

Effect of the North Equatorial Counter Current on the generation and propagation of Internal Solitary Waves off the Amazon shelf (SAR observations)

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We thank Ocean Science editors Eric Delhez and Natascha Töpfer, as well as Professors Leo Maas and Kevin Lamb for their constructive comments and suggestions, which have helped to improve the manuscript. A marked-up version of the manuscript is also provided, following the detailed responses to each of the referees concerns (changes in the main text highlighted in red).

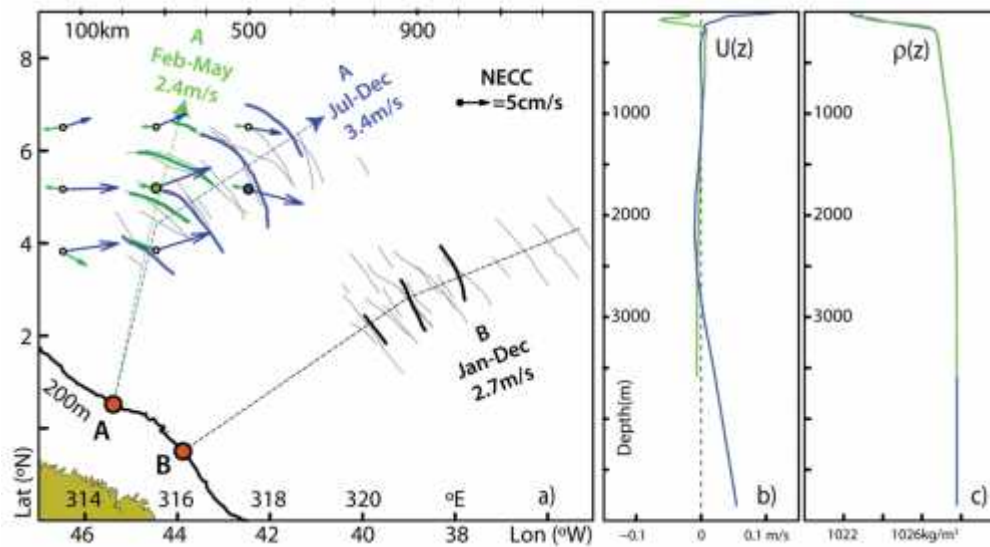
Referee #1

We thank **Referee #1** (L. Maas) for his helpful comments, and appreciate the thoughtful and constructive feedback on the paper, which has helped to improve the manuscript. Appropriate changes are now included in a marked-up version of the manuscript (in red), according to detailed responses to each of referee's concerns (below with referee's comments in *italic*).

Questions raised by referee #1:

1) What profiles for N and U are used to estimate eigenvalue c ?

R: The authors acknowledge that inputs used to solve BVP1 were not explicitly presented. These were computed from climatological monthly means in a similar fashion as described in Section 3.1 (i.e. provided from NOAA/OAR/ESRL PSD, Boulder, Colorado and available at <http://www.esrl.noaa.gov>). May and October means were computed in appropriate locations along the ISWs seasonal propagation paths – highlighted by green and blue filled circles (for May and October, respectively) in a new version of Fig. 9 (shown below, see panel a). The corresponding vertical profiles for stratification (ρ) and shear (U) are also displayed in new Fig. 9 in additional panels (b) and (c).



New Figure 9. (a) Same composite map as in Fig. 1 is shown in thin black lines, with two case studies highlighted in thick coloured lines, corresponding to location A and representative of two different seasons: from February to May (in green and dated 27 May 2009) and from July to December (in blue and dated 3 October 2011). Their averaged propagations speeds are also indicated along with idealized propagation paths (in thin dashed lines). Corresponding NECC monthly means (in green and blue for May and October respectively) are shown to depict its seasonal character along the ISWs propagation paths (see also scaled arrow in black). For reference, the case study shown in Fig. 3b is also highlighted in thick black lines, along with its mean propagation speed. (b) Climatological monthly means for horizontal velocity vertical profiles projected along the ISWs idealized propagation path, taken at green and blue filled circles in panel (a) (for May and October, respectively). (c) Same as panel (b) for potential density vertical profiles. See text for more details.

2) *To what extent is the geographical circumstance that the generation area lies in the freshwater (Amazon) outflow region used as a relevant aspect? Is the density gradient below the top layer (N) e.g. stronger than elsewhere, and does this explain the fact that these ISWs have the largest phase speeds ever observed? Or, is this due to an anticipated (yet not clearly substantiated) extra contribution of the NECC?*

