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Deep Western Boundary Current transport variability in the South Atlantic: preliminary results from a pilot array at 34.5° S

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Abstract

The first direct estimates of the temporal variability of the absolute transport of the Deep Western Boundary Current (DWBC) at 34.5° S in the South Atlantic Ocean are obtained using just under one year of data from a line of four pressure-equipped inverted echo sounders. Hydrographic sections collected in 2009 and 2010 confirm the presence of the DWBC, one of the main deep pathways of the Meridional Overturning Circulation, based on neutral density, temperature, salinity, and oxygen values. Both observations confirm that the DWBC reconstitutes itself after breaking into eddies in the western sub-tropical Atlantic near 8° S. The amplitude and spectral character of the DWBC transport variability are comparable with those observed at 26.5° N, where longer records exist, with the DWBC at 34.5° S exhibiting a transport standard deviation of 25 Sv and variations of ~ 40 Sv occurring within periods as short as a few days. There is little indication of an annual cycle in the DWBC transports, although the observation record is too short to be definitive, and the dominant time scale during the first year of the experiment was about 9–10 days. A “Monte Carlo-style” analysis using 27 yr of model output from the same location as the observations indicates that another 48–60 months of data will be required to encompass a fairly complete span of deep transport variability. The model suggests the presence of an annual cycle in DWBC transport, however the statistical significance of the annual cycle with even 27 yr of model output is low, suggesting that annual period variations in the model are weak as well.

1 Introduction

The role of the Deep Western Boundary Current (DWBC) as a primary pathway for the cold, lower, limb of the Meridional Overturning Circulation has been well documented in the North Atlantic Ocean (e.g., Molinari et al., 1998; Schott et al., 2004; Johns et al., 2008; Meinen et al., 2012), however the pathways and variability of the

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DWBC in the South Atlantic Ocean are less well known. Near 8° S the DWBC appears to break up into rings as it flows southward (Dengler et al., 2004; Schott et al., 2005). It has been suggested that a significant fraction of the DWBC flow transits eastward across the basin in the tropics and/or subtropics and flows southward along the African coast near 30–35° S (e.g., Warren and Speer, 1991; Speer et al., 1995). Most of this southeastward flow appears to continue along the African coast in the Cape Basin to then enter the Indian Ocean (Speer et al., 1995; van Aken et al., 2004). There are few observations of the portion of the DWBC that remains along the western boundary south of 8° S aside from a small number of sections analyzed using an assumed level of no motion (e.g., Zemba, 1991). The only direct current meter estimates available in the region are from off Cabo Frio (22° S) and Cabo Santa Marta (28° S). These 23-month records indicate a very weak southward flow of the DWBC: -0.5 ± 1.6 and -2.8 ± 4.9 Sv, respectively (Müller et al., 1998). The southward geostrophic flow at the western boundary at 28° S is estimated at 10 Sv, but about 4 Sv recirculate northward in the interior (Zangenberg and Siedler, 1998). Given the South Atlantic Ocean's role as a “blender” of water masses in the MOC (e.g., Garzoli and Matano, 2011), and the indications that the DWBC plays a major role in the meridional heat transport at these latitudes (e.g., Dong et al., 2011), it is important to understand how much of the deep cold limb is transiting the basin and reaching the Southern Ocean over time. The purposes of this article are to describe preliminary results from an array of pressure-equipped inverted echo sounders (PIES) deployed across the DWBC at 34.5° S on the western boundary of the South Atlantic Ocean just north of the Brazil-Malvinas Confluence, and also to demonstrate that the pilot array of PIES is successfully observing the DWBC. The variability observed during the first year of deployment will be compared to 27 yr of output from a high-resolution model to evaluate the statistical information contained in a one-year record.

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2 Data and methods

PIES technology has been in use in a variety of forms for several decades, and a large number of previous articles have described the instrument and the data it collects. Only a brief review of the instrument is presented here with appropriate references for more information. In essence a PIES makes two measurements: the bottom pressure and the round-trip travel time required for a 12 kHz acoustic pulse to travel from the bottom-moored instrument up to the sea surface and back. Bottom pressure is measured with a highly precise Paros pressure gauge (e.g., Watts and Kontoyiannis, 1990; Donohue et al., 2010), and the round-trip travel time is determined using a transducer and a high quality crystal clock (e.g., Rossby, 1969; Watts and Rossby, 1977; Tracey and Watts, 1986). The travel time measurements from each PIES are calibrated into daily, full-water-column, profiles of temperature, salinity and specific volume anomaly via hydrography-derived look-up tables using the Gravest Empirical Mode (GEM) technique (e.g., Meinen and Watts, 2000; Watts et al., 2001).

Vertically integrating the specific volume anomaly profiles yields dynamic height anomaly profiles, and differencing dynamic height anomaly profiles between neighboring PIES sites provides geostrophic relative velocity profiles orthogonal to the line between the PIES (e.g., Meinen et al., 2006). Differences in bottom pressure from neighboring PIES sites provide absolute geostrophic velocity variability at the bottom that can be used to reference the relative velocity profiles. Due to the well-known leveling problem,¹ however, the time-mean absolute geostrophic velocity at the bottom cannot be determined from the bottom pressure differences (e.g., Donohue et al., 2010). If an independent estimate of the time-mean bottom velocity is available for the re-

¹In brief, a time-mean pressure difference between two neighboring sensors can occur due to the sensors being on the bottom at different depths, or it can occur with two sensors that are at the same depth but there is a time-mean geostrophic current orthogonal to the line between them. With only the two sensors, there is insufficient information to discriminate between these two scenarios.

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gion, from historical current meter records, for example, or from concurrent ship-based velocity sections, then this mean can be added to the bottom-pressure-derived time-varying absolute velocities. When independent data are not available, such as is the case for the region at 34.5° S, then the time-mean bottom velocity between pairs of PIES must be derived from another source, e.g. from a high quality numerical ocean model. The model used in this study for providing the mean bottom velocity will be discussed shortly. Once a time-mean has been added to the time-varying absolute geostrophic velocities determined from the bottom pressure gauges, the resulting absolute velocity time series can be used to reference the relative geostrophic velocity profiles determined from the travel time and GEM look-up tables. The result is full-water-column time series of absolute velocity perpendicular to the line between each pair of PIES.

