

Abstract

The WOCE Hydrographic Program (WHP) and repeated hydrographic data were used to document overall property changes of the Antarctic Bottom Water (AABW) in the Australian-Antarctic Basin between the 1990s and 2000s. Strong cooling and freshening is observed on isopycnals for layers denser than $\gamma^n = 28.30$. Changes in average salinity and potential temperature below this isopycnal correspond to basin-wide warming of 1300 ± 200 TW and freshening of 24 ± 3 Gt yr⁻¹. While freshening can be explained by freshening of major source waters, i.e., the High Salinity Shelf Water (HSSW) of the Ross Sea and the dense shelf water formed in the Adélie and George V Land (AGVL) region, extensive warming of the AABW cannot be explained by warming of the source waters. A possible cause of warming of the AABW is a decrease in supply of the Ross Sea Bottom Water (RSBW). Hydrographic profiles between the Drygalski Trough of the Western Ross Sea and 150° E were analyzed in the context of a simple advective-diffusive model to assess the causes of the observed changes. The RSBW has also warmed by a larger amount than its source water (the HSSW). The model suggests that the warming of the RSBW observed between the 1970s and 2000s can be explained by a 21 ± 23 % reduction in transport of the RSBW and an enhancement of the vertical diffusion of heat as a result of a 30 ± 7 % weakening of the abyssal stratification. Freshening of the HSSW reduced the salinity and density stratification between the bottom water layer and overlying ambient water. Hence, freshening of the HSSW both directly freshened and indirectly warmed the RSBW by enhancing the vertical mixing. A simple box model suggest that changes in property and volume transport (decrease of 6.7 % is assumed between the year 1995 and 2005) of the RSBW can explain 51 ± 6 % of the warming and 84 ± 10 % of the freshening observed in the AABW. These facts demonstrate that changes in both property and volume transport of the RSBW have contributed to the warming and freshening of the AABW in the Australian-Antarctic Basin.

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1 Introduction

Abyssal layer of the world ocean is supplied by the Antarctic Bottom Water (AABW) which originates from cold dense water formed by the formation of large volumes of sea ice in the Antarctic coastal polynyas. Thermohaline circulation, or overturning circulation, which is driven by the formation of dense water, makes an important contribution to the transport and storage of heat, carbon and other properties that influence climate. Understanding of formation process of the AABW, and its sensitivity to change, is therefore needed to improve projection of future climate change. The AABW which occupies abyssal layers of the Australian-Antarctic Basin is supplied by dense waters exported from the Ross Sea (Ross Sea Bottom Water, RSBW) and from the Adélie and George V Land (AGVL) region (Adélie Land Bottom Water, ALBW) (Rintoul, 1998; Williams et al., 2008, 2010). (As the focus of this paper is on the abyssal waters of the Australian-Antarctic Basin, we use AABW here to refer to this particular variety of bottom water.) The dense shelf waters formed by heat loss and addition of brine during sea ice formation in coastal polynyas entrain the relatively warm Circumpolar Deep Water (CDW) as they descend the continental slope. Both polynyas in the AGVL region and Ross Sea are regions of intense sea ice formation (Tamura et al., 2008) and hence dense water production. The inflows of the ALBW and the RSBW make the Australian-Antarctic Basin one of the best-ventilated basin in the Southern Ocean, based on the concentration of chlorofluorocarbons (CFCs) (Orsi et al., 1999).

The nature and causes of changes in properties of the AABW in the Australian-Antarctic Basin have been clarified in several recent studies. Whitworth (2002) concluded that two distinct modes of the AABW had been produced in the basin, with the cold, fresh mode more prevalent in later observations. Aoki et al. (2005), from comparison of CTD profiles from the same location and the season, reported evident freshening trend at 140° E from at least the early 1970s. Using the sections spanning the entire basin, Rintoul (2007) showed that freshening of the AABW was observed throughout the basin. The distribution of CFC-12 indicated that changes are limited

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to ventilated layer that had been in contact with the atmosphere within 55 yr (Johnson et al., 2008), consistent with the time scale of changes reported by Aoki et al. (2005) and Rintoul (2007). Changes in the abyssal waters of much of the world ocean have been traced back to the AABW, with the largest changes seen close to the Antarctic continent (Purkey and Johnson, 2010). In particular, changes in the AABW of the Australian-Antarctic Basin have been identified as the source of changes in the deep North Pacific (e.g., Kawano et al., 2009; Masuda et al., 2010). Identifying the process driving the changes in the AABW properties is therefore an important step in understanding changes in the overturning circulation.

Changes in the AABW could reflect changes in the properties, volume transport or mixing ratios of the sources of the dens waters to the Australian-Antarctic Basin. For example, the wide-spread freshening observed in recent decades could reflect an increased supply of melt water from ice shelves and glacier tongues. The rapid melt of floating ice in the Amundsen-Bellingshausen Seas has been identified as the likely source of the strong freshening trend observed in the shelf waters of the Ross Sea (Jacobs et al., 2002; Zwally et al., 2005; Jacobs and Giulivi, 2010). The changes in temperature of the AABW have more complicated distribution and are more difficult to explain. The AABW warmed over much of the basin between the years 1936–1993 and 1994–1996 (Whitworth, 2002), but cooled in the coastal region between 140 and 150° E from the year 1950 to 2001 (Jacobs, 2004, 2006) and at 140° E between the 1970s and 2002 (Aoki et al., 2005). These changes in temperature with space and time suggest that changes in the relative contribution of different source waters may have also played a role in observed changes in the AABW.

Several studies suggest that changes in the RSBW, and in its source waters the HSSW, are likely to contribute changes observed in the Australian-Antarctic Basin. Jacobs (2006) and Rintoul (2007) showed consistent freshening of the RSBW between the Ross Sea outflow and 150° E. Freshening of shelf waters in the Ross Sea have been attributed in part to inflow of glacial melt water from the east (Jacobs et al., 2002; Jacobs and Giulivi, 2010). Changes in sea ice formation likely also have played a role in

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recent freshening. Tamura et al. (2008) inferred a 30 % reduction in the volume of sea ice formed in the Ross Ice Shelf Polynya region from the end of the 1990s to the 2000s. Rivalo et al. (2010) reported that the export of dense water from the Western Ross Sea was 45 % lower in 2001 than in 1997. Therefore, change in both the properties and volume transport of the RSBW may have contributed to change in the AABW observed downstream.

In this study, we describe changes in the AABW properties based on CTD profiles of the WOCE Hydrographic Program (WHP) occupation in the 1990s and later repeat occupation in 2000s. Historical observations are used to track changes in the RSBW. A simple advective-diffusive model of the RSBW out flow is then developed to determine the contribution of changes in transport and mixing to the changes observed in the RSBW. Finally, influences of changes in property and volume transport of the RSBW on property of the AABW observe downstream are assessed.