R: A direct influence from the Amazon fresh water outflow in the ISW generation and propagation dynamics is not expected to be significant. Typically a NW Amazon plume is observed to develop and extend westward of 47° W (independently of the wind regime, since it is the NBC the leading forcing term), and hence not necessarily related with the ISWs shown in our Fig. 1, which are eastwards of 46° W (see e.g. Figs. 10 and 13 in Nikiema et al., 2007, and references therein). Another important density front, however, arises as the waves propagate into the NECC, whose influence is of much greater importance. Density gradients along the ISWs propagation paths are indeed strong and are reinforced as the pycnocline vertical extension decreases towards the open ocean (see our Figs. 1, 9 and 11). Nonetheless, density gradients (i.e. N) in this region top at around $2 \times 10^{-2} \text{ s}^{-1}$ (see e.g. our Fig. 11), which are comparable to other regions exhibiting fast propagating ISWs (where N at the pycnocline is also of order of 10^{-2}), such as the South China Sea or the Mascarene Ridges (see e.g. Li and Farmer, and da Silva et al., 2015, respectively). Therefore, the authors agree with the referee as he suggests an extra contribution from the NECC to be more likely to explain the elevated phase speeds reported in the manuscript. This is indeed anticipated in the present study using climatological data and SAR evidence, along with other in situ measurements reported in the literature, but as the referee points out a more substantiated approach is needed in order to challenge or confirm this hypothesis. This is briefly mentioned in the summary as a motivation for further studies including in situ or modelling results.

3) *Can one state something about the other (smaller scale) IWs of $O(1 \text{ km})$ wavelength, visible in zoomed versions of the satellite images?*

R: Please see point C below, concerning questions raised by the referee in the main text.

4) *Can one distinguish interfacial from obliquely propagating internal waves?*

R: Despite ISW radar signatures (propagating along the pycnocline) being well established in the literature (see e.g. Alpers, 1985), no successful attempts have yet emerged for the sea surface signatures of internal tidal beams and their impact on the ocean thermocline. An extensive study has also been undertaken in the Bay of Biscay to survey ISWs resulting from the “local generation” mechanism (see New and da Silva, 2002), which also included this same goal, but without any firm conclusions. We recall that satellite borne SARs are able to detect strain rates (i.e. horizontal current gradients) at the sea surface of the order of 10^{-3} s^{-1} (see e.g. Alpers, 1985), which are typical values for short-period internal waves in ocean (which may raise by an order of magnitude for strong nonlinear ISWs). However, while strain rates for IT beam impacts near the sea surface have never been measured in situ (to the best of our

knowledge), model results reveal these to be nearly an order of magnitude smaller and hence unlikely to be imaged in up-to-date SAR imagery. As an example, we present model results from the MITgcm computed for the Mascarene Ridge in the framework of recent investigations by da Silva et al. (2015). These modelling data reproduced quite well the observational SAR ground truth and are therefore considered to yield robust simulations of the surface velocity fields induced by the ISWs and IT beams. Fig. R1 (below) presents in its top panel a model frame for horizontal velocity (u , color-coded on the right) in which an IT beam may be seen emanating from the western slopes of an underwater bank (in yellow and reddish colours), and reflecting from the sea surface at around 740 km (between the vertical dashed lines). Farther east a well-developed ISW packet may also be seen propagating rightwards near 760 km. The corresponding surface strain rate associated with these phenomena is depicted in the bottom panel, in which the solitary disturbances may be readily identified (see black arrow) with surface gradients peaking at around 10^{-3} s^{-1} . However, strain rates (which essentially dictate the strength of the radar backscatter) in the vicinities of the beam impact are far less and of the order of the remaining background e.g. beyond 770 km, which are therefore unlikely to be detected in SAR images. Nonetheless, new generation satellite SARs with improved spatial and radiometric resolutions may yield more promising results, which should be explored in the future.

