1	Estimation of oceanic sub-surface mixing under a severe cyclonic storm
2	using a coupled atmosphere-ocean-wave model
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# 6 Abstract

A coupled atmosphere-ocean-wave model used to examine mixing in the upper oceanic 7 layers under the influence of a very severe cyclonic storm Phailin over the Bay of Bengal (BoB) 8 9 during 10-14 October 2013. Model simulations highlight prominent role of cyclone induced near-10 inertial oscillations in sub-surface mixing up to the thermocline depth. The inertial mixing introduced by the cyclone played <u>a</u>central role in deepening of <u>the</u> thermocline and mixed layer 11 depth by 40 m and 15 m, respectively. For the first time over the BoB, Aa detailed analysis of 12 inertial oscillation kinetic energy generation, propagation, and dissipation was carried out using 13 an atmosphere-ocean-wave coupled model during a cyclone. at a location in northwestern BoB. 14 The peak magnitude of kinetic energy in baroclinic and barotropic currents found to be 1.2 m<sup>2</sup> s<sup>-2</sup> 15 and 0.3×10<sup>-2</sup> m<sup>2</sup>-s<sup>-2</sup>, respectively. The power spectrum analysis suggested a dominant frequency 16 operative in sub-surface mixing was associated with near inertial oscillations. The peak strength 17 of 0.84 m<sup>2</sup> s<sup>-1</sup> in zonal baroclinic current found at 14 m depth. The baroclinic kinetic energy remain 18 higher (> 0.03 m<sup>2</sup> s<sup>-2</sup>)-during 11-12 October and decreased rapidly thereafter. The wave number 19 rotary spectra identified the downward propagation, from surface up to the thermocline, of energy 20 generated by inertial oscillations. A quantitative analysis of shear generated by the near-inertial 21 baroclinic current showed higher shear generation at 40 80 m depth during peak surface winds. 22 23 Analysis highlights that greater mixing within the mixed layer take place where the eddy kinetic diffusivity was high (>  $6 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}$ ). The turbulent kinetic energy dissipation rate increased from 24 4×10<sup>-14</sup> to 2.5×10<sup>-13</sup> W kg<sup>-1</sup> on approaching the thermocline that dampened mixing process further 25 downward into the thermocline layer. A quantitative estimate of kinetic energy in the oceanic water 26 27 column, its propagation and dissipation mechanisms were explained using the coupled 28 atmosphere-ocean-wave model. The large shear generated by the inertial oscillations found to overcome the stratification and initiate mixing at the base of the mixed layer. Greater mixing was 29 30 found at the depths where the eddy kinetic diffusivity was large. The baroclinic current, holding a

31	larger fraction of kinetic energy than the barotropic current, weakened rapidly after the passage of
32	cyclone. The shear-induced by inertial oscillations found to decrease rapidly with increasing depth
33	below the thermocline. The dampening of mixing process below the thermocline explained
34	through the enhanced dissipation rate of turbulent kinetic energy upon approaching the
35	thermocline layer. The wave-current interaction, non-linear wave-wave interaction were found to
36	affect the process of downward mixing and cause the dissipation of inertial oscillations.

#### 38

#### 39 1. Introduction

40 The Bay of Bengal (BoB), a semi-enclosed basin in the northeastern Indian Ocean, 41 consists of surplus near-surface fresh water due to large precipitation and runoff from the major river systems of the Indian subcontinent (Varkey et al., 1996; Rao and Sivakumar, 2003; Pant et 42 43 al., 2015). Presence of fresh water leads to salt-stratified upper ocean water column and formation of barrier layer (BL), a layer sandwiched between bottom of the mixed layer (ML) and top of the 44 thermocline, in the BoB (Lukas and Lindstrom, 1991; Vinayachandran et al., 2002; Thadathil et 45 46 al., 2007). The BL restricts entrainment of colder waters from thermocline region into the mixed 47 layer thereby, maintains warmer ML and sea surface temperature (SST). The warmer SST together with higher tropical cyclone heat potential (TCHP) makes the BoB as one of the active regions for 48 cyclogenesis (Suzana et al. 2007; Yanase et al. 2012, Vissa et al. 2013). Majority of tropical 49 cyclones generate during the pre-monsoon (April-May) and post-monsoon (October-November) 50 51 seasons (Alam et al., 2003; Longshore, 2008). The number of cyclones and their intensity is highly 52 variable in seasonal and interannual time scales. The stratification of the Ocean is one of the important factor to drive the ocean response of the tropical cyclone. The oceanic response to the 53 tropical cyclone depends on the stratification of the ocean. The BL formation in the BoB is 54 associated with the strong stratification due to the peak discharge from rivers in the post-monsoon 55 season. The intensity of the cyclone largely depends on the degree of stratification (Neetu et al. 56 57 2012; Li et al. 2013). The coupled atmosphere-ocean model found to improve the intensity of 58 cyclonic storm when compared to the uncoupled model over different oceanic regions (Warner et al., 2010; Zambon et al., 2014; Srinivas et al., 2016; Wu et al., 2016). Zambon et al. (2014) 59 compared the simulations from the coupled atmosphere-ocean and uncoupled models and reported 60

61	significant improvement in the intensity of storm in the coupled case as compared to the uncoupled
62	case. The uncoupled atmospheric model produced large ocean-atmosphere enthalpy fluxes and
63	stronger winds in the cyclone (Srinivas et al., 2016). When the atmospheric model WRF was
64	allowed interactions with the ocean model, the SST found to be more realistic as compared to
65	warm bias in the stand-alone WRF (Warner et al., 2010). Wu et al. (2016) demonstrated the
66	advantage of using a coupled model over the uncoupled model in better simulation of typhoor
67	Megi's intensity.