2 Changes in property of the AABW in the Australian-Antarctic Basin

2.1 Changes in property of the AABW

A schematic illustrating the sources and circulation of the AABW in the Australian-Antarctic Basin is shown in Fig. 1. Dense waters formed in the AGVL region and the Ross Sea, after entraining the CDW as they descend the slope, supply the AABW, which flows westward over the continental slope and rise of Antarctica to the Kerguelen Plateau. The AABW turns north along the plateau until it joins the lower part of the Antarctic Circumpolar Current (ACC) to flow to the east (Orsi et al., 1999). The AABW exits the basin through the Princess Elisabeth Trough (PET) and through the Australian-Antarctic Discordance (AAD). The WHP sections (I8S, I9S, and SR3) cross the set of boundary currents transporting the AABW away from the source regions. Each of these sections was repeated in the 2000s (station locations are shown in Fig. 1 and details are given in Table 1). In this section, we describe changes in bottom water properties based on these data.

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The SR3 section crosses the AABW path just downstream of the inflow of the ALBW. Because this line is close to the source, near-bottom temperature and salinity vary with season and with location (e.g., Fukamachi et al., 2000). To minimize aliasing of seasonal and spatial variability, we focus on near-repeats of the same stations and in the same season. For the SR3, we supplement the October 2001 section with stations south of 61° S occupied in 2002 and 2003 in the same season as the original WHP line, to minimize seasonal aliasing (hereafter the combined profiles from 2001–2003 are simply referred to as 2001). The two repeats of the I9S were in the same season. At the I8S in the western side of the basin, the repeats were within 2 to 3 months of each other, and as the section is distant from the source regions we expect seasonal variability to be small. (No clear evidence of seasonal variability was observed in a two-year mooring array close to the I8S section, Fukamachi et al., 2010).

Changes in the AABW are illustrated by comparing the meridional variation of potential temperature, salinity, neutral density and apparent oxygen utilization (AOU) averaged over a 100 m-thick layer at the bottom for the WHP and repeat sections (Fig. 2). At the SR3, the potential temperature is lower in 2001 and 2008 than in the WHP south of 62° S. However, the lowest potential temperature (−0.72 °C) is in 2001 and the potential temperature in 2008 (−0.64 °C) is comparable to that of the WHP. Such a non-monotonic change suggests interannual variability in potential temperature. Cooling is also observed near 63° S at the I9S (−0.06 °C). However, warming is found further downstream along the entire I8S section and north of 61° S along the I9S (by up to 0.1 °C). The reversal of sign in the potential temperature changes is also observed in AOU: cooling is associated with a decrease in AOU on the I9S south of 62° S, and warming is associated with an increase in AOU at the I8S and north of 62° S on the I9S (Fig. 2b). While the changes in salinity do not reverse in sign, the freshening trend is intensified closer to the source regions (near the SR3 and at higher latitudes). In addition, the bottom layer is fresher in 2001 than in 2008 at SR3 (the minimum salinity south of 62° S is lower by 0.01). Hence, salinity also varies interannually. Near the sources of dense water (south of 62° S on the SR3 and near 63° S on the I9S), there

is no clear change of density because cooling and freshening compensate each other. Further downstream (at the I8S and north of 61° S on the I9S), freshening and warming produce a decrease in density ($\sim 0.02 \text{ kg m}^{-3}$).

Changes in the potential temperature-salinity diagram are illustrated in Fig. 3, following Johnson et al. (2008). Cooling and freshening on isopycnals is evident for water denser than $\gamma^n = 28.30$, indicating basin-wide changes in the AABW. Cooling and freshening on isopycnals reflects warming and/or freshening of water masses where the stratification is such that warmer, more saline waters overly cooler, fresher waters (i.e., density ratio > 1 , Bindoff and McDougall, 1994). Thus, the changes in the AABW properties likely indicate warming or freshening of the sources of dense water.

Potential temperature and salinity differences on pressure surfaces are shown in Figs. 4 and 5. At the SR3, the strongest cooling is observed at the sea floor between 63° S and 64° S (Fig. 4d, e). At the I9S, warming dominates below $\gamma^n = 28.30$, although weak cooling is detected over the slope south of 62° S (Fig. 4c). Freshening dominates everywhere below $\gamma^n = 28.30$, with the strongest freshening at the sea floor (Fig. 5c). At the I8S, warming (Fig. 4a, b) and moderate freshening (Fig. 5a, b) are observed below $\gamma^n = 28.30$.

In summary, warming increases downstream and the freshening increases upstream along the pathway of the AABW. To obtain a rough estimate of the heat and freshwater fluxes needed to explain the observed patterns of change, we interpolated the changes observed at each section to the entire basin, using a Gaussian weighting function with an e-folding scale of 300 km and an influence radius of 1000 km. The overall warming of the AABW requires an input of $1300 \pm 200 \text{ TW}$ of heat (or $0.37 \pm 0.05 \text{ W m}^{-2}$) and 24 Gt yr^{-1} of freshwater (Table 2).

2.2 Possible causes of the observed changes

One possible cause of the observed freshening is a decrease in salinity of the dense shelf waters, as observed in the Ross Sea ($-0.03 \text{ decade}^{-1}$ since the 1960s, Jacobs and Giulivi, 2010). No equivalent time series exists to assess changes in the salinity of

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dense shelf waters formed in the AGVL region. On the other hand, it is not straightforward to explain the warming of the AABW. The warming of the HSSW in the Ross Sea is trivial ($2 \times 10^{-3} \text{ }^\circ\text{C decade}^{-1}$) and of the size expected from the change in freezing point temperature induced by the rapid freshening (Jacobs and Giulivi, 2010). Warming of the dense source waters therefore cannot explain the warming of the AABW. Other possible causes of warming include geothermal heat flux, horizontal or vertical diffusion, and reduced supply of cold dense shelf waters.

Among these candidates, the average geothermal heat flux for depths of 3500 to 4500 m is roughly $0.05\text{--}0.1 \text{ W m}^{-2}$ (Stein and Stein, 1992). Even a 50 % increase would explain less than 15 % of the warming of the AABW and there is no reason to expect the geothermal heat flux to increase rapidly on decadal time-scales.

Horizontal diffusion is expected to carry a heat flux of about 0.01 W m^{-2} , using a basin-averaged horizontal temperature gradient at the depth of the $\gamma^n = 28.30$ isopycnal ($8.9 \times 10^{-7} \text{ }^\circ\text{C m}^{-1}$) and a horizontal diffusivity of $2 \text{ m}^2 \text{ s}^{-1}$ (e.g., Ledwell, 1999). Hence horizontal diffusion cannot explain the observed warming. Polzin and Firing (1997) estimated a vertical diffusivity at 55° S on the I8S averaged below 1000 m of $4.4 \times 10^{-4} \text{ m}^2 \text{ s}^{-1}$. The basin-averaged vertical gradient of potential temperature at the depth of the $\gamma^n = 28.30$ isopycnal is $3.8 \times 10^{-4} \text{ }^\circ\text{C m}^{-1}$, implying a vertical diffusive heat flux of 0.67 W m^{-2} . Changes in the basin-averaged density stratification, which can affect both horizontal and vertical diffusion, were less than 10 % during the observational period. An increase in vertical diffusion of 10 % would explain only 19 % of the warming over the entire AABW layer. Hence, none of these processes can account for sufficient change in heat flux to explain the observed warming.