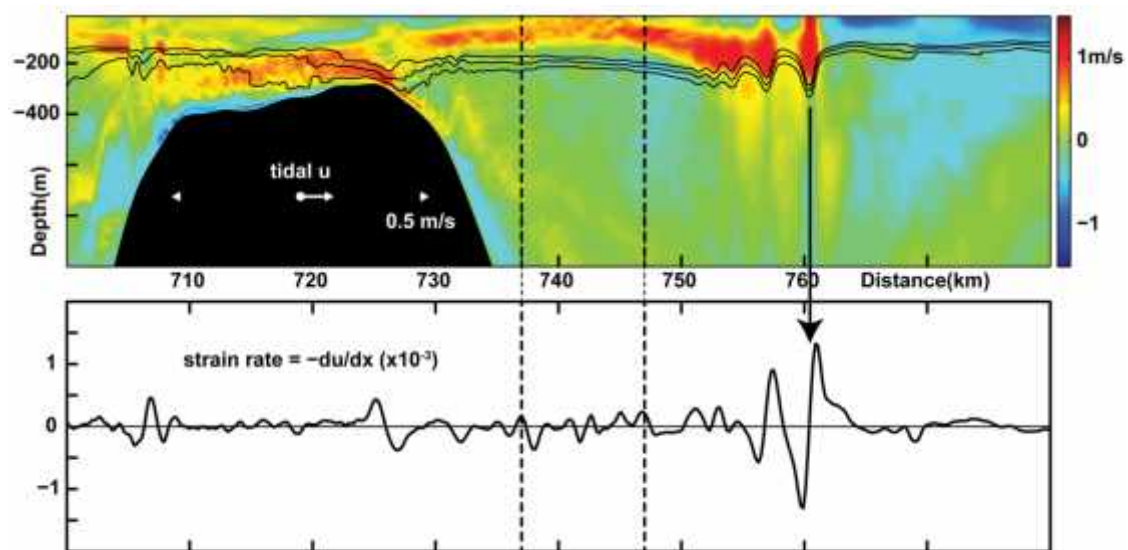


Figure R1. Top panel shows model horizontal velocities (u) along with selected isopycnals to highlight the presence of ISWs near 760 km. Bottom panel shows strain rates computed at the surface level (i.e. $z=5\text{m}$), which yield a proxy of the radar backscatter. Tidal phase is given in white labels in top panel. For reference the region showing the surface reflection of an internal tidal beam is bounded by vertical dashed lines in both panels.

5) Does ‘thermocline shrinking’ also explain the increase of the ISW’s phase speed and its larger penetration distance?

R: As predicted, for instance, by the fully nonlinear Dubreil-Jacotin-Long equation (suited for large amplitude ISWs, see e.g. Dunphy et al., 2011), ISWs will indeed become increasingly fast (hence traveling greater distances) with increasingly narrower waveguides – just as suggested by the referee. However, according to Fig. 9, the ISW sea surface manifestations depicted from SAR (taking generation site A for example) are seen north of 3 to 4° N, propagating offshore for several semidiurnal IT wavelengths, and reaching as far as 7° N. On the other hand, Fig. 11

shows a ‘*thermocline shrinking*’ essentially taking place between 1 and 3 to 4° N, which therefore occurs prior to the ISWs and includes essentially the propagation domain of the IT (i.e. from the shelf-break to 3-4° N) rather than the full propagation range of the ISWs. In fact, these extend farther NE along a pycnocline, whose vertical thickness is fairly constant beyond 4° N (cf. Fig. 11). ISWs from both seasons and emanating from site A are seen extending across the majority of the NECC (in Fig. 1), which is thus more likely to be responsible for the increased phase speeds of the ISWs and penetration distances between July and December (and vice-versa for Feb-May regime).

Questions raised by referee #1 in the main text:

(see also <http://www.ocean-sci-discuss.net/12/C1080/2015/osd-12-C1080-2015-supplement.pdf>)

A) pp. 2500, line 29. How do we know these are 'low modes'? Typically, the beam size is set by the shelf depth, which, when propagating oceanwards is small (thus of "high mode") compared to ocean depth. Also, the outflow of fresh water suggests generation of interfacial waves on the front (thermocline) separating these waters from the salty deep sea.

R: The particular set of ISW packets addressed in this study, whose inter-packet lengths can be estimated from the SAR imagery, are consistent with semidiurnal interfacial internal tides of the fundamental mode propagating in deep-water (similarly to the Bay of Biscay as discussed in New and da Silva, 2002). Typical horizontal wavelengths are of the order of one hundred kilometres and thus consistent with mode 1 vertical structures computed from BVP 1. Nonetheless, we acknowledge that higher modes may also be generated at the shelf-edge and propagate as internal tide beams into the open ocean. However, these cannot account for the present ISW observations, e.g. in the framework of the ‘local generation’ mechanism (see e.g. New and Pingree, 1992) as their ray tracing geometry predicts locally generated waves at approximately 100 km from the nearest shelf-break (assuming M_2 frequencies) – and thus in a location much closer to the shelf-break than the SAR observations (see e.g. our Fig. 3). At the same time, generation of large-scale interfacial tides resulting from fresh water outflows (from the Amazon River mouth) and associated density fronts is also unlikely, since these do not account for the waves’ semidiurnal period, nor is the river plume extension consistent with the ISW locations (see also point 1 above).

B) pp. 2502, line 15. You mean that the ISW signatures in Fig 1 are a composite of individual trains of ISWs observed? (You could add a number to them and a corresponding number to the lines of Table 1).

R: Fig. 1 is aimed at presenting a composite map resulting from the 17 image acquisitions listed in Table 1. We do not map all the individual waves obtained from all the images, but instead (for clarity) only the strongest ISW signatures were considered and only the leading wave is depicted for each ISW packet (in black solid lines, with a total of 59 packets being represented from the total 17 images). We note that all 17 image acquisitions were made in Wide-Swath mode, each covering 400x400km², and hence covering waves generated at multiple tidal cycles (as many as 5 like in Fig. 3a). Therefore adding a reference number to each of the ISW

observations depicted in Fig. 1 would result in an additional 59 labels, ranging from 1 to 17. An attempt to do so was undertaken but soon proved too cumbersome, including in fact several indistinguishable observations.

C) pp. 2503, line 5. Zooming in by 800% on Figs 3a,b, much shorter but often co-aligned IWs (?) are visible, clearly of O (1-2 km) wavelength (definitely shorter than the strong ISW (white bands) of 5-10km according to your estimate). Any explanation? Do they form part of the elusive 'dispersive tail' following (at slower speed) the solitons?