A primary goal of this paper is to discuss the time variability of the DWBC. Defining an integration domain to call the DWBC is somewhat tricky, as the water mass definitions one might use (which will be discussed shortly) in some cases require estimates of dissolved oxygen, which the PIES (as with most moorings) do not provide. For the purposes of this study, therefore, a similar vertical integration domain will be used as was applied in recent work such as at 26.5° N (e.g., Meinen et al., 2012); the DWBC transport will be defined as the integral from 800 dbar to 4800 dbar (or the bottom where it is shallower than 4800 dbar). This allows comparison with the results at other latitudes. The character of the transport time series that will be shown is not sensitive to modest, 100–300 dbar, changes in these integration limits.

The details of the techniques for using PIES to measure the transport of the DWBC have been developed and tested versus other measurement system in the North Atlantic at 26.5° N (Meinen et al., 2004, 2006, 2012). The PIES-derived absolute transports were very similar to those determined from current meters and dynamic height moorings, with correlation coefficients between the absolute transport time series exceeding $r = 0.9$ and root-mean-square differences of a few Sv ($1 \text{ Sv} = 10^6 \text{ m}^3 \text{ s}^{-1}$) over

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a roughly equivalent integration domain (correlations for baroclinic transports using an assumed level of no motion at 800 dbar were similarly good).

The PIES data presented here are from a line of four instruments deployed along 34.5° S at 51.5° W (Site A), 49.5° W (Site B), 47.5° W (Site C), and 44.5° W (Site D) as a pilot array to measure the western boundary components of the MOC (Fig. 1). One instrument was additionally equipped with a single-depth current meter (CPIES), but the current meter data is not crucial for the purposes of this paper and will not be discussed herein. All instruments were deployed in March 2009, and data were acoustically downloaded from the four instruments in July 2010 and again in December 2010 and July 2011. Due to acoustic transmission issues the data record from one instrument (Site B: see Fig. 1) is more limited than the others. For this preliminary study, data from the ~ 10.5 month period when all four records are available (6 May 2009–22 March 2010) will be presented. Data values are at daily resolution, with all records having been low-pass filtered with a cut-off period of three days.

Data from two hydrographic sections completed as part of this pilot study are also presented herein to describe the water masses observed in the region. The sections were occupied during August 20–24, 2009 and July 7–11, 2010. Both cruises were completed onboard the Argentine research vessel Puerto Deseado.

3 Model description

Detailed absolute velocity observations of the DWBC in the South Atlantic are limited. As such, coupling observational knowledge with high-quality numerical model output presents a useful opportunity to advance understanding beyond what either can provide alone. For this study, a numerical model was used for two purposes. First was to supply the near-bottom time-mean absolute velocity that needs to be combined with the PIES bottom-pressure-derived time-varying absolute geostrophic velocity anomalies to yield the full near-bottom absolute geostrophic velocity. Secondly the model was

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used to obtain an estimate of the statistical stability of the deep flows from short records relative to longer-period variations.

The model output selected for this study is from a run of the Ocean general circulation model For the Earth Simulator (OFES; e.g. Sasaki et al., 2008). The OFES model is a massively parallelized implementation of the NOAA/GFDL Modular Ocean Model version 3 (MOM3) being executed by the Japan Agency for Marine-Earth Science and Technology (JAMSTEC). The model equations have been discretized in a Mercator B-grid with a horizontal resolution of 0.1° and 54 vertical z levels. For this study, model fields were provided by JAMSTEC at 0.2° increments (every other horizontal grid point) at 3-day increments (snapshots, not 3-day averages) over the period from 1980–2006. The model was spun up for 50-yr with a monthly climatology derived from NCEP/NCAR reanalysis atmospheric fluxes (Masumoto et al., 2004), and then forced with daily mean NCEP/NCAR reanalysis data from 1950 to 2006 (Sasaki et al., 2008). Only the data from the final 27 yr of the run were used herein. This model has previously been successfully validated against both other models and the limited available observations in the South Atlantic (Perez et al., 2011; Dong et al., 2011; Giarolla, 2010, personal communication).

To obtain the time-mean absolute velocity near the bottom for use with the actual bottom pressure measurements, the velocities from the model were first temporally averaged over the full 27 yr of the model run, and then these time-mean velocities were horizontally averaged between the longitudes of the pilot array moorings. The mean meridional velocities from the model are shown in Fig. 2 along with the nominal locations of the PIES/CPIES discussed in this article. The mean velocity of the three deepest model layers was then averaged in order to obtain the mean reference velocity (the results are not particularly sensitive to different selections of the deepest levels – differences are less than 0.2 cm s^{-1}). This provides a “best estimate” for the geostrophic time-mean absolute velocity at the bottom between each pair of moorings.

To obtain transports from the model, the model velocities were integrated between the nearest grid points to the four pilot array sites for comparison to the observed data.

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Note that transport integration domains in the model are not exactly the same as in the real ocean because the model topography is not identical to the real ocean topography. The mean bottom velocities from the model as well as the model and real ocean depths at the four sites are shown in Table 1.

4 Results

In this region the precise location and variability of the DWBC are not as well known as in some other regions due to the paucity of velocity observations. However, the presence of the DWBC can clearly be demonstrated via hydrographic observations such as temperature, salinity, dissolved oxygen and nutrient sections (e.g., Reid et al., 1977; Zemba, 1991; Piola and Matano, 2001). The higher-than-ambient dissolved oxygen signal at ~ 2500 dbar clearly indicates more recently ventilated waters, although as indicated in the July 2010 section (Fig. 3) the signal is not always as tightly confined along the boundary as might be expected for the DWBC. The selection of the A, B, C and D sites for the pilot array (Fig. 1) was designed based on previous hydrographic and IES observations to capture the DWBC flow and allows for offshore meanders/shifts as far as 44.5° W. Overlaying the neutral density surfaces, calculated following Jackett and McDougall (1997), can help identify water masses being carried meridionally across the array. Based on an analysis of deep water-masses in the northwest Argentine Basin, Preu et al. (2012) proposed the following water mass boundaries/definitions:

- Antarctic Intermediate Water (AAIW): Salinity less than 34.25 psu
- Upper Circumpolar Deep Water (UCDW): neutral density between 27.75 and 27.9 with dissolved oxygen values below 4.5 ml l^{-1}
- North Atlantic Deep Water (NADW): neutral density between 27.9 and 28.1 with salinity greater than 34.8 psu

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- Lower Circumpolar Deep Water (LCDW): neutral density between 28.06 and 28.2 with salinity less than 34.8 psu
- Antarctic Bottom Water (AABW): potential temperature less than 0 °C

Based on these definitions, the deep oxygen maximum observed in July 2010 in Fig. 3 is identifiable as NADW. The oxygen section also clearly depicts a relative minimum ($< 5 \text{ ml l}^{-1}$) below the core of NADW, associated with LCDW, and a near bottom increase to $> 5.1 \text{ ml l}^{-1}$ at depths greater than 4000 dbar, indicative of AABW. Temperature and salinity sections from this cruise (not shown) are consistent with the presence of AAIW, NADW, and AABW.