The spatial distribution of the observed changes is also an important factor. While the AABW warmed over most of the basin, cooling was observed near the source region. The fact that localized cooling is accompanied by a decrease in AOU and by significant freshening suggests a larger input of the cold, fresh source from the AGVL region relative to the input of relatively warm and salty inflow from the Ross Sea.

The rough order of magnitude estimates of the contribution of various heat sources and the spatial distribution of changes in temperature and salinity support the hypothesis that a reduction in the inflow of the RSBW may have contributed to the changes observed in the basin. To examine this possibility in more detail, we consider an advective-diffusive balance for the waters between the Western Ross Sea and 150° E.

3 Changes in water masses property and advection-diffusion processes of the RSBW, and its impact on the AABW

3.1 Changes in property of the RSBW

Summer (January to March) hydrographic profiles along the approximate flow path of the RSBW in the region between the Drygalski Trough of the Western Ross Sea and 150° E were collected from World Ocean Database 2009 (Boyer et al., 2009). A total of 256 profiles reach within 50 m of the bottom, spanning the period from 1969 to 2004. Property changes of the RSBW are shown in Fig. 6. The RSBW freshened consistently from the 1970s and the bottom-intensified salinity maximum is drastically attenuated in the 2000s throughout the study area around 150° E (Fig. 6a), 160° E (Fig. 6b), and in the vicinity of the source region (Fig. 6c). Warming is observed except near the source region. At 160° E, AOU averaged over the 100 m-thick layer at the bottom slightly increased in the 2000s ($118.1 \mu\text{mol kg}^{-1}$) compared to that in 1970s ($114.1 \mu\text{mol kg}^{-1}$; not shown).

Freshening of the RSBW can be explained by freshening of the HSSW, but warming of the RSBW cannot be explained by the observed warming of the HSSW ($2 \times 10^{-3} \text{ }^\circ\text{C decade}^{-1}$). Since decreases in volume transport of the HSSW and hence the RSBW reduces the supply of both negative heat flux and oxygen rich water throughout its flow path, they are consistent with the property changes between the Drygalski Trough and 150° E. In the next subsection, we quantitatively estimate the change in volume transport of the RSBW by using a simple advection-diffusion model.

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3.2 A model of advection-diffusion process of the RSBW

The cold and saline HSSW flows out from the Drygalski Trough of the Western Ross Sea. The outflow sinks down the slope, entraining the ambient warm and relatively fresh (compared to the HSSW) CDW to form the RSBW. The RSBW is advected westward to 150° E with further warming and freshening induced by mixing with the ambient water. Hereafter, the dense water originating from the HSSW and sinking down the slope is referred to as the Modified Shelf Water (MSW) following Orsi and Wiederwohl (2009). Based on observations in the 1970s and 2000s, we will estimate the change in volume transport of the RSBW by analyzing changes in its advection-diffusion processes.

The governing diffusion equation of potential temperature of the RSBW can be written as follows.

$$\theta_t + u\theta_x + v\theta_y + w\theta_z = \frac{\partial}{\partial x}Kh\theta_x + \frac{\partial}{\partial y}Kh\theta_y + \frac{\partial}{\partial z}Kz\theta_z,$$

where x is flow path coordinate (downstream positive), y is the direction orthogonal to the flow path, and z is vertical axis (upward positive). Kz and Kh are vertical and horizontal diffusivities, respectively, and subscripts indicate the first order derivatives.

Starting from this general formula, we will simplify this equation based on the settings and approximations applicable in this case. As shown in Fig. 7, flow paths of the RSBW are bounded by the continental slope and sea floor in the direction of y and z , respectively, and fluxes through the two boundaries are zero. Hence, horizontal diffusion term of y axis and vertical diffusion terms are replaced by fluxes from the other side divided by the width L_0 and thickness H of the RSBW, respectively.

$$\theta_t + u\theta_x + v\theta_y + w\theta_z = \frac{\partial}{\partial x}Kh\theta_x + \frac{1}{L_0}Kh\theta_y + \frac{1}{H}Kz\theta_z. \quad (1)$$

The third and fourth terms in the left hand side can be neglected since velocity orthogonal to the flow path and vertical velocity are negligible. The first term on the right

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hand side can be neglected because the horizontal gradient in the x -direction is negligible compared to gradients in the y and z directions. Equation (1) can therefore be simplified to:

$$\theta_t + u\theta_x = \frac{1}{L_0}Kh\theta_y + \frac{1}{H}Kz\theta_z. \quad (2)$$

Here, we will consider relative importance between the first and second terms in the right hand side. From Fig. 7, width L_0 and thickness H of the RSBW can be estimated as 200 km and 200 m respectively and horizontal and vertical gradients were estimated as $\theta_y \approx 2.6 \times 10^{-6} \text{ }^\circ\text{C m}^{-1}$ and $\theta_z \approx 7.0 \times 10^{-4} \text{ }^\circ\text{C m}^{-1}$, respectively. No information on vertical diffusivity in the study area is available (Muench et al. (2009) estimated it in the vicinity of the source region but not for the remaining area. We will discuss this point later). However, Kunze et al. (2006) estimated average vertical diffusivity in the Southern Ocean by using strain profiles from CTD profiles and vertical velocity shear from LADCP. From Fig. 16 of Kunze et al. (2006), vertical diffusivity of near bottom can be the orders of 10^{-5} – $10^{-4} \text{ m}^2 \text{ s}^{-1}$. By adopting Osborn's (1980) model, vertical diffusivity Kz can be written in the following form of

$$Kz = 0.2 \frac{\epsilon}{N^2} \quad (3)$$

where 0.2 is mixing efficiency, ϵ is energy dissipation rate, and N is Brunt–Väisälä frequency. ϵ is expected to be constant in time scale of our interest since dynamical background that drives turbulence remains unchanged. N^2 can be calculated from vertical profiles of observed data in following form of

$$N^2 = g(\beta S_z - \alpha \theta_z); \quad g = -9.8 \text{ m s}^{-2} \quad (4)$$

Though absolute value of Kz is unknown due to the lack in observation of ϵ , rate of temporal change of Kz can be estimated by calculating N^2 for each period. Here, by using vertical diffusivity of Kunze et al. (2006), vertical flux takes value of 3.5×10^{-11} to $3.5 \times 10^{-10} \text{ }^\circ\text{C s}^{-1}$. By using $Kh = 2 \text{ m}^2 \text{ s}^{-1}$ (e.g., Ledwell, 1999), horizontal fluxes

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takes value of $2.6 \times 10^{-11} \text{ }^\circ\text{C s}^{-1}$. Hence, the vertical diffusion term is likely one order larger than horizontal diffusion term. Even if the vertical diffusion does not dominate the horizontal diffusion, both vertical and horizontal diffusion are expected to act on the property of the RSBW in the same sense since the RSBW similarly mixes with the water of the CDW origin for both vertical and horizontal diffusion. Thus, by neglecting horizontal diffusion term, we obtain formula as

$$\theta_t + u\theta_x = \frac{1}{H}Kz\theta_z.$$

Now we will consider the uniform current ($u > 0$) throughout the flow path of the RSBW. There is no direct observation of u , but Fukamachi et al. (2000) measured current speed of 14.1 to 19.8 cm s^{-1} at slightly downstream (140° E). Pure downstream current speed could be somewhat smaller than the above value since current direction was not observed because the mooring was too close to the magnetic pole. But, if we assume the above value for downstream current speed, the time scale that advection carries the signal throughout the flow path of the 1000 km is 58 to 82 days. Hence, on the time scale we consider, temperature of this system will swiftly adjusted to the temperature induced at $x = 0$. Hence, for the case of multi-decadal variability, we obtain a formula as

$$\frac{\theta_z}{\theta_x} = \frac{uH}{Kz} \quad (5)$$

for each period of interest.

Equation (5) implies that the ratio of vertical gradient to horizontal gradient is determined by ratio of volume transport to vertical diffusivity. Horizontal velocity u determines residence time in the study area and, thus, determines the time scale which the RSBW is affected by vertical diffusion. Thickness H determines response of potential temperature of the RSBW against vertical diffusion. Both u and H proportionally contribute to maintain water mass property of the RSBW against the vertical diffusion. Hence, when volume transport uH takes large (small) value, horizontal gradient θ_x

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component and fluctuation which is assumed to increase in proportion to time, i.e., $V_{tr} = \overline{V_{tr}} + \delta V_{tr}t$ and $\Delta\theta = \overline{\Delta\theta} + \delta\Delta\theta t$. Substitution of these into Eq. (8) gives,

$$-\frac{1}{V} \left(\overline{V_{tr}}\overline{\Delta\theta} + \overline{V_{tr}}\delta\theta t + \overline{\Delta\theta}\delta V_{tr}t + \delta\theta\delta V_{tr}t^2 \right) \quad (9)$$

The first term in the bracket is steady state component which is assumed to be balanced with steady state heat flux in the basin due to vertical and horizontal diffusion and geothermal heat. Here we are interested in the fluctuation components which deviate from the balance, and thus, we neglect the first term and the fourth infinitesimal term. By integrating the above equation with respect to time, warming of the AABW due to the RSBW influence ($\Delta\theta_{RSBW}$) by temporal changes in property and volume transport of the RSBW can be considered. Corresponding freshening of the AABW due to the RSBW influence (ΔS_{RSBW}) is also given by the same procedure. For the integrating time T , we will have

$$\begin{aligned} \Delta\theta_{RSBW} &= -\frac{1}{V} \left(\overline{V_{tr}}\delta\theta + \overline{\Delta\theta}\delta V_{tr} \right) \frac{T^2}{2} \\ \Delta S_{RSBW} &= -\frac{1}{V} \left(\overline{V_{tr}}\delta S + \overline{\Delta S}\delta V_{tr} \right) \frac{T^2}{2} \end{aligned} \quad (10)$$

The first term in the right hand side of Eq. (10) is the component due to the changes in property of the RSBW and the second term is the component due to the decrease in volume transport.

Contributions of temperature and salinity changes from the RSBW are then estimated using Eq. (10) for the period of 10 yr from 1995 to 2005, following the observational period in the Australian-Antarctic Basin. The trend in potential temperature and salinity of the RSBW is evaluated by the regression line for bottom water property averaged over 100 m-thick layer at bottom around 150° E (Fig. 10).

Potential temperature and salinity of the AABW in the year 1995 and 2005 are calculated by basin average estimated by using all the data shown in Table 1 and the trends

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shown in Table 2a, respectively. Above average are estimated by the same interpolation procedure as that in the Sect. 2.1. Then, $\delta\theta$ and δS are estimated by changing rate of difference between the RSBW and the AABW as

$$\delta\theta = \frac{\Delta\theta_{2005} - \Delta\theta_{1995}}{T}, \quad \delta S = \frac{\Delta S_{2005} - \Delta S_{1995}}{T}.$$

Steady state component of volume transport \overline{V}_{tr} of the RSBW is required in this formulation. Temporal and spatial mean potential temperature (-0.36°C) and salinity (34.70) around 150°E (Fig. 10) gives approximate mixing ratio of the HSSW and the LCDW as 1 : 6. This implies that the HSSW exported from the Drygalski Trough entrained the LCDW whose volume is 6 times larger than the HSSW during advection. Then, adopting the export value of the HSSW from the Drygalski Trough (0.98 Sv) estimated by Whitworth and Orsi (2006), \overline{V}_{tr} at 150°E is assumed to be 7 Sv.

With the above \overline{V}_{tr} and observed $\delta\theta$ and δS , the change in property of the RSBW can explain $44 \pm 5\%$ and $63 \pm 8\%$ of overall warming and freshening of the AABW. The change rate in volume transport δV_{tr} is the largest unknown in Eq. (10). Hence, we consider possible two scenarios. For the scenario 1, we assume that the decrease of volume transport of the RSBW from 1995 to 2005 is 6.7%. This value is based on suggested decrease in volume transport of the RSBW ($\approx 20\%/30\text{yr}$) derived in the Sect. 3.2. For the scenario 2, we assume that the decrease of volume transport of 20%, which reflects the rapid decrease in sea ice production pointed out by Tamura et al. (2008).

Hereafter we evaluate $\Delta\theta_{\text{RSBW}}$ and ΔS_{RSBW} using Eq. (10) according to two scenarios for decrease in volume transport (Table 4). For the scenario 1, $\Delta\theta_{\text{RSBW}}$ can explain $51 \pm 6\%$ of the overall warming of the AABW. Contribution of the change in property (first term) by far exceeds that of the change in volume transport (second term). ΔS_{RSBW} can explain $84 \pm 10\%$ of the overall freshening of the AABW and similarly as in the case of the $\Delta\theta_{\text{RSBW}}$, contribution of the first term by far exceeds that of the second term. For the scenario 2, the second term largely increased and $\Delta\theta_{\text{RSBW}}$

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and ΔS_{RSBW} take value $63 \pm 8\%$ and $131 \pm 16\%$ of the overall warming and freshening of the AABW, respectively. In this case, the second term is comparable to the first term as for freshening, while it is approximately only a half of the first term as for warming.

Hence, the change in properties of the RSBW can explain nearly a half of the observed change of the AABW, and the decrease in volume transport of the RSBW can also be significant especially with the larger decrease with time.

4 Discussion

Property changes observed in the AABW require a heat and freshwater input of 1300 TW and 24 Gt yr^{-1} , respectively (the Sect. 2). The inferred rate of warming ($0.37 \pm 0.05 \text{ W m}^{-2}$) of the AABW that supplies the abyssal basin deeper than 3000 m is comparable to that of the world ocean warming (0.20 W m^{-2}) between 1955 and 1998 (Levitus et al., 2005), a large part of which occurs above 700 m depth. The observed freshening is about half of the freshening rate on the shelf and over the gyre reported by Jacobs et al. (2002). Errors are estimated from the multiple occupations of SR3, as only one repeat of the other sections is available.

Changes in the RSBW are inferred to be one of the major factors driving the changes observed in the AABW (the Sect. 3). Both enhanced vertical diffusion as a result of weakening stratification and a decrease in volume transport of the RSBW are likely to contribute to the warming and increase in AOU of the RSBW, and to the changes in AABW observed downstream. However, this conclusion rests on numerous assumptions. Here we discuss the robustness and consistency of the results.

Horizontal diffusion was assumed to be small (Eq. 2). The validity of this approach can be tested by confirming consistency between the observed vertical diffusivity (Muench et al., 2009) and that we adopted in this study. Based on Eq. (6), K_z can be expressed as:

$$K_z = \frac{uH}{\theta_z/\theta_x} = \frac{uH}{S_z/S_x}$$

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trends in the properties of the RSBW, used in the evaluation of the influence of the RSBW changes on the AABW, are based on observations between the 1970s and 1990s since no data were available from after 1996 near 150° E.

Long-term changes in the dense shelf water formed in the AGVL region may also be important, but could not be estimated because of the lack of long-term observations in the area. Furthermore, it is pointed out from model study that the calving of the Mertz Glacier tongue in this region, which occurred on 12 to 13 February 2010, can reduce export of the Dense Shelf Water by 23% with an unignorable decrease in density of exported water (Kusahara et al., 2010). Given that this type of event occur on decadal to centenary time-scales, influence on both the Dense Shelf Water and the AABW can be extensive. Hence, it is highly desirable to continue ongoing measurements programs in this area.

As the steady state volume transport of the RSBW is a key parameter that determines the influence of the RSBW change on the AABW, better observations of flow rates would help to clarify the cause of changes observed in the AABW. Further observations along the flow path from the Ross Sea to the Australian-Antarctic Basin are needed to monitor the change and estimate the impact on the AABW.

Appendix A

Application of the OMP analysis

In this study we used Optimum Multiparameter (OMP) analysis to discern profiles which are not influenced by the HSSW in vicinity of the slope by defining water type of the HSSW and the LCDW. The OMP analysis is introduced by Tomczak (1981), and Tomczak and Large (1989) to estimate spatial distribution of water masses and their mixing process. As for potential temperature of the LCDW, we adopt 0°C for both periods of the 1970s and 2000s and corresponding salinity of 34.70 and 34.685 are given from θ - S diagram. As for potential temperature and salinity of the HSSW, those of each

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period are adopted from regression lines shown in Fig. 3 of Jacobs and Giulivi (2010), which presented long term trend of both potential temperature and salinity of the HSSW (as for the 1970s: $\theta = -1.89^\circ\text{C}$, $S = 34.83$, as for the 2000s: $\theta = -1.87^\circ\text{C}$, $S = 34.74$).

Acknowledgements. This work was supported by the fund from Grant-in-Aids for Scientific Research (20221001 and 21310002) of the Japanese Ministry of Education, Culture, Sports, Science and Technology.

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Table 2a. Basin and section averaged trends in property of the AABW in the Australian-Antarctic Basin. Total volume of layers below isopycnal of $\gamma^n = 28.30$, mean potential temperature trend, required heat flux, mean salinity trend, required freshwater input in the Australian-Antarctic Basin.

	Australian-Antarctic Basin
Volume (km ³)	2.54×10^6
$\Delta\theta$ (°C decade ⁻¹)	$4.09 \pm 0.53 \times 10^{-2}$
Heat input	1322 ± 172 (TW)/ 0.37 ± 0.05 (W m ⁻²)
ΔS (decade ⁻¹)	$-2.67 \pm 0.42 \times 10^{-3}$
Fresh water input (Gt yr ⁻¹)	24.0 ± 3.2

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Table 2b. Basin and section averaged trends in property of the AABW in the Australian-Antarctic Basin. Section averaged thickness of layers below isopycnals of $\gamma^{\theta} = 28.30$, potential temperature trends, required heat fluxes, salinity trends, required freshwater inputs for respective sections. Error ranges for the SR3 were derived from 95 % confidence limit of linear regression coefficient.

	I8S (40–60° S)	I8S (60–67° S)	I9S	SR3
Thickness (m)	761	840	838	695
$\Delta\theta$ ($^{\circ}\text{C decade}^{-1}$)	7.65×10^{-2}	4.22×10^{-2}	3.72×10^{-2}	$2.14 \pm 2.41 \times 10^{-2}$
Heat flux (W m^{-2})	0.75	0.46	0.40	0.19 ± 0.22
ΔS (decade^{-1})	-1.15×10^{-3}	-2.82×10^{-3}	-3.14×10^{-3}	$-4.22 \pm 2.24 \times 10^{-3}$
Fresh water input ($\text{mm m}^{-2} \text{yr}^{-1}$)	3.5	7.5	8.3	9.0 ± 4.6

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Table 3. Mean vertical salinity gradients, horizontal salinity gradients, their ratios, N^2 and uH^* for the RSBW in the 1970s and 2000s from region between 150 km and 900 km downstream from the Drygalski Trough.

	1970s	2000s
S_z (m^{-1})	$-8.94 \pm 0.71 \times 10^{-5}$	$-3.71 \pm 0.24 \times 10^{-5}$
S_x (m^{-1})	$-3.86 \pm 0.59 \times 10^{-8}$	$-2.90 \pm 0.62 \times 10^{-8}$
S_z/S_x	2300 ± 400	1300 ± 300
N^2 (s^{-2})	$1.91 \pm 0.17 \times 10^{-6}$	$1.34 \pm 0.05 \times 10^{-6}$
uH^* (s^2)	$12.1 \pm 2.3 \times 10^8$	$9.5 \pm 2.2 \times 10^8$

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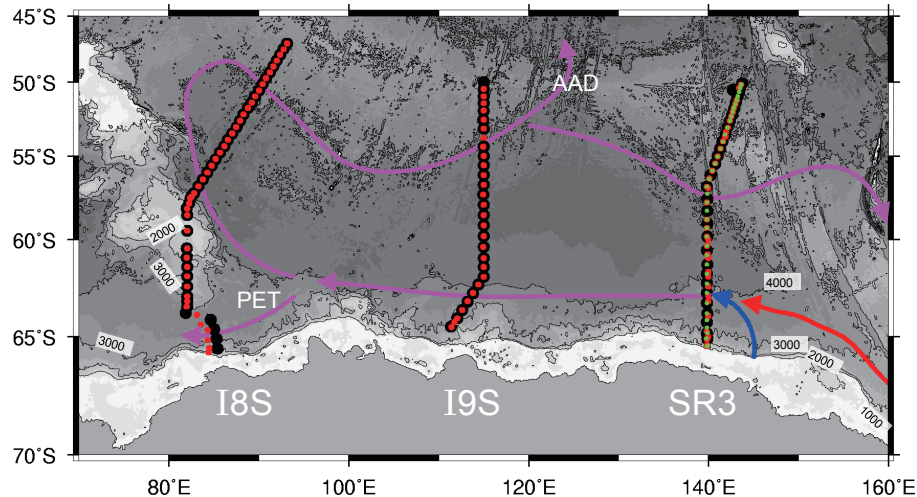


Fig. 1. Locations of the WOCE WHP (black circles) and their repeat (red and green circles) meridional hydrographic sections crossing the Australian-Antarctic Basin with schematic view of spreading path of the AABW. Red circles on the SR3 show locations of repeat occupations in the year 2001 (50–61° S), 2002 (61–66° S), 2003 (61–66° S) and green circles show those in 2008, respectively. All the red circles on both the I9S and I8S show repeat occupations in the 2000s. Red and blue arrows indicate the RSBW from the Ross Sea and ALBW from the AGVL region respectively and the AABW is indicated by purple arrows.

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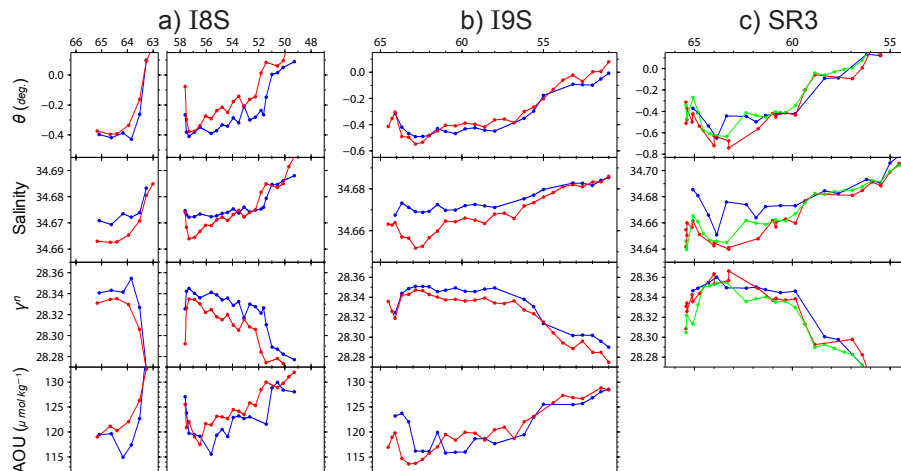


Fig. 2. Meridional variations of the water property averaged over the 100 m-thick layer at the bottom for **(a)** the I8S, **(b)** the I9S and **(c)** the SR3. Blue circles and lines are from the WHP occupations in the 1990s. Red circles and lines in **(c)** are from repeat occupations in the year 2001 (50–61° S), 2002 (61–66° S), 2003 (61–66° S) and green circles and lines in **(c)** are from those in 2008. All the red circles and lines in both **(a)** and **(b)** are from repeat occupations in the 2000s.

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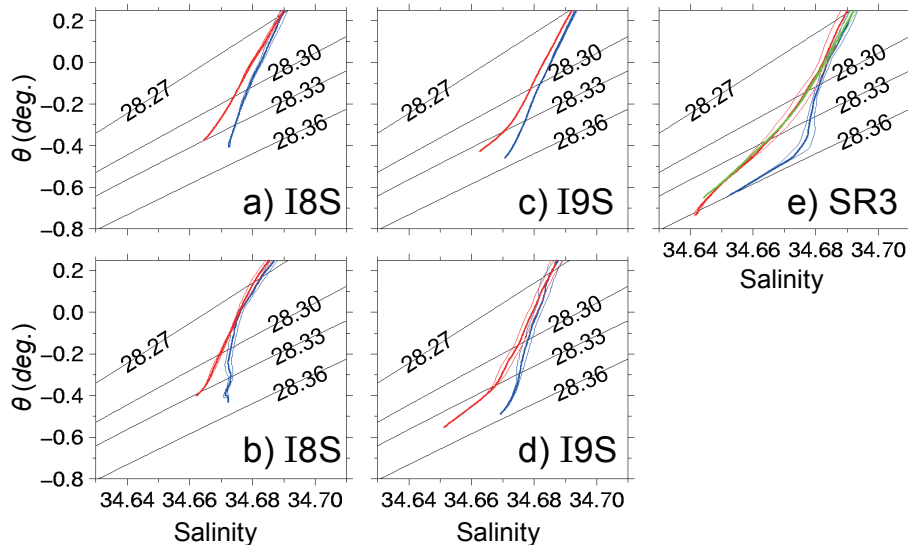


Fig. 3. Variations of $\theta - S$ curves for **(a)** the I8S (46–60° S, the Australian-Antarctic Basin), **(b)** the I8S (60–67° S, the PET), **(c)** the I9S (50–61° S), **(d)** the I9S (61–66° S) and **(e)** the SR3. Data are averaged on the isopycnal surfaces. Means are shown by thick lines and one standard deviation envelopes are shown by thin lines. Blue lines are from the WHP occupations in the 1990s. Red lines in e are from repeat occupations in the year 2001, 2002, 2003 and green lines in e are from in the year 2008. All the red lines in other panels are from repeat occupations in the 2000s. Labeled black lines are contour of neutral density.

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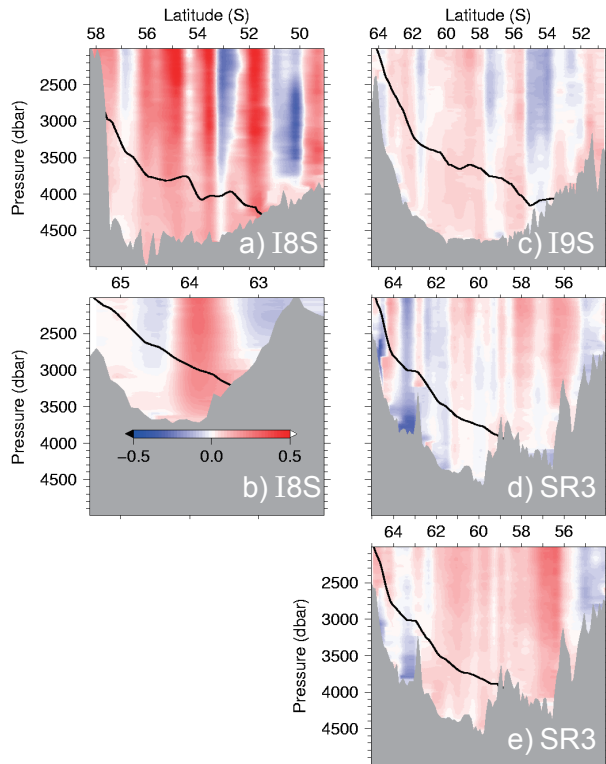


Fig. 4. Difference in isobaric potential temperature between the WHP and repeat occupations with the mean isopycnal of $\gamma^n = 28.30$. Red areas indicate warming and blue areas indicate cooling. **(a)** the I8S (the Australian-Antarctic Basin), **(b)** the I8S (the PET), **(c)** the I9S, **(d)** the SR3 (obtained from subtracting the WHP from combined repeat occupations in the year 2001, 2002, and 2003), **(e)** the SR3 (obtained from subtracting the WHP from repeat occupations in the year 2008). Mean isopycnals of $\gamma^n = 28.30$ are shown by solid lines.

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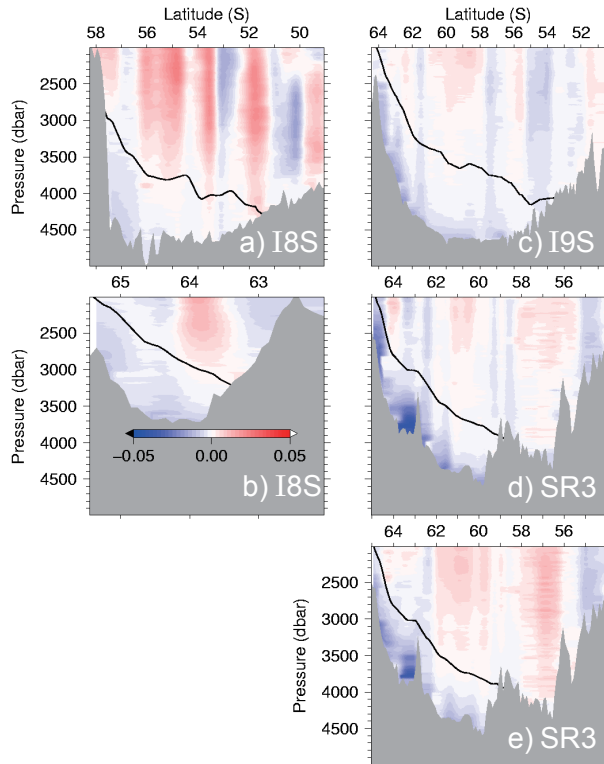


Fig. 5. Difference in isobaric salinity between the WHP and repeat occupations with the mean isopycnal of $\gamma^n = 28.30$. Red areas indicate increase in salinity and blue areas indicate freshening. **(a)** the I8S (the Australian-Antarctic Basin), **(b)** the I8S (the PET), **(c)** the I9S, **(d)** the SR3 (obtained from subtracting the WHP from combined repeat occupations in the year 2001, 2002, and 2003), **(e)** the SR3 (obtained from subtracting the WHP from repeat occupations in the year 2008). Mean isopycnals of $\gamma^n = 28.30$ are shown by solid lines.

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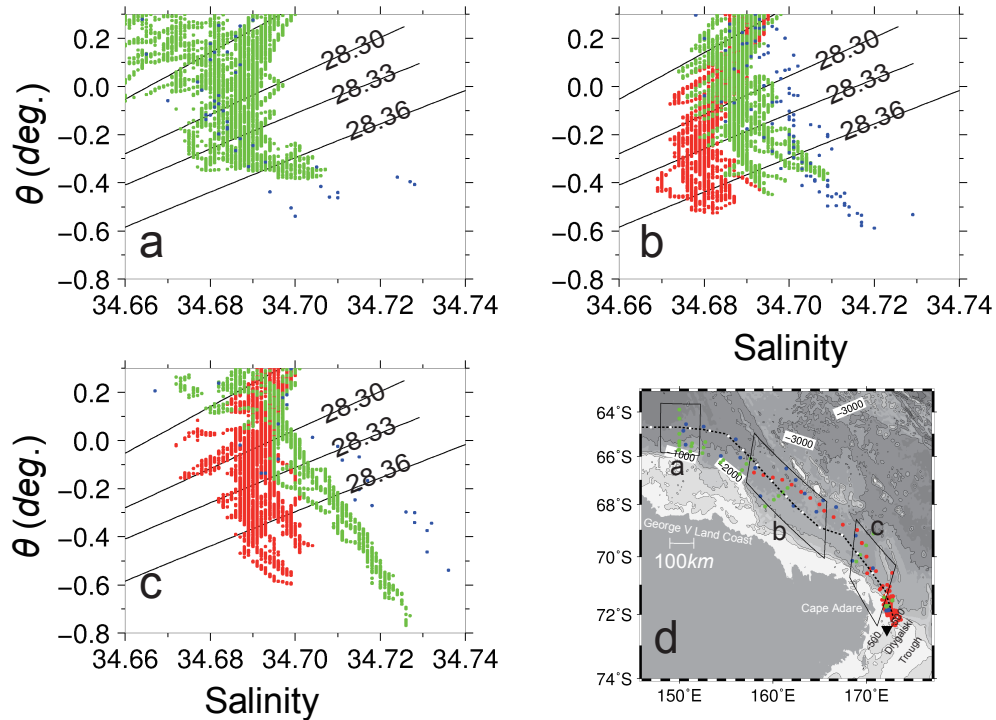


Fig. 6. θ – S diagrams of region between the Drygalski Trough and 150° E, (a) around 150° E, (b) around 160° E, (c) immediately downstream of the Ross source region, (d) locations of stations. Blue circles are from occupations in the 1970s, green circles are from 1990s and red circles are from 2000s. Labeled black lines in (a), (b), and (c) are contour of neutral density. Dashed line in d indicates approximate position of flow path of the RSBW and white circles are marked on 100 km interval.

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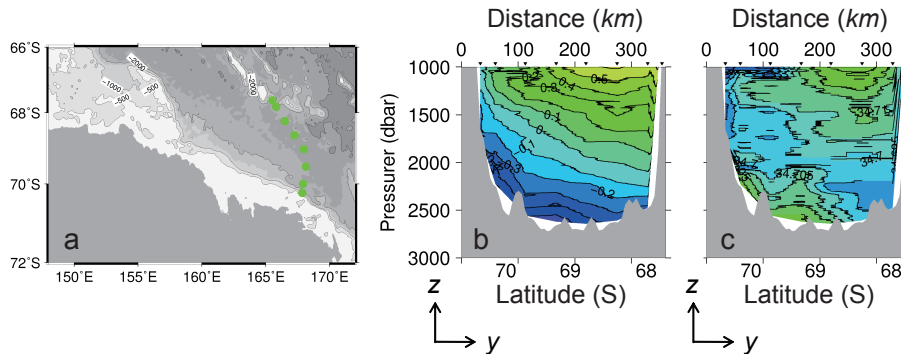


Fig. 7. Thickness and width (cross shore) of the RSBW from occupations in the 1990s. **(a)** Location of stations, **(b)** vertical cross section of potential temperature, **(c)** vertical cross section of salinity. y is cross shore coordinate (offshore positive), and z is vertical axis (upward positive). Thickness and width of cold and saline the RSBW are approximately 200 m and 200 km, respectively.

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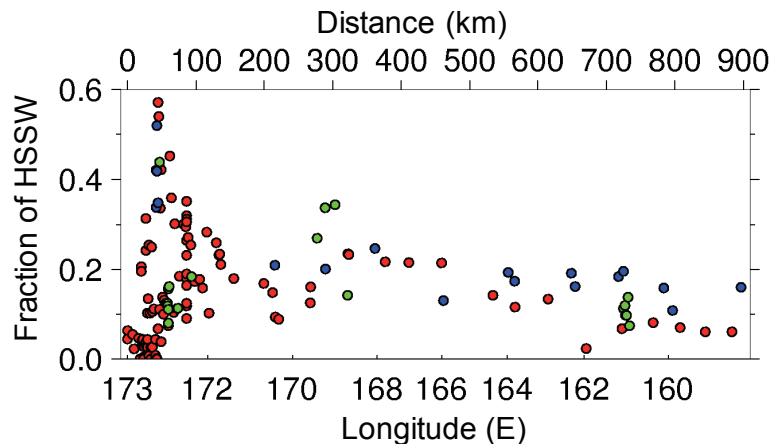


Fig. 8. Distribution of fraction of the HSSW within 100 m from bottom along flow path of the RSBW estimated by OMP analysis. Blue circles are from occupations in the 1970s, green circles are from 1990s and red circles are from 2000s.

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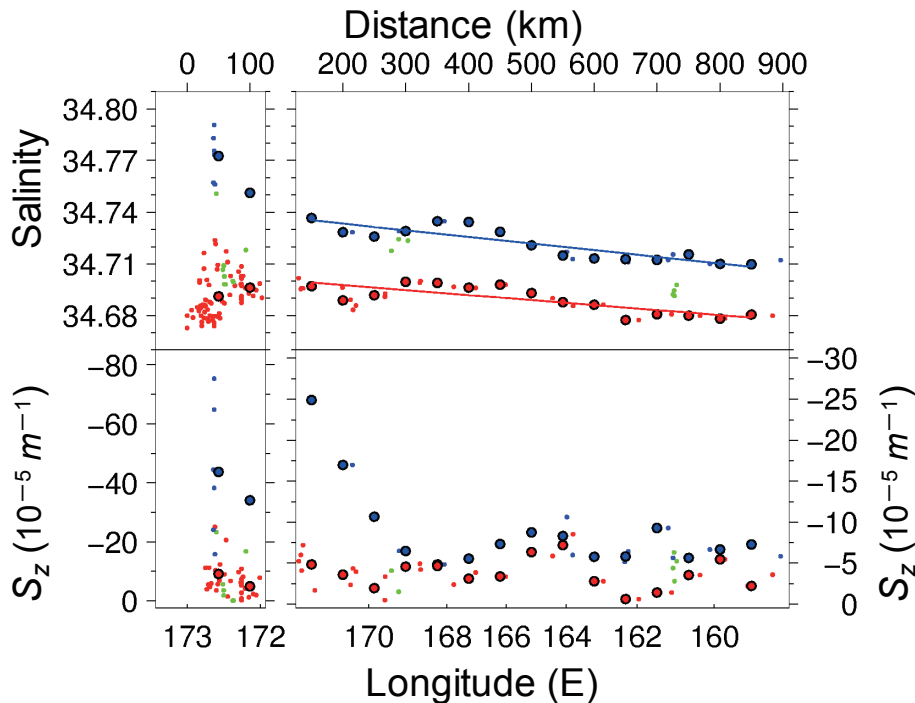


Fig. 9. Distribution of salinity averaged over the 100 m-thick layer at the bottom (upper panel) and vertical salinity gradient (lower panel) along flow path of the RSBW. Small circles are from occupations in the 1970s (blue), 1990s (green) and 2000s (red). In the estimation of vertical gradient, vertical gradients of 150 m scale are calculated at every 1 m interval by linear least square fit for each profile. The largest vertical gradients within 300 m from bottom are adopted as vertical gradient formed between the MSW/RSBW and the ambient water. Large circles are bin averages for the 1970s and 2000s and same colors as those for small circles are given. Solid lines in upper panel are regression lines for bin averages in region between 150 km and 900 km downstream from the Drygalski Trough.

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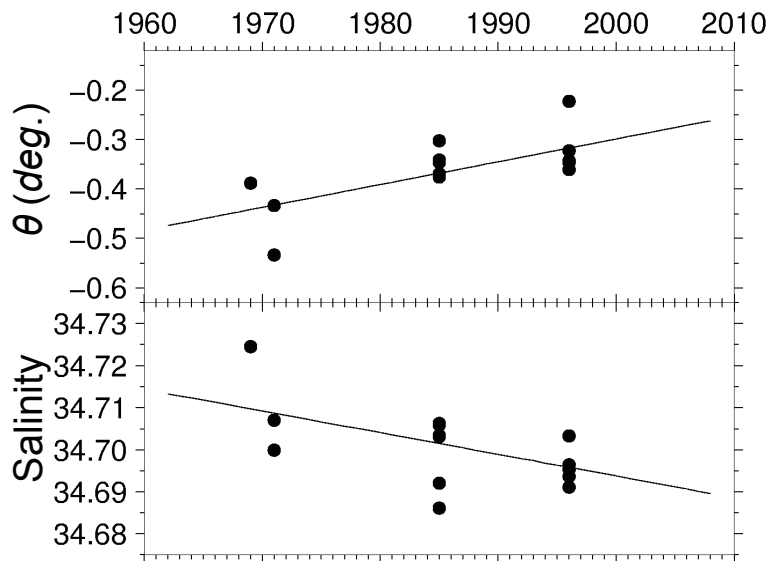


Fig. 10. Potential temperature (upper panel) and salinity (lower panel) averaged over the 100 m-thick layer at the bottom vs. year with regression lines from region around 150° E. Profiles located within the box a in Fig. 6d are used here.

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