R: We appreciate the referee for highlighting these less evident observations of much shorter-scale ISWs, which were familiar to the authors. In fact, they are a part of the typical ISW field as seen from SAR imagery, frequently appearing alongside their larger mode-1 companions. Several explanations are still being investigated. One possibility includes shorter-scale mode-1 'wave-tails' resulting from mode-2 solitary-like waves traveling deeper down the water-column. Similar observations have been recently documented in independent studies for the South China Sea and in the Mascarene Ridge (see e.g. Guo et al., 2012; and da Silva et al., 2015, respectively). Fig. R2 illustrates this phenomenon in more detail as compared with a typical SAR observation in the Mascarene Ridge. Note that in between the larger mode-1 ISWs, there are other packets (different in nature) separated by approximately half the mode-1 semidiurnal IT wavelengths. A different hypothesis, as the referee points out, includes the sea surface manifestations of dispersive tails originating from the larger ISWs, as discussed for instance in Grimshaw et al. (2014). However, based on the current SAR collection and without further data sources (e.g. in situ or modelling) we can only speculate at this point, and hence prefer to discuss these separate observations as a different subject entirely within forthcoming investigations.

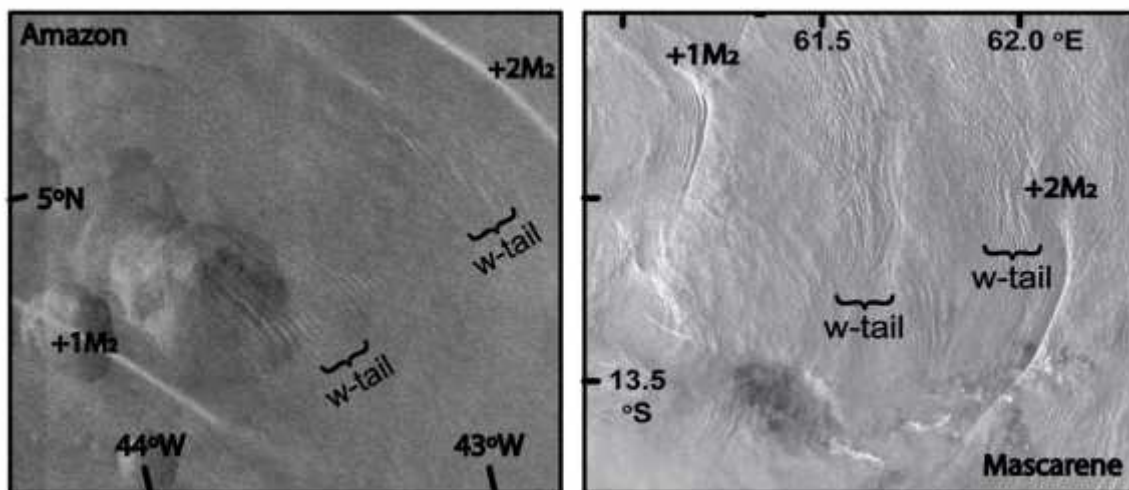


Figure R2. Left panel shows zoomed version of our Fig. 3a, highlighting the presence of shorter-scale 'wave-tails' in between the larger mode-1 ISWs. Right panel shows a subset of an ENVISAT-ASAR image of the Mascarene ridge of the Indian Ocean used in da Silva et al. (2015) and dated March 28, 2009. It shows a typical view of the Mascarene Ridge propagating ISW field. Sea surface signatures of two mode-1 ISW packets and 2 "wave-tails" are delimited within the image subset.

D) pp. 2506, line 11. But the additional 5 cm/s does not at all seem enough to explain the increase in propagation speed from 2.7 to 3.4 m/s... This invites for a deeper wave-current interaction consideration.

R: We agree with the referee since that the additional 5 to 10 cm/s resulting from the NECC climatological means between May and October seem small when comparing the ISW propagation speeds between the same periods (i.e. 2.4 m/s for May and 3.4 in October). However, our goal at this stage of the paper is simply to illustrate that there is in fact a seasonal reverse in the NECC, and that it is in qualitative agreement with both the climatological data and the SAR observations. The next two paragraphs, as the referee suggests are indeed aimed at quantifying how the NECC influences the waves' propagation speeds. This is done in a first approach by means of linear theory and BVP 1, which includes vertical shear, and hence the contribution of the climatological NECC. As discussed in the paper, a reasonable agreement is found for waves in May, but the linear phase speeds using NECC data for October (i.e. 2.5 m/s) yield far less than SAR estimated propagation speeds (around 3.4 m/s). However, we note that measurements performed in situ in November at this same location by Brandt et al. (2002, see their Fig. 1) revealed currents to be much greater than their climatological mean value – between 1 and 1.5 m/s within the ISWs propagation path (see also our Fig. 1). Running BVP 1 including a proxy of these results (instead of the climatological NECC) yields phase speeds between 3 and 3.5m/s), which are in better agreement with SAR estimated propagation speeds. Further studies of wave-current interactions are planned in a near future on the basis of additional high-resolution numerical modelling, as was done for the Mascarene Ridge (see da Silva et al., 2015). See also point (F) below summarizing additional evidence of measured NECC strength (as opposed to its seasonal average).

E) pp. 2506, lines 5 and 8. What profile is measured/assumed here? c seems to be much larger than u . Hence 'resonance' (vanishing of the denominator) is not very likely to occur. Not clear how you obtain this number? You need to specify $N(z)$ (from Fig. 11?) and $U(z)$ before an eigenvalue, c , can be computed.

R: We again agree with the referee since the vertical profiles for stratification and horizontal currents had not been specified in the manuscript. These followed from climatological data (as described in Section 3.1) for May and October in appropriate locations along the ISWs propagation paths, which are now highlighted in a new version of Fig. 9. The corresponding vertical profiles for stratification (ρ) and shear (U) are also displayed in additional panels (b) and (c). As the referee points out, a critical layer (i.e. a layer in which the speed of the background flow matches that of the long linear wave) is unlikely in this case, since linear wave speeds are roughly between 2 and 3 m/s, and the NECC contribution (i.e. $U(z)$) is at most of the order of 1 m/s.

F) pp. 2507, line 18. Any further info from literature on actual NECC strength (as opposed to its seasonal average)?

R: Complementary results to those reported in Brandt et al. (2002) obtained by drifter observations along the NECC confirm that the climatological data (provided from NOAA/OAR/ESRL PSD, Boulder, Colorado and available at <http://www.esrl.noaa.gov>) may be underestimated. According to Lumpkin and Garzoli (2005) (see also their Figs. 8 and 9), typical values for April and November not only confirm the NECC reversing character, but also reinforce that corresponding monthly means for near surface currents are of the order of 1 m/s, especially during November. This additional information was added at the end of Section 3.1.

G) pp. 2508, line 18. Without further explanation, idiosyncratic terminology.

R: Terminologies including solutions 6.1 and 18.5 (i.e. concerning the HYCOM simulations) were removed from the text, keeping only the corresponding references.

H) pp. 2510, line 13. Explain f and ω .

R: Definitions for ' f ' and ' ω ' are now included in Section 3.3.

I) Table 1. If the range (first line) is 2.25 m, (amplitude 1.12m) how can we distinguish the ascending or descending tide from a measured tidal height of -0.83m? Also: what does it mean that the tidal range varies so much (e.g. 1.34 m at next A, fourth line). This is a function of phase in the fortnightly modulation?

R: An additional column indicating time after high water (TAHW – computed correspondingly for sites A or B) for each image has been added to Table 1 to clarify tidal phases during satellite acquisitions (see also tidal heights and ranges as well as Fig. 2). The referee is correct since images are acquired in different phases of the fortnightly cycle, ranging from neap tides to spring tides (see also Fig. 2). Usually satellite acquisitions are aimed at spring tides, when conditions favour ISW formation (e.g. in the Bay of Biscay or the Mascarene Ridge). However, in the Amazon case ISWs appear to be imaged throughout the entire fortnightly cycle, even despite its strong modulation of the order of 1 m (see our Fig. 2).

J) Figure 11. Seems more representative of different locations.

R: Figure caption was clarified to make it clear that profile P1 is set just beyond the shelf-break where the pycnocline thickness is still wide and where ITs are being generated, while profile P2 is set just prior to the SAR ISWs observations, where thermocline thickness has already decreased substantially (roughly by 100 m) – and hence belonging to different locations as suggested by the referee.

Additional references (not included in the main text):

Oumarou **Nikiema**, Jean-Luc Devenon, Malika Baklouti, Numerical modeling of the Amazon River plume, *Continental Shelf Research*, Volume 27, Issue 7, 1 April 2007, Pages 873-899. Doi: 10.1016/j.csr.2006.12.004.

Qiang **Li** and David M. **Farmer**, 2011: The Generation and Evolution of Nonlinear Internal Waves in the Deep Basin of the South China Sea. *J. Phys. Oceanogr.*, 41, 1345–1363. Doi: 10.1175/2011JPO4587.1.

New, A. L., and R. D. **Pingree** (1992), Local generation of internal soliton packets in the central Bay of Biscay, *Deep Sea Res., Part A*, 39, 1521– 1534. Doi: 10.1016/0198-0149(92)90045-U.

Roger **Grimshaw**, Chuncheng Guo, Karl Helfrich, and Vasiliy Vlasenko, 2014: Combined Effect of Rotation and Topography on Shoaling Oceanic Internal Solitary Waves. *J. Phys. Oceanogr.*, 44, 1116–1132. Doi: 10.1175/JPO-D-13-0194.1.

Rick **Lumpkin**, Silvia L. **Garzoli**, 2005. Near-surface circulation in the Tropical Atlantic Ocean Deep Sea Research Part I: Oceanographic Research Papers, Volume 52, Issue 3, Pages 495-518. Doi: 10.1016/j.dsr.2004.09.001.

Referee #2

We thank **Referee #2** (K. Lamb) for reviewing the manuscript, and for his thorough recommendations to improve the paper. Appropriate changes are now included in a marked-up version of the manuscript (in red), according to detailed responses to each of the referee's concerns (below with referee's comments in italic).

Points raised by referee #2:

1. *There are a number of places where the English could be improved.*

(a) *abstract, lines 10–11: "... ever recorded. The main characteristics of the ISWs are ...".*

(b) *page 3, line 8: "have benefited from"*

(c) *page 3, line 16: "... concerning global ..., in particular that owing to ..."*

(d) *page 3, line 21: "...significantly far apart ..."*

(e) *page 3, line 23: how does satellite altimetry confirm that IWs transport mass?*

(f) *page 4, line 1: "... by other remote sensing methods to propagate considerable ..."*

(g) *page 4, line 7: "feature the Amazon shelf break as an important hospot for the ..."*

(h) *page 4, line 16: what is meant by 'admittedly' here? Do you mean 'presumably'. Later it should read "transfers to smaller-scale"*

(i) *page 4, lines 17–18: "'Brandt et al. (2002) confirmed the existence ..."*

(j) *page 4, line 22: "inter-packet distances ...".*

(k) *page 5, line 23: "... implications for the redistribution"*

(l) *page 6, line 6: "... Section 2 describing ..."*

(m) *page 7, lines 3–4: '... preferences, since favorable stratification conditions for ISWs occur year-round in equatorial regions ...'*

(n) *page 11, line 15: "associated with the increase in their average ..."*

(o) *page 11, line 19: "... related to the ..."*

(p) *page 12, lines 13–16: "Note also that (delete comma) changes in ... because B generated waves only partially intersect the NECC (i.e. along"*

(q) *page 13, line 2: "- represents the ..."*

(r) *page 13, line 18: "underestimating its true influence...."*

(s) *page 13, line 21: I suggest deleting the first clause to start with "We now turn to ..."*

(t) page 14, line 16: “... useful for guiding ...”

(u) page 16, line 2: “... a significant reduction in C occurs between A and B”

(v) page 17, line 22: “vertical extent”

(w) page 18, line 11: “Two vertical profiles ..”

(x) page 20, line 5: what is mean by Ht performing?

R: We appreciate all the referee’s notes, and changes were made as instructed. In particular, for points (e), (h) and (x) please note that:

(e) page 3, line 23: how does satellite altimetry confirm that IWs transport mass?

R: We acknowledge that this sentence may be confusing, since satellite altimetry is confirming only the long propagations distances for the internal tides. The sentence now reads: “Satellite altimetry studies (see e.g. Ray and Cartwright, 2001) have indeed confirmed propagating distances of the order of 1000 km for the long Internal Tides (ITs, i.e. IWs of tidal frequency), and shorter-scale Internal Solitary Waves (ISWs) have been shown by other remote sensing methods to propagate considerable distances as well (e.g. da Silva et al., 2011; Guo et al., 2012)”.

(h) page 4, line 16: what is meant be ‘admittedly’ here? Do you mean ‘presumably’.

R: We agree that ‘presumably’ conveys the appropriate meaning (and not ‘admittedly’), since these energy transfers to smaller-scale waves were only suggested based on visual observations made by the crew. The change was made accordingly.

(x) page 20, line 5: what is mean by Ht performing?

R: The term was replaced by ‘decreasing’ so that the sentence reads: “... with Ht decreasing accordingly (see blue line in Fig. 12a).”

2. I think it would be helpful to give more information on water depths where the waves are observed.

R: We agree with the referee and depth contours for 200, 500, 1000, 2000, 3000, 4000, and 4500 m have been added in thin grey lines to Fig. 9a. Note the 200m and 1000m contours are hardly distinguishable from the 500 m isobath (highlighted in a black thick line), and thus representative of a steep shelf-break, where the waves are believed to be generated.

3. Page 13, equation (3). Should be $w(-H) = 0$.

R: Eq. (3) was corrected as instructed in the manuscript.

4. ISW widths are given in Figure 6. How exactly was the wave width defined?

R: The characteristic soliton widths presented in Fig. 6 are defined from their corresponding SAR sea surface signatures, as shown in Fig. 4a (labelled ' L ') for a normalized backscatter intensity profile taken across a representative solitary wave shown in Fig. 3b. This is a typical parameter estimated in SAR studies involving ISWs (see e.g. New and da Silva, 2002; da Silva et al., 2011), which is meant as a measure of the individual waves horizontal scales (along their propagation direction). It is important to note that ISWs propagating on the near surface thermocline (or pycnocline) are capable of producing sea surface roughness patterns that are observable by satellite-borne SARs when the wind speed is not excessively strong (e.g. Alpers, 1985). A schematic view is shown in Fig. R1 (bellow), according to hydrodynamic modulation theory (as described in Alpers, 1985), where it can be seen how the isopycnal displacements below (and corresponding dimensions) translate to the SAR images above. According to this figure, mode-1 ISW interfacial waves travelling along the pycnocline (assuming waves of depression, which is often the case in deep waters) will be revealed in SAR images by bright bands preceding darker ones in their direction of travel. In essence, this follows from surface velocity fields induced by travelling ISWs, which create different convergence and divergence patterns, modulating the surface roughness and thus the intensity of the radar backscatter signal. In practice, however, ISW sea surface signatures are modulated by other ambient factors, such as the local wind speed and direction, the presence of organic surface films, and possible resonant interaction with surface waves (see Alpers, 1985; da Silva et al., 1998; Brandt et al., 1999; Araújo et al., 2002; and Kudryavtsev et al., 2005 for a detailed discussion). In these cases the SAR intensity variations, which are usually symmetrically distributed about a mean value close to zero, can be shifted towards more positive or negative values depending if the waves are seen as bright bands in a darker background or vice-versa. Therefore, all our estimates for L yield only a proxy for the waves' horizontal dimensions. Further notes, conveying a summary of the above information were inserted in Section 2 (along with appropriate references), as additional information to clarify this issue for the broader oceanographic community. We stress that the length scale L is defined as the distance where the normalised backscatter is greater than I_0 (see Section 2) along the ISW propagation direction in the cases of positive-only image contrasts, and the distance corresponding to positive followed by negative backscatter variations (relative to I_0) in the cases of positive/negative image contrasts (see Figs. 4a and 4b, respectively). Hence, the reference is I_0 and no noise level is considered since our procedure has been done manually for each ISW.

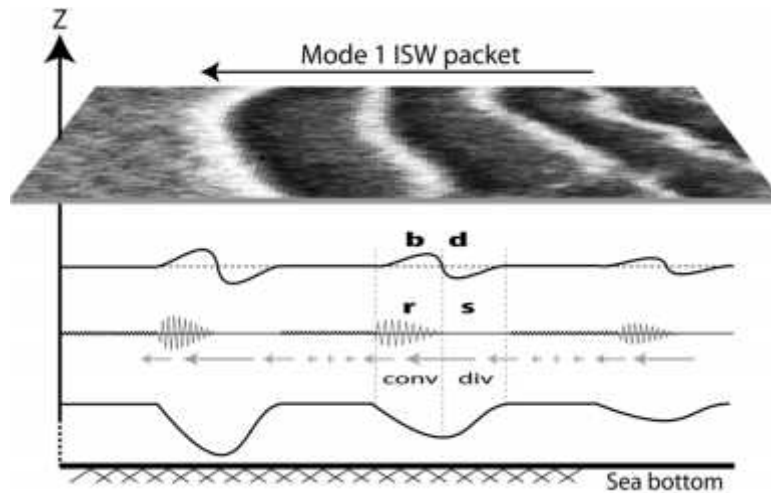


Figure R1. Example of a SAR image showing sea surface signatures of a mode-1 ISW packet. The ISWs are assumed to be moving from right to left. From top to bottom the horizontal profiles represent the following features: SAR intensity profile along the ISWs, with bright enhanced backscatter (b) preceding dark reduced backscatter (d) in the direction of propagation; surface roughness representation indicating how rough (r) and smooth (s) the surface is along an ISW wave packet; surface current variability induced by ISWs (note indication of convergence and divergence fields near the surface); isopycnal displacements produced by ISW propagation.

5. Referring to HYCOM simulations 6.1 and 18.5 is not very informative. I don't see any value if even mentioning 'simulation 6.1'. Can you give more information regarding resolution, how long the simulation was for, does it include the NECC and the North Brazilian Current etc.?

R: Terminologies including solutions 6.1 and 18.5 (i.e. concerning the HYCOM simulations) were removed from the text, keeping only the corresponding references. We agree with the referee in adding further details concerning the HYCOM simulations used in paper. These was included in Section 3.2, summarizing that HYCOM is a realistically forced three-dimensional global ocean model (3D global HYbrid Coordinate Ocean Mode, see Bleck 2002), with both tidal and atmospheric forcing, and hence including the ocean's major current systems, such as the NBC or the NECC. The simulations used in this manuscript refer to an annual period (from October 2011 to September 2012) and their tidal forcing includes the largest semi-diurnal and diurnal constituents (i.e. M_2 , S_2 , N_2 , K_2 , and K_1 , O_1 , P_1 , Q_1), where a nominal horizontal resolution is set to $1/12^\circ$ at the equator with 32 layers in the vertical direction. Hourly 3D time series were monthly averaged, whereas Fig. 10 displays the depth integrated and time-mean (over one full year) conversion rates (C , left panel) and energy fluxes (F_E , right panel) for the filtered semi-diurnal tides. A more detailed description of the HYCOM simulations used in this paper may be found in Buijsman et al. (2015) and references therein (see also Metzger et al., 2010 and Shriver et al., 2012).

6. Page 16, first paragraph. Is there a reason for the reduced generation between generation sites A and B? Reduced tidal currents?

R: A tidal ellipse map computed from a regional solution (using the Oregon Tidal Prediction Software, OTIS, at a resolution of $1/12^\circ$, see Egbert and Erofeeva, 2002) is presented in Fig. R2. No significant changes in the along-shelf tidal ellipses can be seen, in particular in between A

and B. We note these ellipses are representative of neap-tides, whereas a similar view still holds for spring tides. However, the weaker conversion rates given at the promontory near 315.5°E and 0°N (see our Fig. 10a) are in fact consistent with our SAR observations, which show no significant ISW activity in between A and B. Therefore, at this point, we prefer to discuss the separation between distinct generation sites A and B on the basis of the conversion rates yielded by the HYCOM simulations, since these systematically (i.e. throughout the year) present similar results to those present in Fig. 10a. Nonetheless, we appreciate raising this particular aspect of the study region and possible reasons, such as the convex (as seen from the open ocean) geometry of the promontory (scattering energy rather than focusing it), are set to be further investigated within forthcoming 2D high-resolution simulations in several across-shelf transects between A and B, as well as in the frame of 3D configurations (using the fully nonlinear and non-hydrostatic MITgcm).

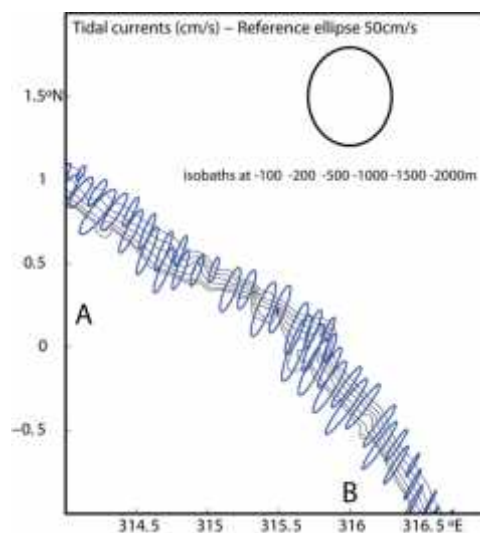


Figure R2. Ellipse tidal map computed from a least-square fit to the 1/12° regional tidal model OTIS run for one tidal cycle on 3 October 2011 (i.e. close to neap-tides). A 50 cm/s reference ellipse is also shown.

7. page 19, line 8: Shouldn't use 'precisely' here. This is not very precise. Why was a 30 m amplitude internal tide used? What if it was 25 m or 20 m?

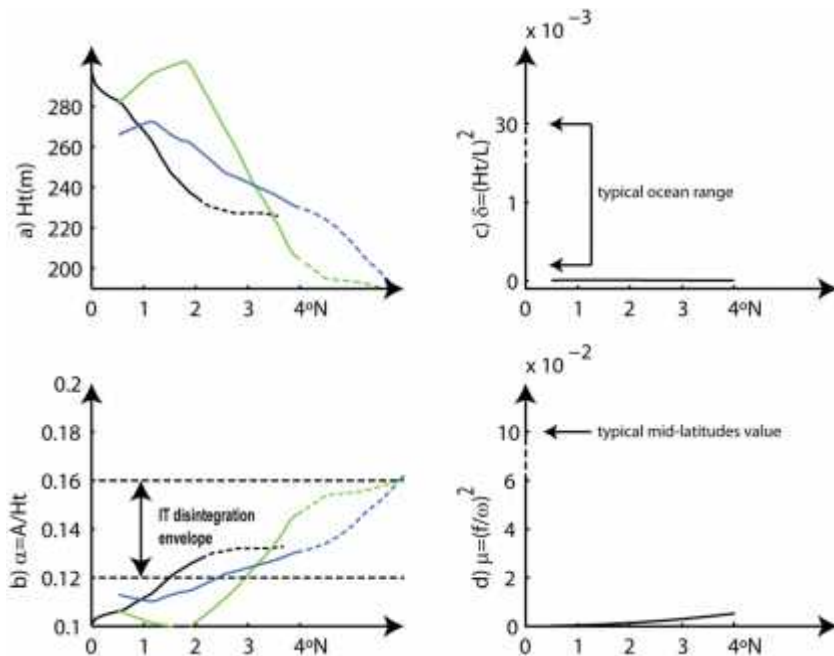
R: We agree with the referee and have replaced 'precisely' with 'approximately' in the new version of the manuscript. The nominal value for the IT amplitude was inferred from Vlasenko et al. (2005) and their Fig. 3.3, which shows low frequency internal oscillations measured in the study region with amplitude of approximate 30 m. We acknowledge that this value may change between generation sites and amongst different seasons, hence changing the overall agreement between the ISW observations and the disintegration envelope in Fig. 12b. However, we note that the key parameter governing nonlinearity in Fig. 12b is the thermocline thickness (i.e. H_t in Fig. 12a), and in this case the agreement between SAR and climatological data would still remain – i.e. ISWs are only observed after ITs propagate along an increasingly narrow pycnocline where nonlinearity effects increase significantly.

8. For several of the figures the labelling of the horizontal axis could be improved. For example in Figure 5 'km' appears directly under the fifth column. It would be better to have the number 300 there with 'km' below the numbers along the axis. