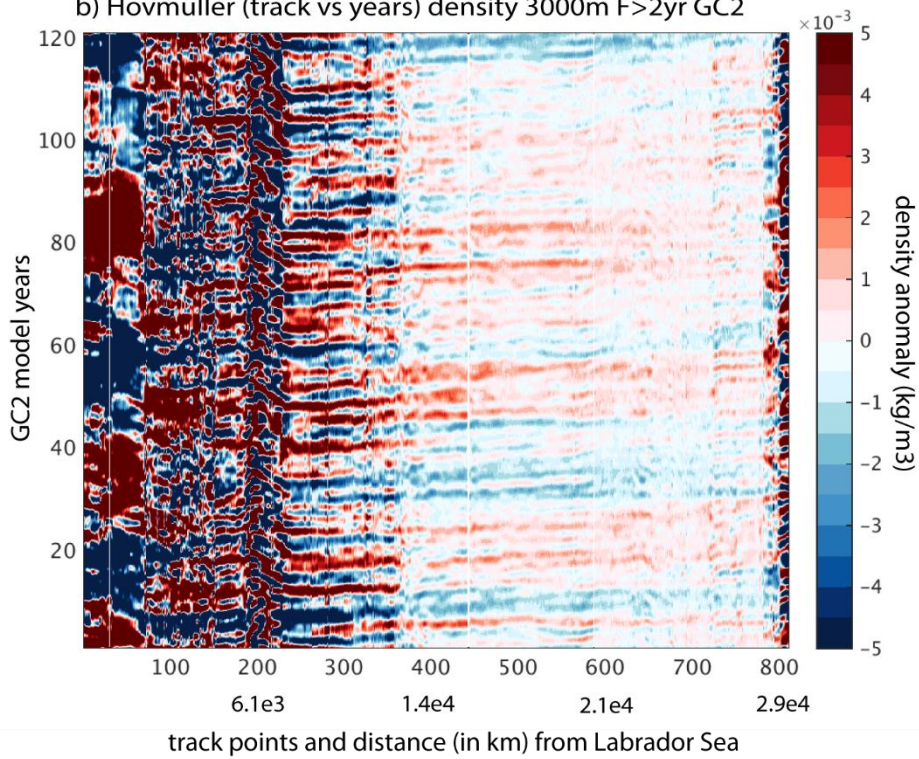


Figure 11. Hovmöller diagram of density along the wave track in the GC2 experiment. a) The track-points correspond to the first point before the coast following the bathymetry at 3138 m in GC2. b) Hovmöller diagram of density anomalies at 3000m along the wave track, after filtering the density timeseries retaining >2 years. x-axis corresponds to track points in Fig. 11a and y-axis corresponds to model year.

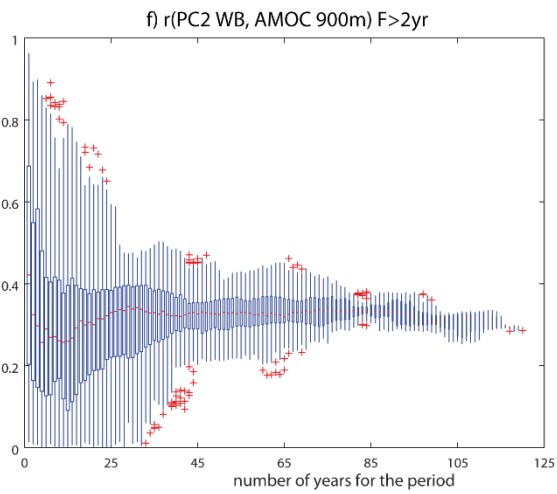
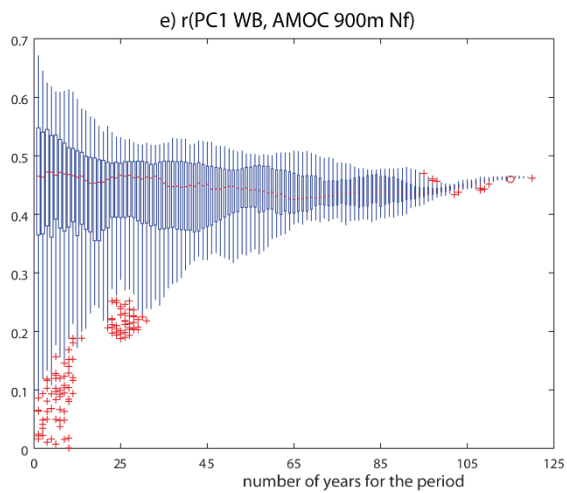
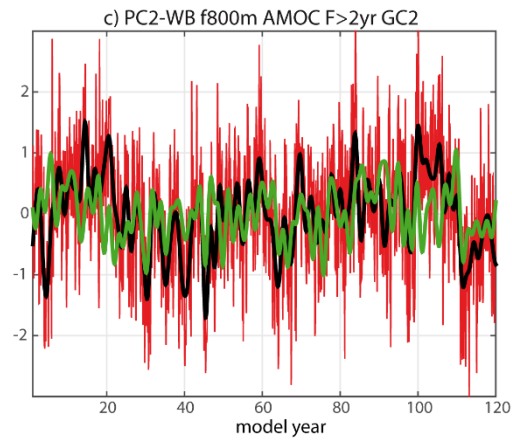
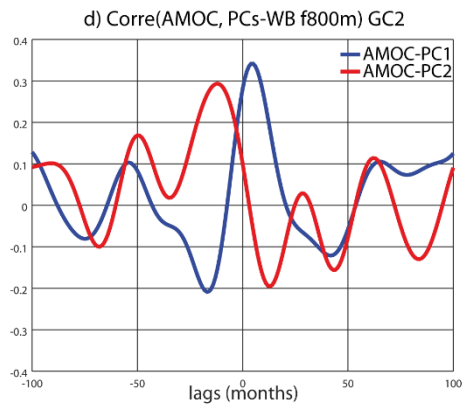
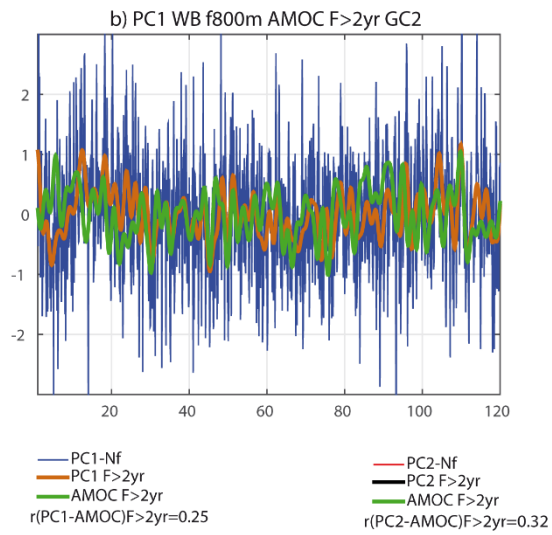
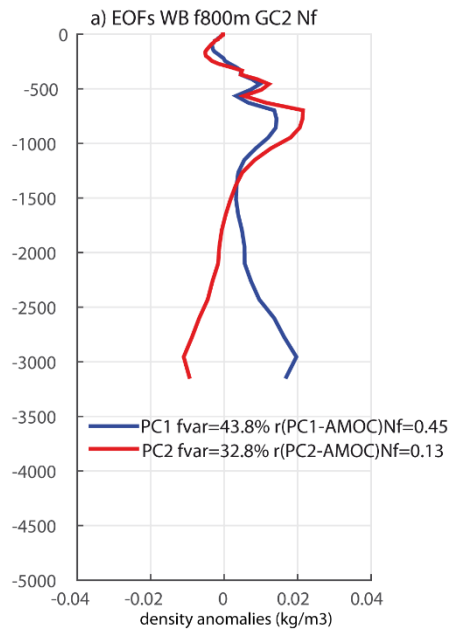


Figure 12. EOFs of WB density profiles at 26N in the GC2 experiment. a) Leading (blue line) and second (red line) EOF from WB at 26N and 76.75W in the GC2 experiment. Percentage of explained variance and correlation with the AMOC timeseries defined at 26N and 925m depth are displayed. b) Timeseries associated with leading EOF-WB (PC1-WB, blue line), PC1-WB filtered timeseries (PC1-F>2yr, brown line) and the AMOC filtered timeseries (AMOC-F>2yr, green line). c) Same as b) but for the PC2-WB. d) Correlation between AMOC (at lag 0) and the first two PCs-WB (lagged). e) Correlation between PC1-WB and AMOC timeseries. Box-plot similar to Fig. 10c, x-axis represents number of years used in the periods for the correlation, y-axis represents correlation score. f) Same as e) but for the correlation between PC2-WB and AMOC filtered with periodicity > 2 years.

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