In August 2009 another CTD section was collected along the mooring line over the course of five days (20–24 August 2009). The potential temperature and salinity sections, both with the neutral density surfaces overlain, shows the clear presence of AAIW, NADW, and AABW along the section with no clear indication of LCDW (Fig. 4, left panels). Unfortunately the oxygen data from this cruise is problematic due to a sensor problem, thus, UCDW in this section can only be identified based on neutral density. The strong preponderance of NADW between the 27.9 to 28.1 neutral surface layers suggests that within the domain of the section the bulk of the 800–4800 dbar waters at the time of this section were of North Atlantic origins.

Comparing the mean potential temperature and salinity sections from the PIES data over those five days (Fig. 4, right panels) to the actual CTD section data (Fig. 4, left panels) illustrates how well the PIES can estimate the general water mass patterns and the layer interfaces (albeit at an admittedly lower horizontal resolution). In general the agreement is quite good. In the upper 100–200 dbar there are some differences related to the reduced seasonal signal that comes from the application of the GEM technique to PIES data. These differences can be reduced in the future through the application of a “seasonal GEM correction” (e.g., Watts et al., 2001). The seasonal differences have no impact on the velocities, however, as the latter are based on horizontal density gradients and not the density at any given point. Below the seasonally-affected layer, the

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agreement through the main thermocline/halocline layer is excellent, while at depths below 2000 dbar there are some more noticeable differences (although some of the difference results from the dissimilar horizontal resolution). Small fluctuations at depth are probably beyond the capability of the PIES-GEM technique to capture – these fluctuations have been shown to have little or no impact on the volume transports (e.g., Meinen et al., 2004), but they do have great importance with regards to understanding the source regions of the water masses (e.g., Molinari et al., 1998; van Sebille et al., 2011).

4.1 Absolute velocity and transport

The time-mean and the temporal standard deviation of the absolute velocity profiles (Fig. 5) indicate large variability throughout the array during the first 10.5 months of the experiment, with the standard deviation generally exceeding the mean. The strong, highly varying, flow in the upper water column on the western side of the array is associated with the Brazil Current, which is thought to be confined above ~ 800 dbar (e.g., Garzoli and Garraffo, 1989), while the upper water column flow on the east side is associated with the meandering of retroflected waters coming from the Brazil/Malvinas confluence to the south (e.g. see velocity vectors in Fig. 1). The estimated near surface velocities ($\sim 15 \text{ cm s}^{-1}$) from the PIES are significantly lower than the $35 \pm 14 \text{ cm s}^{-1}$ mean surface velocity estimated from surface drifters at this location (Oliveira et al., 2009). This is most likely due to two factors: the broad smoothing which results from calculating geostrophic velocity over the 2° – 3° longitudinal spans between PIES; and possibly also due to the drifter velocities including the Ekman flow, which is absent from the PIES transports. Below 800 dbar there is still significant mean flow and variability on the western side of the array, which is associated with the DWBC based on the aforementioned water mass evaluation and historical hydrographic observations in the region (e.g., Zemba, 1991). The time-mean DWBC was found to be strongest between Sites A and B immediately beneath the southward flowing Brazil Current (note that part of the Brazil Current is missed west of site A). The statistical standard error of the mean

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ranges from $\sim 1 \text{ cm s}^{-1}$ below 2000 dbar to roughly 4 cm s^{-1} in the upper 500 dbar between sites B and C, which suggests that the mean pattern is fairly robust except below ~ 1500 dbar offshore between sites B,C and C,D. However, the mean profiles are highly dependent on the OFES model time mean bottom velocities used, so the focus of this analysis is on the time variability. Offshore in the upper water column the mean flow reversal and high variability suggests the presence of strong anti-cyclonic circulations most-likely associated with the retroflexion of the Brazil Current just to the south at the Brazil-Malvinas Confluence (e.g., Garzoli and Garraffo, 1989). This retroflexion appears to be much tighter than was suggested in some earlier studies (e.g., Peterson and Stramma, 1991; Stramma and England, 1999), but it resembles the circulation pattern of satellite-derived mean dynamic topography (Rio and Hernandez, 2004) and the OFES numerical simulation (Fig. 2).² At the DWBC depths (below 800 dbar) the strongest velocity variability is between sites B and C.

These results expressed as velocities might be somewhat deceptive in the sense that they represent horizontal averages over the distance between sites, and those distances are not all the same. The span between sites C and D is 50 % larger than the spans between the other pairs of instruments. As such, it is also instructive to focus on transports integrated between pairs of sites, as the transport integration eliminates this issue. Integrating the absolute transport between each pair of sites and between 800 dbar and 4800 dbar (or the bottom where it is shallower than 4800 dbar) indicates that the deep flow is quite variable in all three spans (Fig. 6). The standard deviations of the daily time series³ are 12 Sv, 27 Sv, and 28 Sv for the A-to-B, B-to-C and C-to-D spans, respectively. This indicates that the transport variability is equally high offshore in the C-to-D span as it is in the B-to-C span nearer the slope. The variability in the A-to-B span is significantly weaker than that in the B-to-C span. If the observed time variability is based primarily on zonal movement of quasi-stable velocity signals, e.g.

²Recall that only the deepest levels in the model were used to add to the measured bottom pressure gradients, so the shallower levels are independent.

³Recall that the PIES records have been low-pass filtered with a 3-day cutoff.

the meandering of a fairly stable DWBC, then moving the velocity signal across integration domains of significantly different size will result in transport amplitude fluctuations of different magnitudes, and the difference in amplitudes between the A-to-B and B-to-C spans is roughly consistent with the smaller integration domain of the A-to-B span in the deep layer due to the sloping topography (e.g. see Fig. 5).

The transport within each span shows variability on time scales ranging from a week to a month or two, and each exhibits transport changes exceeding 20 Sv on extremely fast time scales (Fig. 6). While there are certain events that suggest anti-correlation between the more highly variable offshore B-to-C and C-to-D time series (e.g. the southward transport maximum in the B-to-C span in December 2009 and the corresponding northward transport maximum in the C-to-D span), those time series are not correlated with one another in a statistically significant way ($r = -0.38$). The two more inshore spans, A-to-B and B-to-C, are significantly anti-correlated at the 95 % level ($r = -0.53$). Note this significance evaluation is based on the estimated integral time scale of 9–10 days for the three records (calculated via the methods described in Emery and Thomson, 1997) and the requirement for two integral time scales per degree of freedom based on the lag integration limits (see Appendix B in Meinen et al., 2009 for more information). Despite its significance, the anti-correlation between the A-to-B and B-to-C deep transports is quite modest, and a linear relationship between the two would only explain $\sim 25\%$ of the observed variance. This lack of strong correlation in the presence of very high amplitude variations is similar to that observed at 26.5°N (e.g., Meinen et al., 2012), and it illustrates the importance of integrating over a fairly large domain in order to average out small-scale features that are likely not related to overall DWBC variations.

Integrating over the complete array from sites A to D in the deep layer (800–4800 dbar) yields a transport that varies from a northward maximum of +46 Sv to a southward maximum of –83 Sv, with variations exceeding 40 Sv over very short time scales (~ 1 week) and a standard deviation of 25 Sv (Fig. 7). The mean southward transport of –17 Sv is dependent on the mean bottom velocities used from OFES, and

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as such should not be a focus here. The standard deviation of the transport is roughly comparable with that found over 5 yr of data from an array of moorings stretching a similar distance offshore at 26.5° N (Meinen et al., 2012).

Perhaps the most interesting result from this analysis of the pilot array data becomes evident when the absolute transport (black solid line in Fig. 7) is compared to the baroclinic transport relative to an assumed level of no motion (cf. at 800 dbar; dark gray dashed line in Fig. 7). These two time series are completely uncorrelated with one another ($r = -0.23$), and at times can disagree by as much as 50–100 Sv (e.g. early December 2009, and early March 2010). This strongly illustrates the point, raised previously in analyses of XBT data and numerical model output (e.g., Baringer and Garzoli, 2007; Garzoli and Baringer, 2007), that the barotropic flows are strong and need to be measured to study the absolute flow near the western boundary near 34.5° S. If the transport associated with the true velocity at the assumed level of no motion is integrated over the DWBC domain (light gray dash-dot line in Fig. 7) it has a significantly higher standard deviation (32 Sv) than that of the true absolute transport (25 Sv) or the baroclinic transports (18 Sv).

Independent validation of the velocity and transport data is difficult, as there are no other in situ observations at this latitude, however the hydrographic observations (Figs. 3 and 4) suggest that the absolute velocity section (Fig. 5) is quite realistic:

1. The low oxygen water above the NADW, associated with UCDW, is higher ($\sim 4.42 \text{ ml l}^{-1}$) on the boundary at ~ 1365 dbar, where the flow is southward, than at the easternmost station (4.18 ml l^{-1} at 1400 dbar), where the flow is northward. Thus the westernmost UCDW core is recirculating southward along the boundary and increasing its O_2 concentration by vertical mixing.
2. The high oxygen ($> 5.76 \text{ ml l}^{-1}$) NADW core at ~ 2500 dbar close to the boundary is part of the southward flowing DWBC while the offshore ~ 2400 – 2800 dbar core (oxygen $> 5.68 \text{ ml l}^{-1}$) located near 47° W is part of the northward recirculation (Figs. 3 and 5). These cores are separated by a low oxygen region ($< 5.3 \text{ ml l}^{-1}$).

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This is further confirmed by the salinity distribution, as the bulk of the high salinity deep water (> 34.94 psu) is only observed in the westernmost station (Fig. 4).

3. The lowest oxygen (< 4.8 ml l^{-1}) observed below the NADW core, associated with LCDW, is found at ~ 3000 – 3500 dbar at the easternmost station (Fig. 3), where the mean flow is weakly northward (Fig. 5).
4. Finally, the high oxygen bottom core, associated with AABW, is located mostly east of 49° W, and is therefore in the region of northward flow below the recirculated NADW.

Thus to the extent that the hydrographic observations made on these cruises are representative of the mean over the first year of the study, the mean velocity section (utilizing the deep OFES means) is consistent with the hydrographic information.

5 Discussion

At this early stage of analysis of the pilot array, the most important results are likely to come from a joint analysis of the data with the output from a high-quality, high-resolution general circulation models such as OFES. The standard deviation of transport integrated over the same domain within the three-day subsampled model output over the full 27 yr of the run used herein is 16 Sv (Fig. 8, upper panel). This is only about two-thirds the standard deviation of the observed absolute transports, however the time period is quite different (10.5 months of data versus 27 yr of model output). A simple Monte Carlo-style test using 1000 random 10.5 month subsets of the 27 yr model record suggests that a record of 10.5 months length in the model could have standard deviations between 8 and 24 Sv, with the mean and median standard deviations of records of that length being 14 and 13 Sv, respectively. Because the largest standard deviation for a 10.5 month record in the model (24 Sv) is roughly equal to that of the observed data from the pilot array (25 Sv), it cannot be definitively stated that the

variability of the model is too low, however it seems likely that the model is underestimating the true transport variability since the true data only overlaps with one extreme end of the model range.

While the comparison of the short data records to the much longer OFES transports suggests the latter may be underestimating the true amplitude of the variability, it is still potentially instructive to evaluate the spectrum of the long model run to evaluate character of the variability that might be expected to be observed once the pilot array records are longer (Fig. 8, lower panel). Using a 5-yr window for the spectral calculations, the largest energy peak is at a period of about 202 days, close to the semiannual period, and there is a much lower, much broader peak close to the annual period. An annual climatology of the OFES transport record (Fig. 9) suggests a weak (± 5 Sv) annual cycle with the maximum southward transport in late October (austral spring) and maxima in northward transport in February and May–June (austral summer and late austral fall). The scatter of the daily values is quite high, however this small annual cycle is (barely) statistically significant from zero at the 95% confidence level (dashed lines in Fig. 9). There is no obvious indication of a significant semiannual or annual period to the observed record (Fig. 7), but, given the large higher frequency variability, with only just under one year of data from the pilot array it is premature to draw a conclusion with regards to the presence or lack of an annual cycle in the DWBC transports at this latitude. An analysis of the DWBC transports integrated a similar distance offshore at 26.5° N using a 5-yr record illustrated no statistically significant annual or semiannual cycle at that location, however the calculation was somewhat dependent on the distance integrated offshore (Meinen et al., 2012). When integrated over a smaller domain nearer the shelf a stronger annual cycle appears at 26.5° N during 2004–2009. Even this stronger annual cycle at 26.5° N is not statistically significant from zero, however, based on the observed scatter at that location. Several years of additional data will be required to evaluate whether a significant annual cycle exists in the real ocean at 34.5° S.

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The long model record can also be used to determine the length of record needed to encapsulate the bulk of the expected variability at 34.5° S (at least to the extent that the model spectra reproduces the real ocean spectrum – comparison of the model spectrum and data spectrum from the pilot array, not shown, finds significant differences at periods of less than 100 days). A similar 1000-sample Monte Carlo-style calculation was made examining the variance of 1, 10, 18, 36, 60 and 120-month subsets as compared to the variance of the full 27-yr model record (Fig. 10). Not surprisingly as the record length of the Monte Carlo-style subsample increases the observed variance asymptotes to the full-record variance of 232 Sv² for the full 27-yr record (see the black solid line converging to the horizontal dotted line in Fig. 10). The full range of the 1000 subsamples (gray filled area in Fig. 10) and the standard deviation of the estimated variances (red cross-hatched area in Fig. 10) become smaller with increasing subsample record length. Assuming the spectral distribution of energy in the model closely approximates that of the real ocean, this suggests that 48–60 months of data will be required before the observed variance would approximate that of a longer-term (i.e. decadal to multi-decadal) record. Given the potential disagreement between the model and real ocean variability that is hinted at by the amplitude difference between model and reality, and the lack of any obvious suggestion of annual or semi-annual energy in the first 10.5 months of data, it is likely that this 4–5 yr requirement is a lower bound for the needed record length.

6 Conclusions

Just under a year of data from a pilot array of pressure-equipped inverted echo sounders (PIES) have been used to provide a first glimpse at the time varying absolute flow of the Deep Western Boundary Current (DWBC) at 34.5° S. The 10.5 months of PIES data illustrate a high degree of DWBC variability in this short time span, with a transport variance of roughly 230 Sv² (15 Sv standard deviation) when integrated from 800–4800 dbar (or the bottom) and from 51.5° W to 44.5° W. This high transport

variance is comparable to that found in the subtropical gyre in the North Atlantic at 26.5° N, which is somewhat surprising given the expected deep circulation pattern believed to exist in the South Atlantic at 34.5° S, where a significant fraction of the signal is expected to be on the eastern boundary (e.g., Warren and Speer, 1991; Speer et al., 1995). It is clear that the variability of the DWBC must be studied in the context of other features (e.g., Rossby waves, etc.) that exist in the basin.

Ultimately the goal of this pilot array is to build, with international collaborative projects, a trans-basin monitoring array for the Meridional Overturning Circulation along 34.5° S. The limited nature of historical/independent observations of the absolute deep transports at 34.5° S (e.g., Zenk et al., 1999) makes validation at this location difficult, however the technique has been carefully validated at other locations and comparison with output from a high-quality, high-resolution, ocean general circulation model at 34.5° S indicates that the observed variability from the pilot array is comparable, if slightly larger, than that predicted by the model. Future augmentation with additional instruments (e.g., Perez et al., 2011), and longer time series at these existing sites, will lead to better understanding of the DWBC variability at this location, and ultimately to the relationship between DWBC variability and those of the basin-wide Meridional Overturning Circulation.

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suggestions for improving an earlier draft of this manuscript, and their help is gratefully acknowledged. ARP acknowledges the support of grant CRN2076 from the Inter-American Institute for Global Change Research, which is supported by the US National Science Foundation (GEO-0452325). CSM, RCP and SLG also acknowledge support from AOML and the NOAA Climate Program Office.

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Table 1. Time-mean near-bottom meridional velocity from the OFES model averaged over the horizontal span between the indicated pairs of PIES/CPIES and over the deepest three layers above the model ocean bottom. Negative velocity indicates southward flow. Also shown are the real ocean depths at each of the actual PIES/CPIES sites and the model ocean depths at the nearest model grid points.

Site/span	Mean velocity	Real ocean depth	Model ocean depth
Site A		1360 m	1429 m
Span from A-to-B	-5.8 cm s^{-1}		
Site B		3535 m	3831 m
Span from B-to-C	-0.1 cm s^{-1}		
Site C		4540 m	4760 m
Span from C-to-D	$+0.2 \text{ cm s}^{-1}$		
Site D		4757 m	4760 m

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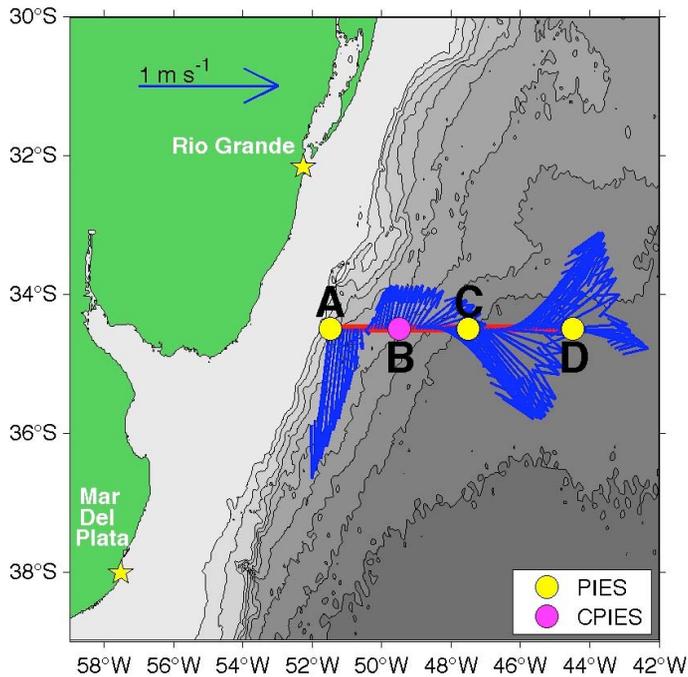


Fig. 1. Map indicating the location of the four PIES/CPIES making up the pilot array. Blue vectors indicate the water velocity at 21 m measured via shipboard acoustic Doppler current profiler on the Brazilian naval research vessel N. H. Cruzeiro do Sul in March 2009 during the array deployment cruise. Black letters indicate site names – instrument types are noted in legend.

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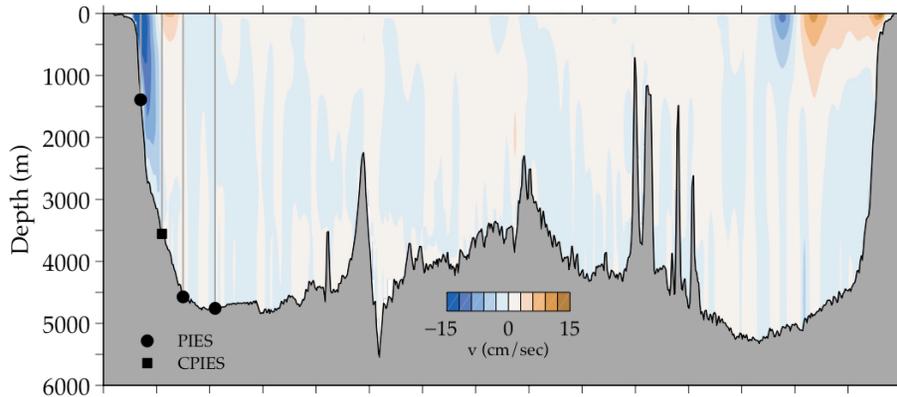


Fig. 2. Contour plot of the OFES model record-length (27-yr) mean meridional velocity, with oranges indicating northward flow and blues indicating southward flow. Also shown are the nominal locations of the PIES/CPIES deployed in the pilot array on the western boundary. Model (OCCAM) 1/10 degree bottom topography is shown in gray.

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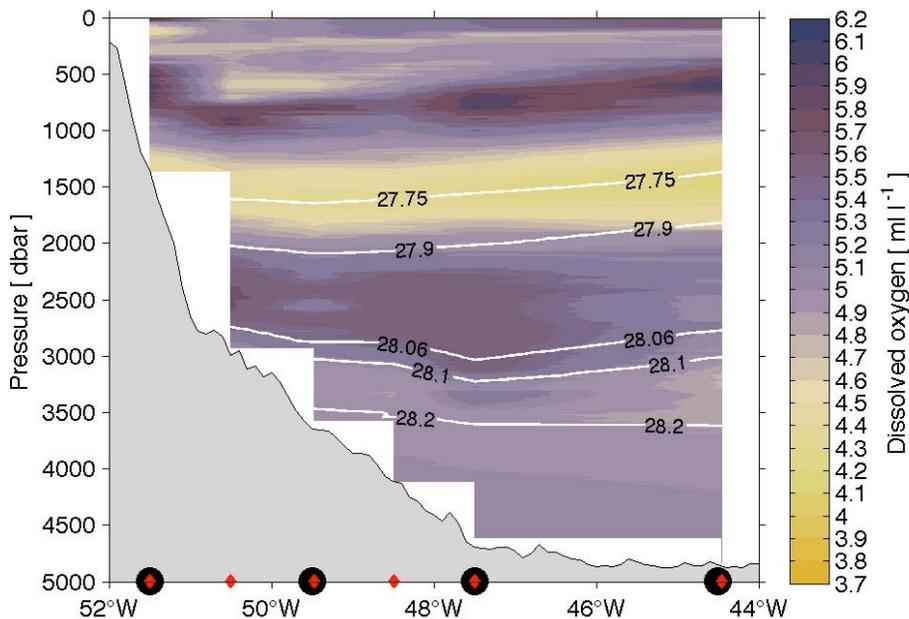


Fig. 3. Dissolved oxygen section collected along the PIES/CPIES line during 7–11 July 2010 on the Argentine research vessel Puerto Deseado. Red diamonds along bottom axis indicate locations of the CTD profiles, and black dots indicate the PIES/CPIES sites. White contours with labels indicate neutral density surfaces. Gray-shading indicates bottom topography from the Smith-Sandwell data set (Smith and Sandwell, 1997).

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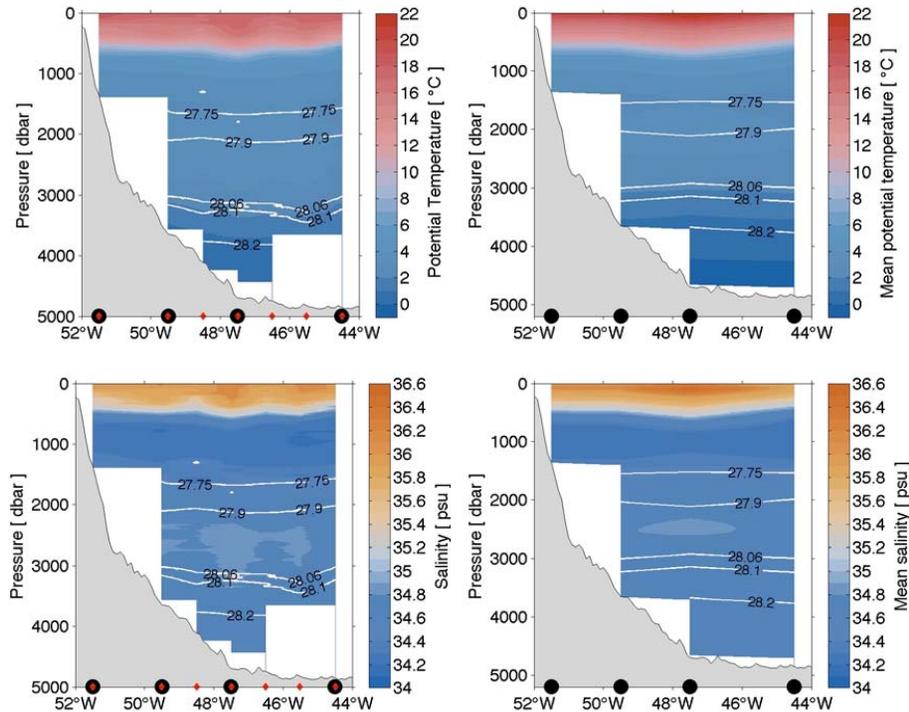


Fig. 4. Comparison of potential temperature and salinity measured during the 20–24 August 2009 CTD section (left, top and bottom, respectively) and the PIES-GEM estimated potential temperature and salinity averaged over the same five days (right, top and bottom, respectively). Conductivity measurements by the CTD were noisy due to a sensor problem; the data have been smoothed vertically to remove small artificial vertical structures. Gray shading indicates bottom topography. Large black dots on bottom axes indicate the PIES/CPIES locations, while small red diamonds indicate the location of the CTD casts. White contours with labels indicate neutral density surfaces.

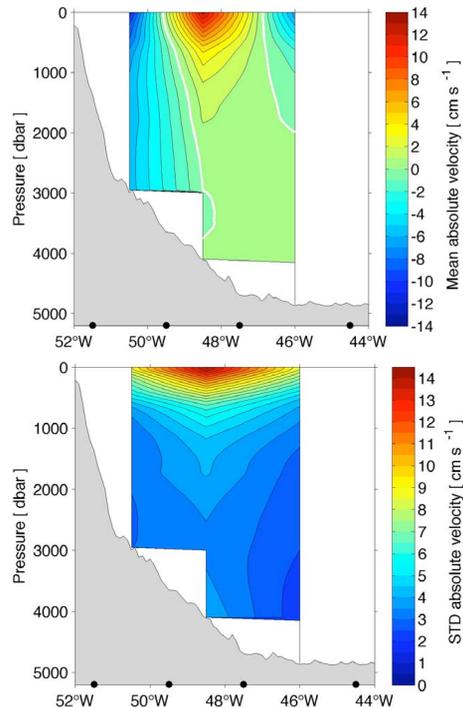


Fig. 5. Upper panel – time mean absolute meridional velocity determined between pairs of PIES/CPIES along the array. Lower panel – standard deviation of the time-varying absolute meridional velocity determined between the pairs of PIES/CPIES along the array. Black dots along lower axis denote locations of the PIES/CPIES.

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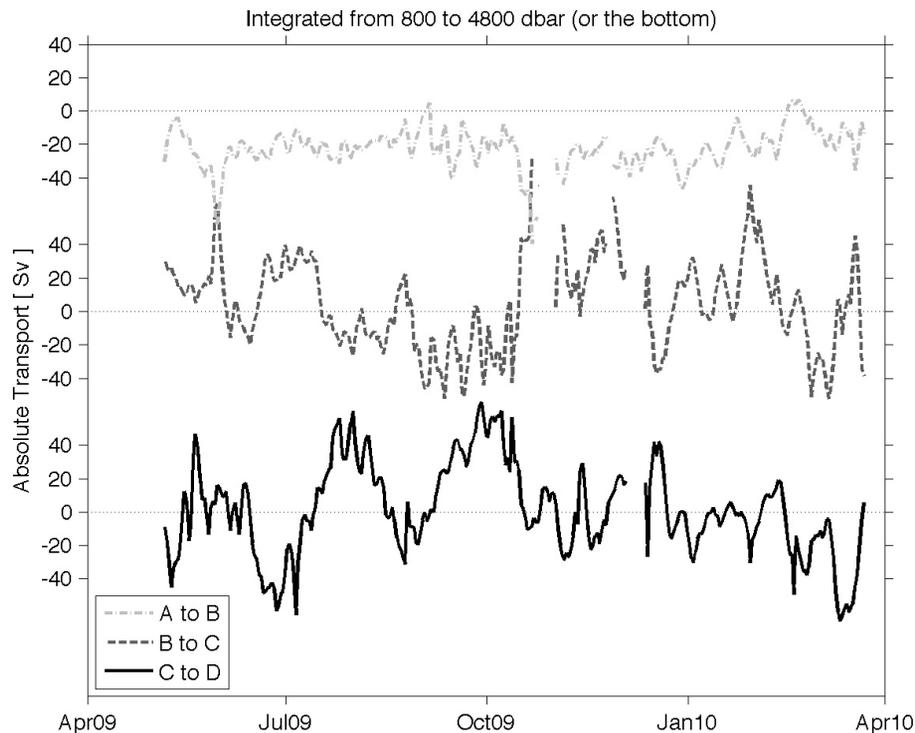


Fig. 6. Absolute transport integrated within the DWBC layer (800–4800 dbar) and between the indicated pairs of PIES/CPIES (see Fig. 1 for locations). Standard deviations of the three records are 12, 27, and 28 Sv for the A-to-B, B-to-C and C-to-D spans, respectively.

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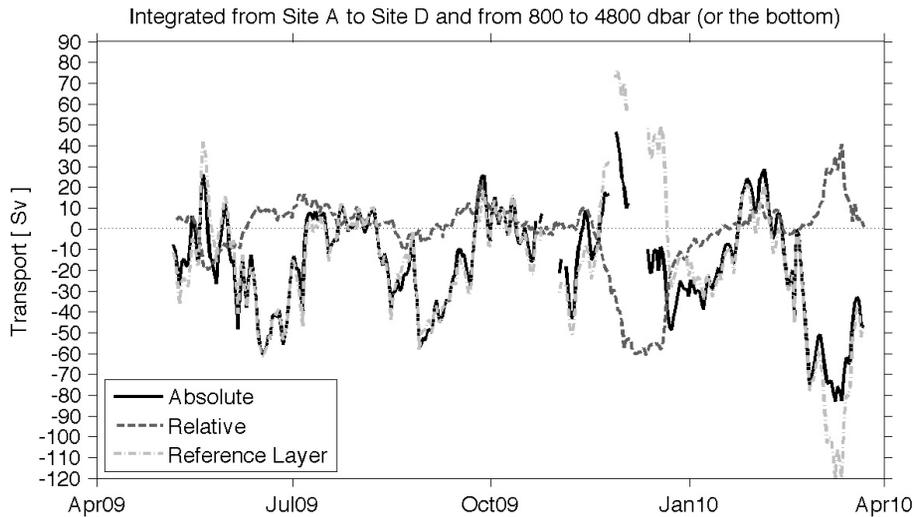


Fig. 7. Absolute transport (black line) integrated between sites A and D within the nominal DWBC layer (800–4800 dbar). Also shown are the components of the absolute transport associated with velocity relative to an assumed level of no motion at 800 dbar (dark gray dashed) and with the absolute velocity actually observed at the level of no motion reference layer (light gray dash-dot).

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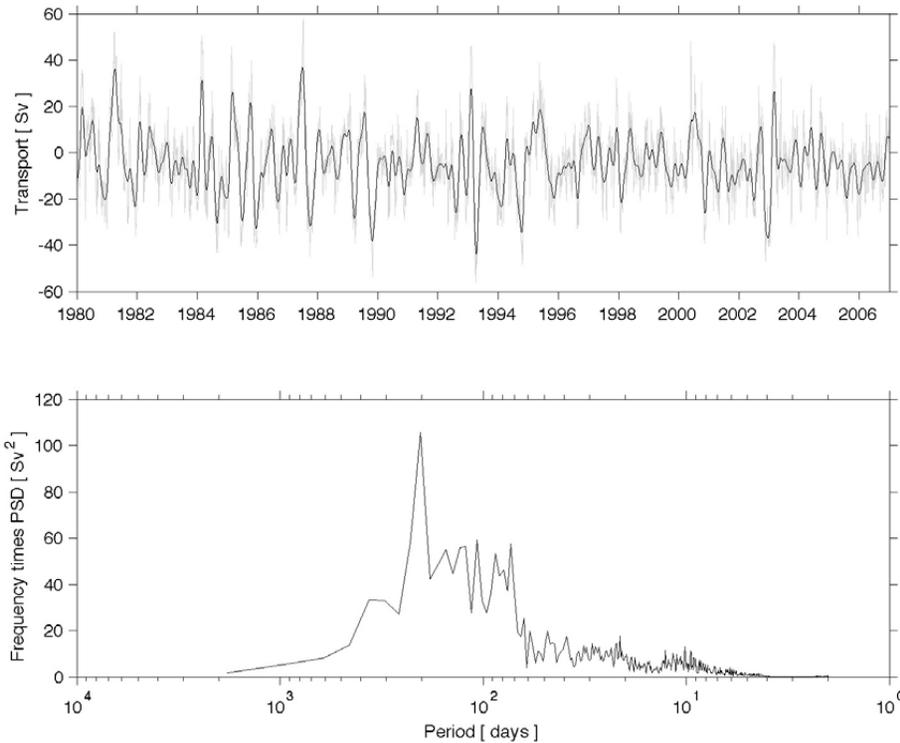


Fig. 8. Upper panel: DWBC transport from the OFES model integrated over the same span as the real observations. Gray line is the 3-day subsampled model output, while the black line is the 90-day low-pass filtered record. Lower panel: Spectra of the DWBC transport from OFES (using the 3-day subsampled, unfiltered data).

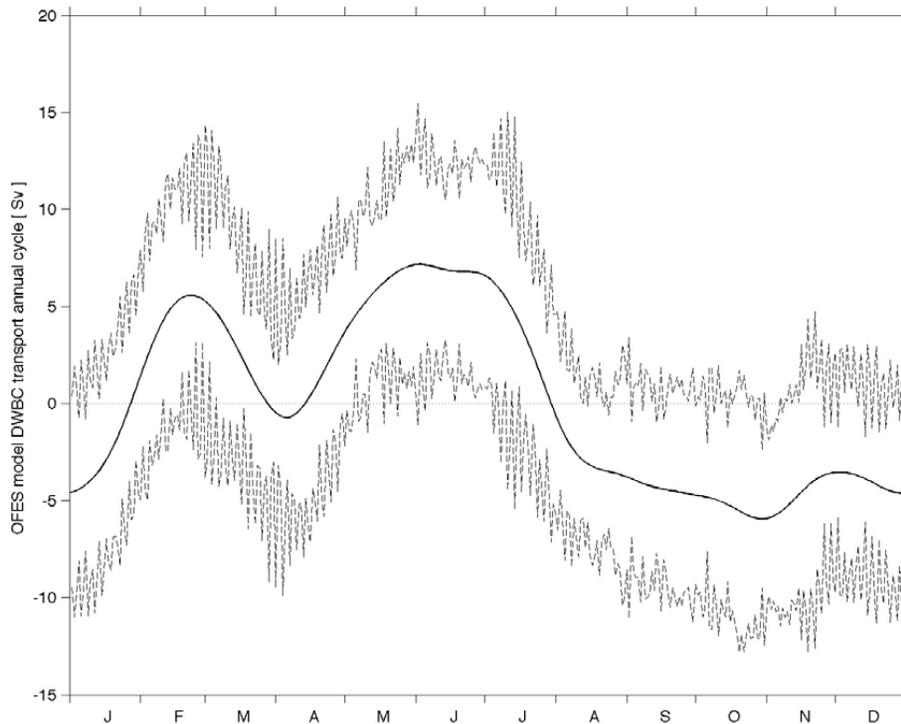


Fig. 9. Annual cycle of the DWBC transport calculated from the 27-yr of OFES model output. Annual cycle was determined as a daily climatology that was smoothed with a 60-day 2nd order Butterworth low pass filter. Dashed lines indicate plus and minus two standard errors (95 % confidence limits).

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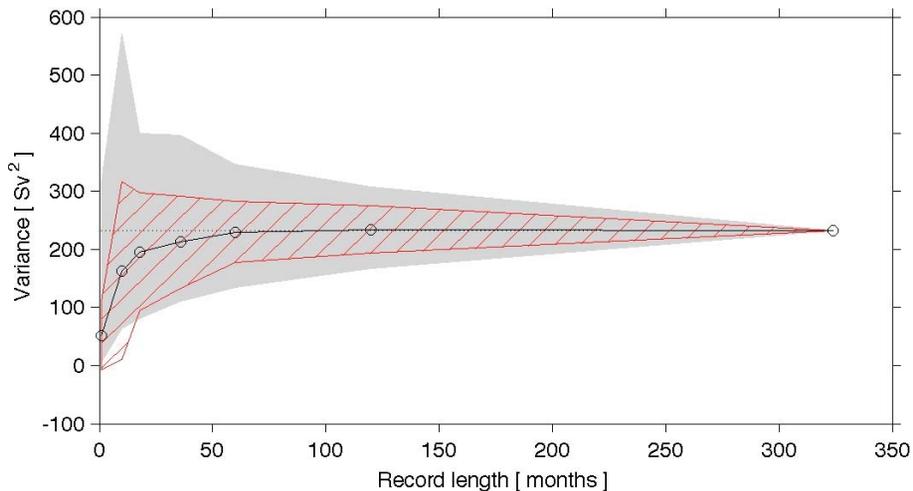


Fig. 10. Median DWBC transport variance determined via a Monte Carlo-style calculation using 1000 random selections of the indicated record lengths of the OFES time series shown in Fig. 8. Complete record variance is shown as the right-most point and the horizontal dotted line. Gray region illustrates the maximum and minimum range observed for all random subsamples of a given record length, while the red cross-hatched area indicates the median values plus and minus one standard deviation of the subsample variances.

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