68 Mixing in the water column has an important role in energy and material transference. Mixing in the ocean can be introduced by the different agents such as wind, current, tide, eddy, 69 70 and cyclone. Mixing due to tropical cyclones is mostly limited to the upper ocean but the cyclone-71 induced internal waves can affect the subsurface mixing. Several studies have observed that the 72 mixing in the upper oceanic layer is introduced due to the generation of near--inertial oscillations 73 (NIO) during the passage of tropical cyclones (Gonella, 1971; Shay et al., 1989; Johanston et al., 74 2016). This mixing is responsible for deepening of ML and shoaling of the thermocline (Gill, 75 1984). The vertical mixing caused by storm--induced NIO has a significant impact on the upper 76 ocean variability (Price, 1981). The NIO are also found to be responsible for the decrement of SST 77 along the cyclone track (Chang and Anthes, 1979; Leipper, 1967; Shay et al., 1992; Shay et al., 78 2000). This decrease in SST is caused by the entrainment of cool subsurface thermocline water in 79 the mixed layer into the immediate overlying layer of water. This cooling of surface water is one of the component of the decaying mechanism of the stormy event reason for the decay of cyclone 80 81 (Cione and Uhlhorn, 2003). There is a remarkable difference in the The magnitude of surface this 82 cooling of surface temperature differs largely depending on the degree of stratification at the rightward to the cyclone track (Jacob, 2003; Price et al., 1981). moving on the highly stratified to 83 less or weakly stratified bay locations those are falling at the rightward to the cyclone track (Jacob, 84 2003; Price et al., 1981). 85

The near\_-inertial process can be analyzed from the baroclinic component of currents. The vertical shear of horizontal baroclinic velocities that is interrelated to buoyancy oscillations of surface layers-are\_is utilized in various studies to have an adequate understanding of the mixing associated with high frequency oscillations i.e. NIO (Zhang et al., 2014). The shear generated due to NIO is <u>an\_one of the-</u>important factor <del>other than the wind stress</del> for the intrusion of the cold 91 thermocline water into the ML during near--inertial scale mixing (Price et al., 1978; Shearman, 2005; Burchard and Rippeth, 2009). The alternative upwelling and downwelling features of the 92 93 temperature profile are an indication of the inertial mixing. The Kinetic energy bounded with these components of current shows a rise in magnitude at the right side of cyclone track (Price et al., 94 1981; Sanfoard et al., 1987; Jacob, 2003). The reason for this high magnitude of kinetic energy is 95 linked with strong wind and rotating wind vector condition of the storm. The spatial distribution 96 of near-inertial energy is primarily controlled by the boundary effect for inertial oscillations (Chen 97 et al., 2017). The NIO energy found to decline with the decreasing depth and vanish in the coastal 98 99 regions (Schahinger, 1988; Chen et al., 2017).

In several studies (Chang et al., 2008; Lin et al., 2008; Shang et al., 2008; Lin et al., 2003; 100 101 Zhao et al., 2009), upper Ocean response for various cyclonic events is also inspected and proved 102 for the enhancement of primary productivity during post cyclone state of the Ocean. At the time 103 when storm is active and prior to it, the surface concentration of chlorophyll a (Chl-a), a proxy for 104 the concentration of primary productivity is comparatively lower than that of the post cyclonic state of ocean surface (Sarangi, 2011, Latha et al., 2015). This increment in the chlorophyll is 105 106 dependent on the relative entrainment of the cool subsurface water, enriched in nutrient under the influence of energetic near inertial wave mixing caused by the tropical cyclones. 107

108 The Aimaim of this paper is to understand and quantify the near-inertial mixing due to the 109 very severe cyclonic storm Phailin in the BoB., a very severe cyclonic storm (VSCS) in the BoB. 110 Phailin was developed over the BoB in the northern Indian ocean in October 2013. The landfall of Phailin occurred on 12 October 2013 around 15:30 GMT near Gopalpur district of Odisha state on 111 112 the east coast of India. After the 1999 super cyclonic event of the Odisha coast, Phailin was the 113 second strongest cyclonic event that made landfall on the east coast of India (Kumar and Nair, 114 2015). The low-pressure system developed in the north of the Andaman Sea on 7th October 2013, which transformed into a depression on 8th October at 12 °N, 96 °E. This depression got converted 115 116 to a cyclonic disturbance on 9th October and further intensified while moved to east-central BoB and opted the maximum wind speed of 200 km h<sup>-1</sup> at 03:00 GMT on 11<sup>th</sup> October. Finally, landfall 117 occurs at 17:00 GMT 12th October. More details on the development and propagation of Phailin 118 can be found in the literature (IMD Report, 2013; Mandal et al. 2015). The performance of the 119

120 coupled atmosphere-ocean model in simulating the oceanic parameters temperature, salinity, and 121 currents during the Phailin is accessed in Prakash and Pant (2017). 122 Most of the past studies on the oceanic mixing under cyclonic conditions were carried out 123 using in-situ measurements, which are constrained by the spatial and temporal availability. To the 124 best of our knowledge, the present study is first of its kind that utilizes a coupled atmosphere-125 ocean-wave model over the BoB to estimate the cyclone-induced mixing and associated energy propagation at the cyclone track and a location of maximum surface wind stress during the period 126 of peak intensity of the cyclone. The study also focuses on analyzing the subsurface distribution 127 of NIO with its vertical mixing potential. Further, the study quantifies the shear generated mixing 128 129 and the kinetic energy of these baroclinic mode of horizontal current varying in the vertical section 130 at a selected location during the active period of the cyclone. The dissipation rate of NIO and 131 turbulent eddy diffusivity are quantified.

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# 133 2. Data and Methodology

### 134 2.1 Model details

135 Numerical simulations during the period of <del>VSCS</del> Phailin were carried out using the 136 Ocean-Atmosphere-wave-Sediment ocean-atmosphere-wave-sediment transport coupled (COAWST), described in detail by Warner et al. (2010). COAWST modeling system couples the 137 138 three-dimensional oceanic model 'Regional Ocean Modeling System' (ROMS), the atmospheric 139 model 'Weather Research and Forecasting' (WRF), and the wind wave generation and propagation 140 model 'Simulating Waves Nearshore' (SWAN). ROMS model used for the study is a free surface, 141 primitive equation, sigma coordinate model. ROMS is a hydrostatic ocean model that solves finite 142 difference approximations of the Reynolds averaged Navier-Sstokes equations (eChassignet et al., 143 2000; Haidvogel et al., 2000, Haidvogel et al., 2008; Shchepetkin and McWilliams, 2005). The atmospheric model component in the COAWST is a non-hydrostatic, compressible model 144 145 'Advanced Research Weather Research Forecast Model' (WRF-ARW), described in Skamarock 146 et al., (2005). It has different schemes for representation of boundary layer physics and physical parameterizations of sub-grid scale processes. In the COAWST modeling system, appropriate 147 148 modifications were made in the code of atmospheric model component to provide an improved

bottom roughness from the calculation of the bottom stress over the ocean (Warner et al., 2010).
Further, the momentum equation is modified to improve the representation of surface waves. The
modified equation needs the additional information of wave energy dissipation, propagation
direction, wave height, wave-length that are obtained from wave component of the COAWST
model.

154 The spectral wave model SWAN, used in the COAWST modeling system, is designed forshallow water. The wave action balance equation is solved in the wave model for both spatial and 155 156 spectral spaces (Booij et al. 1999). In the COAWST, the ocean model ROMS simulated free 157 surface elevations (ELV), and current (CUR) are provided to the wave model SWAN. The Kirby 158 and Chen (1998) formulation has been used for the computation of current. The Model Coupling 159 Toolkit (MCT) used as a coupler in the COAWST modeling system to couple different model 160 components (Larson et al., 2004; Jacob et al., 2005). A parallel-coupled approach is utilized by the coupler that permits the transmission and transformation of various distributed parameters between 161 162 component models. MCT coupler facilitates exchanges of prognostic variables from one model to another model components. Further details on various parameters exchanged among the 163 164 component models of COAWST modeling system can be found in Warner et al. (2010). The SWAN 165 model used in the COAWST system includes the wave-wind generation, wave-breaking, wave-dissipation, 166 and nonlinear wave-current-wind interaction. The 'Model Coupling Toolkit' (MCT) used as a coupler in 167 the COAWST modeling system to couple different model components (Larson et al., 2004; Jacob et al., 168 2005). The coupler utilizes a parallel-coupled approach to facilitate the transmission and transformation of 169 various distributed parameters among component models. MCT coupler exchanges prognostic variables 170 from one model to another model component as shown in Figure 1. The WRF model receives sea surface 171 temperature (SST) from the ROMS model and supplies the zonal (Uwind) and meridional (Vwind) 172 components of 10-m wind, atmospheric pressure (Patm), relative humidity (RH), cloud fraction (Cloud), 173 precipitation (Rain), shortwave (Swrad) and longwave (Lwrad) radiation to the ROMS model. The SWAN 174 model receives Uwind and Vwind from the WRF model and transfers significant wave height (Hwave) and 175 mean wavelength (Lmwave) to the WRF model. A large number of variables are exchanged between 176 ROMS and SWAN models. The ocean surface current components (Us, Vs), free surface elevations  $(\eta)$ , 177 and bathymetry (Bath) provided to the SWAN from ROMS model. The wave parameters i.e. Hwave, 178 Lmwave, peak wavelength (Lpwave), wave direction (Dwave), surface wave period (Tpsurf), bottom wave 179 period (Tmbott), percentage wave breaking (Qb), wave energy dissipation (DISSwcap), and bottom orbit

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velocity (Ubot) provided from the SWAN to ROMS model through the MCT coupler. Further details on
 the COAWST modeling system can be found in Warner et al. (2010).

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# 183 2.2 Model configuration and experiment design

The coupled model was configured over the BoB to study the VSCS-Phailin during the 184 period of 10 to 15 October 2013 00 GMT 10 October - 00 GMT 15 October 2013. The setup of 185 COAWST modeling system used in this study included fully coupled atmosphere-ocean-wave 186 187 (ROMS+WRF+SWAN) models but the sediment transport is not included. A non-hydrostatic, 188 fully compressible atmospheric model with a terrain-following vertical coordinate system, WRF-189 ARW (version 3.7.1) was used in the COAWST configuration. The WRF model used with 9 km horizontal grid resolution over the domain 65 °E-105 °E, 1°N-34 °N and 30 sigma levels in the vertical. 190 191 The WRF was initialized with 'National Centre for Environmental Prediction' (NCEP) 'Final Analysis' 192 (FNL) data (NCEPFNL, 2000) at 00 GMT 10 October 2013. The lateral boundary conditions in WRF were 193 provided at 6 hour interval from the FNL data. The atmospheric model-We used the parameterization schemes for calculating boundary layer processes, precipitation processes, and surface radiation 194 195 fluxes. The Monin-Obukhov scheme of surface roughness layer parameterization (Monin and Obukhov 1954) was activated in the model. The Rapid Radiation Transfer Model (RRTM) and 196 cloud-interactive shortwave (SW) radiation scheme from Dudhia (1989) were used. The planetary 197 198 boundary layer scheme YSU-PBL, described by Noh et al. (2003), was used. At each time step, the calculated value of exchange coefficients and surface fluxes off the land or ocean surface by 199 the atmospheric and land surface layer models (NOAH) passed to the YSU PBL. The Grid-scale 200 precipitation processes were represented by WRF single-moment (WSM) six-class moisture 201 microphysics scheme by Hong and Lim (2006). The sub-grid scale convection and cloud 202 203 detrainment were taken care by Kain (2004) cumulus scheme.

204 <u>A terrain following ocean model ROMS with 40 sigma levels in the vertical used in this</u> 205 <u>study. The ROMS model domain used with zonal and meridional grid resolutions of 6 km and 4 km</u>, 206 <u>respectively. This high resolution in ROMS enables to resolve mesoscale eddies in the ocean. The vertical</u> 207 <u>starching parameters i.e.  $\theta_{s}$  and  $\theta_{p}$  were set at 7 and 2, respectively. The northern lateral boundary in ROMS</u> 208 <u>was closed by the Indian subcontinent. The ROMS model observed open lateral boundaries in the west</u>, 209 <u>east, and south in the present configuration. The initial and lateral open boundary conditions were derived</u>

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210 from the 'Estimating the Circulation and Climate of the Ocean, Phase II' (ECCO2) data (Menemenlis et 211 al., 2005). The ocean bathymetry was provided from the 2-minute gridded global relief (ETOPO2) data 212 (National Geophysical Data Center, 2006). There was no relaxation provided to the model for any correction in the temperature, salinity, and current fields. A terrain following ocean model ROMS with 213 214 40 sigma levels used in this study. The Generic-Length-Scale (GLS) vertical mixing scheme parameterized as the K- $\varepsilon$  model used (Warner et al., 2005). Tidal boundary conditions were 215 derived from the TPXO.7.2 (ftp://ftp.oce.orst.edu/dist/tides/Global) data, which includes phase 216 and amplitude of the M2, S2, N2, K2, K1, O1, P1, MF, MM, M4, MS4, and MN4 tidal constituents 217 along the east coast of India. The tidal input was interpolated from TPXO.7.2 grid to ROMS 218 computational grid. The Shchepetkin boundary condition (Shchepetkin, 2005) for the barotropic 219 220 current was used at open lateral boundaries of the domain which allowed the free propagation of 221 astronomical tide and wind--generated currents. The domains of atmosphere and ocean models 222 which were part of the COAWST modeling system are shown in Figure 24. The domain for SWAN 223 model was similar to the domain of ROMS model The ROMS and SWAN were configured over 224 the common model domain shown with the shaded bathymetry data in Figure 2. The two locations used for the time series analysis are marked with stars in Figure 2. These two locations, one on-225 226 track and another off-track, were selected in the vicinity of the region of maximum surface cooling and wind-stress during the passage of Phailin. The wave model SWAN was forced with the WRF 227 228 computed wind field. We used 24 frequency (0.04 - 1.0 Hz) and 36 directional bands in SWAN model. The 229 boundary conditions for SWAN were derived from the 'WaveWatch III' model. The atmospheric model WRF had 9 km horizontal grid resolution over the domain 65 °E 105 °E, 1°N-34 °N with 30 sigma 230 levels in vertical. WRF was initialized with National Centre for Environmental Prediction (NCEP) 231 232 Final Analysis (FNL) data (NCEP FNL, 2000) on 10<sup>th</sup> October 2013 at 00 GMT. Lateral boundary conditions in WRF provided at 6 h interval from the FNL data. The ROMS model domain had 233 234 zonal and meridional grid resolutions of 6 km and 4 km, respectively. The northern lateral boundary in ROMS was closed and the model observed open boundaries in rest of the sides. The 235 oceanic initial and lateral open boundary conditions were derived from the Estimating the 236 237 Circulation and Climate of the Ocean, Phase II (ECCO2) data. Ocean bathymetry was derived 238 from 2-minute gridded global relief data (ETOPO2). In the COAWST system, the free surface 239 elevations (ELV) and current (CUR) simulated by ocean model ROMS are provided to the wave model 240 SWAN. The Kirby and Chen (1998) formulation was used for the computation of currents. The surface 241 wind applied to the SWAN model (provided by WRF) used in the Komen et al. (1984) closure model to

transfer energy from the wind to the wave field. The baroclinic time step used in ROMS model was 5 s. 242 243 The SWAN and WRF models used with time steps of 120 s and 60 s, respectively. The coupled modeling 244 system allows the exchange of prognostic variables among the atmosphere, ocean, and wave models at every 600 s. The SST simulation at high spatial and temporal resolutions enables 245 246 accurate heat fluxes at the air-sea interface and exchange of heat between oceanic mixed layer and 247 atmospheric boundary layer. The surface roughness parameter calculated in the WRF model based on Taylor and Yelland (2001), which involved parameters from the wave model. The Advanced 248 249 Very High Resolution Radiometer (AVHRR) data was used for the validation of model simulated SST. 250

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#### 252 2.3. Methodology

The baroclinic current component was calculated by subtracting the barotropic component 253 254 from the mean current with a resolution of 2 m in the vertical. The power spectrum analysis was performed on the zonal and meridional baroclinic currents along the depth section of the selected 255 locations by using periodogram method (Auger and Flandrin, 1995). The continuous wavelet 256 257 transform using Morlet wavelet method (Lilly and Olhede, 2012) carried out to analyze the 258 temporal variability of the baroclinic current at a particular level of 14 m. The near-inertial baroclinic velocities were filtered by the Butterworth 2nd order scheme for the cutoff frequency 259 range of 0.028 to 0.038 cycle hr<sup>-1</sup>, of 0.033 to 0.043. The filtered zonal  $(u_f)$  and meridional  $(v_f)$ 260 261 inertial baroclinic currents were used to calculate the inertial baroclinic kinetic energy ( $E_f$ ) in m<sup>2</sup> 262  $s^{-2}$  and inertial shear (S<sub>f</sub>) following Zhang et al. (2014) using equation (1).

$$S_f^2 = \left(\frac{\partial u_f}{\partial z}\right)^2 + \left(\frac{\partial v_f}{\partial z}\right)^2 \tag{1}$$

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 $S_{f} = \left(\frac{\partial u_{f}}{\partial z}\right)^{2} + \left(\frac{\partial v_{f}}{\partial z}\right)^{2} \tag{1}$ 

As the stratification is a measure of oceanic stability, the buoyancy frequency (N) was calculatedusing equation (2)

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268	$N^2 = -\frac{g}{\rho} \frac{\partial \rho}{\partial z}$	(2)	
269	Where $\rho$ is <u>the</u> density of sea-water and g is <u>the</u> accelerat	ion due to gravity.	
270	The analysis of generation of the inertial oscillat	ions and their dissipation was performed	
271	on the basis of turbulent dissipation rate ( $\epsilon$ ) and turbuler	tt eddy diffusivity ( $k_{\rho}$ ). These parameters	
272	were calculated by using following formula (Mackinn	on and Gregg, 2005; van der Lee and	
273	Umlauf, 2011; Palmer et al., 2008; Osborn, 1980)		
274	$\varepsilon = \varepsilon_0 \left(\frac{N}{N_0}\right) \left(\frac{S_{lf}}{S_0}\right)$	(3)	
275	$k_ ho=0.2x\left(rac{arepsilon}{N^2} ight)$	(4)	
276	Where S <sub>lf</sub> is <u>the</u> low shear background velocity, Values	of $N_0 = S_0 = 3$ cycle per hour and $\varepsilon_0 =$	
277	$10^{-8} \text{ W kg}^{-1}$ .		
278			
279	3. Results and Discussion		
280	3.1. Details of VSCS Phailin		
281	Phailin, a very severe cyclonic storm (VSCS) w	vas developed over the BoB in northern	
281 282	Phailin, a very severe cyclonic storm (VSCS) w Indian Ocean in October 2013. The landfall of Phailin oc	vas developed over the BoB in northern curred on 12 October 2013 around 15:30	
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281 282 283 284 285 286 287 288 289 289 290	Phailin, a very severe cyclonic storm (VSCS) w Indian Ocean in October 2013. The landfall of Phailin oc GMT near Gopalpur district of Odisha state at the east cor- event of the Odisha coast, Phailin was the second strong east coast of India (Kumar and Nair, 2015). The low press Andaman Sea on 7 <sup>th</sup> October 2013, which transformed in 96 °E. This depression got converted to a cyclonic of intensified while moved to east central BoB and opted th 03:00 GMT on 11 <sup>th</sup> October. Finally, landfall occur at 11 the development and propagation of VSCS Phailin can be	vas developed over the BoB in northern curred on 12 October 2013 around 15:30 ast of India. After the 1999 super cyclonic gest cyclonic event that made landfall at sure system developed in the north of the ato a depression on 8 <sup>th</sup> October at 12 °N, disturbance on 9 <sup>th</sup> October and further are maximum wind speed of 200 km h <sup>-1</sup> at 7:00 GMT 12 <sup>th</sup> October. More details on e found in literature (IMD Report, 2013;	

292 <u>3.1.</u>3.2 Validation of coupled model simulations

293 The WRF model simulated track of Phailin was validated against the India Meteorological 294 Department (IMD) reported best-track of the cyclone. A comparison of model simulated track with 295 the IMD track is shown in Figure 32. Solid circles marked on both the tracks represent the 3-hourly positions of the cyclone's center, as identified by the minimum surface pressure. The daily 296 positions of the centre of Phailin are labelled with the date. WRF model in the coupled 297 configuration does a fairly good job in simulating the track, translational speed, and landfall 298 299 location of Phailin. The positional track error was about 40 km when compared to IMD track of 300 Phailin. The stand-alone WRF model (not shown here) was found to simulate Phailin track almost 301 similar to the WRF in coupled configuration. However, the intensity (surface wind speed) in WRF 302 stand-alone model was higher as compared to the coupled model. Figure 4 shows the comparison 303 of stand-alone WRF and coupled model simulated mean sea level pressure (MSLP), wind speed, 304 and wind direction at a buoy (BD09) location (marked with a blue circle in Figure 3). It can be 305 inferred from the figure that stand-alone WRF simulated larger pressure drop and higher wind 306 speed as compared to buoy measurements. The WRF in coupled model configuration shows better 307 performance in simulating the surface wind speed and pressure during Phailin. The exchange of wave parameters with the WRF model in coupled configuration provides realistic sea surface 308 309 roughness that resulted in improvement of surface wind speed.

310 The SST simulated by the ROMS model in coupled and stand-alone configurations 311 simulated SST was was validated against the Advanced Very High Resolution Radiometer 312 (AVHRR) satellite data on each day for the period of Phailin passage over the BoB. Figure 53 313 shows that the coupled model captures the SST spatial pattern reasonably well with about -0.5°C bias in northwestern BoB on 13-14 October. This order of bias in SST could be resulted from the 314 315 errors in initial and boundary conditions provided to the model. simulations are capturing the SST features as well as the magnitude of cooling associated with the storm. The maximum cooling of 316 the sea surface observed on 13<sup>th</sup> October in the northwestern BoB in both, <u>coupled</u> model and 317 318 observations. This post-cyclone cooling primarily associated with the cyclone-induced upwelling resulting from the surface divergence driven by the Ekman transport. Thus, the coupled model is 319 320 reproducing dynamical processes and vertical velocities reasonably well. The stand-alone ROMS 321 model forced with the WRF winds in un-coupled mode overestimates the cyclone-induced cooling with -2.2 °C bias in SST on 13-14 October (Figure 5). The stronger surface winds in stand-alone 322 323 WRF cause the larger cold bias in stand-alone ROMS model.

#### 324 <u>3.2.</u>3.3. Cyclone\_-induced mixing

The coupled atmosphere-ocean-wave simulation is an ideal tool to understand air-sea 325 326 exchange of fluxes and their effects on the oceanic water column. Surface wind sets up currents 327 on the surface as well as initiate mixing in the interior of the upper ocean. In order to examine the 328 strength of mixing due to VSCS Phailin, the model simulated vertical temperature profile together 329 with the surface wind speed, zonal and meridional components of windcurrent, and kinetic energy 330 at the on-track and off-track locations a location 18.75 °N, 86.66 °E are plotted in Figure 64. Comparatively stronger zonal and meridional currents observed at the off-track location than the 331 on-track location on 12 October. The larger kinetic energy available at the off-track location leads 332 333 to greater mixing resulting into deeper mixed layer on 12 October as compared to the on-track 334 location. The surface wind speed at the on-track location shows a typical temporal variations of a 335 passing cyclone. The wind speed peaks, drops, and attains second peak as the cyclone approaches, 336 crosses over, and depart the location. The surface currents forced by these large variations in wind 337 speed and direction at the on-track location results into comparatively weaker magnitude than the 338 off-shore location.

The thermocline, defined as the depth of maximum temperature gradient, is usually 339 340 referred to a location dependent isotherm depth (Kessler, 1990; Wang et al, 2000). Over the BoB region, the depth of 23°C isotherm (D23) found to be an appropriate representative depth of the 341 342 thermocline (Girishkumar et al., 2013). Based on the density criteria, we calculated the oceanic 343 mixed layer depth (MLD) as the depth where density increased by 0.125 kg m<sup>-3</sup> from its surface value. The inertial mixing introduced by the cyclone play central role in deepening of D23 and 344 MLD on 12th October 2013. The warmer near-surface waters mixed downward when the cyclone 345 346 crossed over this location. After the passage of cyclone, shoaling of D23 and MLD observed as a 347 consequence of cyclone induced upwelling that entrain colder waters from the thermocline into the mixed layer. The Ttemperature of the upper surface water (25-m -30 m) decreased by 3.5°C 348 from its maximum value of 28 °C after the landfall of the cyclone on 12-13th October at the off-349 350 shore location (Figure 6g4a). In response to the strong cyclonic winds, the depth of 23 °C isotherm 351 (D23) deepening by 40 m (from 50 m to 90 m) was observed during 04-12 GMT on 12 October. 352 At the same time, the mixed layer depth (MLD), denoted by a thick black line in Figure 6g, deepens 353 by about 15 m. On the other hand, the on-track location showed cooling at the surface only for a short time on 13 October and the deepening of D23 and MLD were 20 m and 10 m, respectively.
 The inertial mixing introduced by the cyclone play central role in deepening of D23 and MLD on
 12<sup>th</sup> October 2013. To examine the role of cyclone induced mixing in modulating the thermohaline
 structure of upper ocean upper ocean, we carried out further analysis on the coupled model
 simulations as discussed in the following sections.

359

# 360 <u>3.2.1.</u>3.3.1. Kinetic energy distribution

361 During the initial phase of <del>VSCS</del> Phailin, the zonal and meridional currents were primarily 362 westward and southward, respectively (Figures 6c, 6d, 6h, and 6i). 4b, 4c). However, on and after 12<sup>th</sup> October when cyclone attains peak intensity and crosses over the location, alternative temporal 363 364 sequences of westward/eastward-movement in zonal current and southward/northward-flow in 365 meridional current were noticed in current profiles (Figures 6).4b, 4c). The Ffrequency of these 366 reversals in zonal and meridional currents are recognized as near-inertial frequency generated from 367 the storm at these locations. is location (18.75 °N, 86.66 °E). The direction and magnitude of currents represent a variability within 16-24 hr that corresponds to the presence of near-inertial 368 369 oscillations at time period for the selected locations. The Kkinetic energy (KE) of currents at 370 various depths is a proxy of energy available in the water column that becomes conducive forto 371 turbulent and inertial mixing. Time series of KE associated with the barotropic and depth--averaged 372 baroclinic components of current at the two point locations (18.75° N, 86.66° E) areare illustrated 373 in Figure 6e (on-track) and 6j (off-track). 4d. The KE associated with the baroclinic component 374 found to be much higher than the barotropic component of current at the both on-track and off-375 shore locations. The depth-averaged baroclinic and barotropic current components' KE also depict the impinging oscillatory behavior. The peak magnitude of KE in baroclinic and barotropic 376 currents at the off-shore location found to be 1.2 m<sup>2</sup> s<sup>-2</sup> and 0.3×10<sup>-2</sup> m<sup>2</sup> s<sup>-2</sup>, respectively on 12<sup>th</sup> 377 378 October at 08:00 GMT. Whereas the magnitude of KE in baroclinic and barotropic currents at the 379 on-shore location were smaller than the off-shore location during the peak intensity of cyclone. 380 The peak magnitude of kinetic energy in baroclinic current at the off-track location was more than 381 double to that of on-track location. The comparatively smaller magnitude of KE at the on-shore 382 location could be associated with the rapid variations in wind speed and direction leading to complex interaction between subsurface currents in the central region of the cyclone. It is worth 383

noting that the time of peak KE in baroclinic currents coincide with the deepening of MLD and
 D23. Therefore, the KE generated in NIO is responsible for sub-surface mixing that acts to deepen
 the mixed layer. The analysis suggests that energy available for mixing process in the water column
 was mostly confined into the baroclinic currents at various depths.

388

# 389 <u>3.2.2.3.3.2.</u> Primary frequency and depth of mixing

390 The power spectrum analysis was performed on the time series profiles at the two selected locations selected point location (18.75 °N, 86.66 °E) to get a distribution of all frequencies 391 392 operating in the mixing process during the passage of Phailin. eyclone. As found in the previous 393 section, the KE associated with baroclinic currents are dominated over the barotropic currents, the 394 The power spectrum analysis performed on the zonal and meridional components of the baroclinic 395 current profile and is shown in Figure 75. It is clear from the figure Figure 5-that the tidal (M2, the 396 semidiurnal component of tide) and near-inertial oscillations (f) are the two dominant frequencies 397 on the surface during the cyclone Phailin. Further, the near-inertial frequency is smaller than the tidal frequency on the surface. Under the influence of cyclonic winds, the NIO signal was stronger 398 399 (0.84 m2s-2) at the off-track than the on-track location. The depth penetration of NIO was up to 400 50 m and 35 m at the off-track and on-track location, respectively. The tidal frequency (M2) and inertial frequency (f) bands shown in the Figure 7 implies that the inertial oscillations were 401 dominant over the tidal constituent in zonal and meridional baroclinic currents. To analyze the 402 mixing potential of the NIO, power spectrum method was applied at the profile of baroclinic 403 404 current component (Figure 5). At the off-track location, Thethe largest power of the NIO was 405 noticed at 14 m depth but the tidal oscillations were almost absent in the along the whole vertical 406 section of baroclinic current (Figure 7). This finding motivated us to analyze the significance and 407 distribution of these this sub-surface variability that resulted into an anomalous deepening of MLD. 408 The Hhighest power of this signal at the off-track location was associated within 0-15 m with the 409 magnitude of 0.84 m<sup>2</sup> s<sup>-1</sup> in zonal baroclinic current and within 0-38 m with the magnitude of 0.76  $m^2 s^{-1}$  in meridional baroclinic current. These signals, however, weaken with increasing depth and 410 411 almost disappeared around 120 m depth. These NIO were are the strongest signals at the 14 m depth in presence of local wind stress that and dominating dominated the mixing compared to any 412 other process.by any other process than the local wind stress. Other processes include the 413

414	background flows, the presence of eddies, variations in sea surface height, non-linear wave-wave
415	and wave-current interactions (Guan et al., 2014; Park and Watts, 2005).
416	The second order butterworth filter was applied on the baroclinic current components in
417	order to get the strength of NIO in the frequency range of 0.028 to 0.038 cycles h <sup>-1</sup> at the selected
418	locations. The filtered baroclinic current was further utilized to calculate the filtered inertial
419	baroclinic KE ( $E_f$ in $m_1^2 s_2^{-2}$ ). The daily profiles of baroclinic KE were analysed at the two selected
420	locations and shown in Figure 8. The peak baroclinic KE differs from 0.14 m <sup>2</sup> s <sup>2</sup> at the on-track to
421	0.23 m <sup>2</sup> s <sup>-2</sup> at the off-track location on 12 October. As shown in Figures 6 and 7, the filtered
422	baroclinic KE profiles (Figure 8) confirm the dominant presence of NIO at the off-track location
423	as compared to the on-track location. The decay of NIO with the increasing depth was noticed at
424	both the locations. However, the NIO baroclinic KE penetrated up to 80 m in case of off-track as
425	compared to only 50 m at the on-track location. The analysis, therefore, suggests that the NIO
426	generated during the Phailin were more energetic at the selected off-track location, which was also
427	the location of maximum surface cooling as noticed in Figure 5. Therefore, the further analysis in
428	the subsequent sections is limited to the off-track location only. In order to analyze the time
429	distribution of the strong NIO, wavelet transform analysis was applied on the zonal and meridional
430	baroclinic currents at 14 m depth. The Scalogram, shown in Figure <u>96</u> , depicts the generation of
431	NIO signal at the off-track location on 12 <sup>th</sup> October that subsequently got strengthen and attains
432	its peak value on the mid of 13 <sup>th</sup> October. The energy percentage of <u>the</u> meridional component was
433	always lower than the zonal component. The peak values of energy percentage was found in the
434	time periods between 1-1.3 days.25-28 hr marked with a white dashed line in Figure 6. A
435	Butterworth 2 <sup>nd</sup> -order band pass filter was applied at the corresponding cutoff frequency interval
436	of 0.033 - 0.043 to filter the NIO signal of the baroclinic zonal and meridional current. Figure 7
437	shows profiles of near inertial zonal $(U_f)$ and meridional $(V_f)$ baroclinic current together with the
438	kinetic energy $(E_f)$ of near-inertial flow. The maximum strength of inertial baroclinic current was
439	0.3 m s <sup>-1</sup> with the signature of an alternate directional reversal of current signals. Presence of these
440	inertial currents were up to 70 m depth with peak value of kinetic energy $E_F$ being 0.048 m <sup>2</sup> s <sup>-2</sup> . It
441	can be noticed from Figure 7c, the baroclinic kinetic energy remains higher (> 0.03 m <sup>2</sup> s <sup>-2</sup> ) only
442	from mid of 11 <sup>th</sup> October till the end of 12 <sup>th</sup> October and thereafter the energy rapidly decreases
443	and almost disappeared after 13th October. This indicates the period of prominent mixing due to

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444 NIO was 11–12 October 2013. The daily averaged values of baroclinic kinetic energy (not shown 445 here) also confirms the maxima in  $E_{\rm f}$  on 11–12 October with the vertical extent up to 80 m depth.

# 446 <u>3.2.3.3.3.3.</u> Role of downward propagation of energy

447 To investigate the energy propagation from the surface to the interior layers of upperocean, we derived the rotary spectra (Gonella, 1972; Hayashi, 1979) of near-inertial wave numbers 448 449 and shown in Figure 10-8. The daily averaged vertical wave-number rotary spectra provides a clear 450 picture of wind energy distribution in the sub-surface water. The anticyclonic spectrum (A<sub>m</sub>) is dominating over the cyclonic spectra  $(C_m)$  for the entire duration of the cyclone. This feature 451 indicates that the energy is propagating downward generated by these inertial oscillations. The 452 Mmagnitude of these oscillations increased from initial stage up to 12<sup>th</sup> October and remained at 453 454 high energy density for the rest of the cyclone period. This downward directed energy initiated a process of mixing between the mixed layer and the thermocline. This energy helps to deepen the 455 456 mixed layer against oceanic stratification by introducing a strong shear. The buoyancy of stratified ocean was overcome to some extent by the shear generated that assist in mixing process during the 457 very severe cyclone. For the current case, kinetic energy (Figure 7c) represents the analogical 458 behavior as reported by Alford and Gregg (2001). Their study Alford and Gregg (2001) highlighted 459 460 that, in most of the cases, the energy of inertial oscillations potentially penetrates the mixed layer but suddenly drops down as it touches the thermocline. The energy dissipation mechanism studied 461 462 in few other studies (Chant, 2001; Jacob, 2003).

463 The 2-layer model described by Burchard and Rippeth (2009) illustrated the process of 464 generation of sufficient shear to start mixing near the thermocline. Their simple model ignored the effect of the lateral density gradient, mixing, and advection. Burchard et al. (2009) mentioned four 465 important parameters for the shear generation, i.e. surface wind stress ( $P_SS^2$ ), bed stress ( $-D_bS^2$ ), 466 interfacial stress (- $D_1S^2$ ), and barotropic flow ( $P_mS^2$ ). Utilizing simulations from our coupled 467 468 atmosphere-ocean-wave model, we calculated individual terms as suggested by Burchard et al. 469 (2009) and presented in Figure 910. It is clear from the figure that the surface wind stress term 470 plays most significant role Surface wind stress found to be the most dominating term in modulating 471 the magnitude of bulk shear during the stormy event. Rest of the terms were Other terms were found to be relatively weaker and, therefore, contributing only marginally in to the variability of 472 473 the bulk shear.

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474 To examine the generation and dissipation of these inertial oscillations, waves, the shear generated by the near-inertial baroclinic current  $(S_{\ell}^2)$  and turbulent kinetic energy dissipation rate 475 476  $(\varepsilon)$  were calculated and analyzed. The shear produced by inertial oscillations-was increasing 477 increased at 20-80 m depth from 40-80 m depth and higher magnitude was associated with peak 478 wind speed of cyclone (Figure 10a12a). This shear overcome the stratification (Figure 10b12b). 479 represented by buoyancy frequency  $N_{a}^{2}$ , and played important role in mixing and deepening of the thermocline and mixed layer on 12<sup>th</sup> October, that was weak at this depth compared to the shear 480 of the near inertial waves. The value of kinetic energy dissipation rate ( $\varepsilon$ ) increased from  $4 \times 10^{-10}$ 481 <sup>14</sup> to  $2.5 \times 10^{-13}$  W kg<sup>-1</sup> on approaching the thermocline (Figure 10e12c). The increase in  $\varepsilon$  indicates 482 the weakening of the shear generated by the inertial waves leading to the fast disappearance of 483 484 these baroclinic instabilities from the region. The non-linear interaction between the NIO and 485 internal tides together with the prevailing background currents cause rapid dissipation of kinetic 486 energy in the thermocline. Guan et al. (2014) also reported an accelerated dampening of NIO 487 associated with the wave-wave interactions between NIO and internal tides. The background currents found to modify the propagation of NIO (Park and Watts, 2005). The magnitude of the 488 489 turbulent eddy diffusivity ( $K_a$ ), shown in Figure 10d12d, implies that the greater mixing takes place within the mixed layer-place where  $K_0$ -was high  $(6.3 \times 10^{-11} \text{ to } 1.2 \times 10^{-11} \text{ m}^2 \text{ s}^{-1})$ . The daily 490 averaged values of  $\varepsilon$  and  $K_{\rho}$  were  $1.2 \times 10^{-13}$  W kg<sup>-1</sup> and  $1.5 \times 10^{-10}$  m<sup>2</sup> s<sup>-1</sup>, respectively on  $12^{\text{th}}$ 491 492 October, which were higher as compared to the initial two days of the cyclonic event. Therefore, 493  $+\mathbf{R}$  esults from the present study, as well as the conclusions from the past studies, indicate that wave-current interaction, mesoscale processes, and wave-wave interaction can affect the process 494 495 of downward mixing and cause the dissipation of inertial oscillations.

### 496 4. Conclusions

497 Processes controlling the sub-surface mixing were evaluated under the high wind speed 498 regime of a severe cyclonic storm Phailin over the BoB. A coupled atmosphere-ocean-wave 499 (WRF+ROMS+SWAN) model as part of the COAWST modeling system was used to simulate 490 atmospheric and oceanic conditions during the passage of Phailin cyclone. A detailed analysis of 501 model simulated data revealed interesting features of generation, propagation, and dissipation of 502 kinetic energy in the upper oceanic water column. Deepening of the MLD and thermocline by 15 503 m and 40 m, respectively were explained through the strong shear generated by the inertial Formatted: Superscript

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504 oscillations that helped to overcome the stratification and initiate mixing at the base of the mixed 505 layer. However, there was a rapid dissipation of the shear with increasing depth below the thermocline. The peak magnitude of kinetic energy in baroclinic and barotropic currents found to 506 be  $1.2 \text{ m}^2 \text{ s}^{-2}$  and  $0.3 \times 10^{-2} \text{ m}^2 \text{ s}^{-2}$ , respectively. The power spectrum analysis suggested a dominant 507 frequency operative in sub-surface mixing that was associated with near-inertial oscillations. The 508 peak strength of 0.84 m<sup>2</sup> s<sup>-1</sup> in the zonal baroclinic current found at 14 m depth at a location in 509 northwestern BoB. The baroclinic kinetic energy remains higher (> 0.03 m<sup>2</sup> s<sup>-2</sup>) during 11-12 510 October and decreased rapidly thereafter. The wave-number rotary spectra identified the 511 512 downward propagation, from the surface up to the thermocline, of energy generated by inertial oscillations. A quantitative analysis of shear generated by the near-inertial baroclinic current 513 514 showed higher shear generation at 20-80 m depth during peak surface winds. Analysis highlights 515 that greater mixing within the mixed layer takes place where the eddy kinetic diffusivity was high  $(> 6 \times 10^{-11} \text{ m}^2 \text{ s}^{-1})$ . The turbulent kinetic energy dissipation rate increased from  $4 \times 10^{-14}$  to  $2.5 \times 10^{-14}$ 516 <sup>13</sup> W kg<sup>-1</sup> on approaching the thermocline that dampened mixing process further down into the 517 thermocline layer. Kinetic energy associated with baroclinic currents were about two order of 518 magnitudes higher than in barotropic component. The peak strength of 0.84 m<sup>2</sup> s<sup>4</sup> in zonal 519 520 baroclinic current was found at 14 m depth at a location in northwestern BoB. The wave-current 521 interaction, mesoscale processes, and wave-wave interaction increased the dissipation rate of shear 522 and, thereby, limited the downward mixing up to the thermocline. were found to affect the process 523 of downward mixing and cause the dissipation of inertial oscillations. The coupled model found 524 to be a useful tool to investigate air-sea interaction, kinetic energy propagation, and mixing in the upper-ocean- and oceanic sub-surface processes. The results from this study highlight the 525 importance of atmosphere-ocean coupling for better understanding of oceanic response under the 526 527 strong wind conditions. The proper representation of kinetic energy propagation and oceanic 528 mixing have applications in improving the intensity prediction of cyclone, storm surge forecasting, 529 and biological productivity.

530

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Figure 24:-COAWST model domain (65°-105 °E, 1°-34 °N) overlaid with GEBCO bathymetry (m).
 Locations used for time-series analysis are marked with a-stars.





Figure <u>32</u>:- Validation of VSCS Phailin track simulated by coupled model (black) with IMD
 reported track (red), Tracks of Phailin simulated by the coupled model (black) and IMD reported
 (red). The 3-hourly positions of the center of Phailin marked with solid circles and the daily
 position at 00 hr are labelled with the dates. Location of buoy BD09 is marked with a blue circle.





Figure 53:- The daily averaged sea surface temperature Sea Surface Temperature (SST) in °C 807 simulated by the coupled model (upper panel), stand alone ROMS model (middle panel), and observed from AVHRR sensor on the satellite (lower panel). and simulated by model (upper panel).











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Figure 75:- The power spectrum analysis (m<sup>2</sup>s<sup>-1</sup>) performed on the simulation period at the on-track
(upper panel) and off-track (lower panel) locations for (a, c) baroclinic zonal current and (b, d)
baroclinic meridional current.







83411 OCT12 OCT13 OCT14 OCT15 OCT835Figure 26:- The scalogram in percentage at -40\_14 m depth by continuous wavelet transform (CWT)836method. Wavelet scalogram shown for the zonal baroclinic current (upper panel) and for the837meridional baroclinic current (lower panel). The white dashed line indicates the peak percentage of838energy.





**Figure 108:** The daily averaged vertical wave-number rotary spectra of near inertial oscillations.











Figure 1240:- Profiles of (a) velocity shear  $log_{10}(S^2)$ , (b) buoyancy frequency  $log_{10}(N^2)$ , (c) turbulent kinetic energy dissipation rate  $log_{10}(\varepsilon)$ , (d) turbulent eddy diffusivity  $log_{10}(K\rho)$ , (e) and (f) are daily averaged turbulent kinetic energy dissipation rate and turbulent eddy diffusivity respectively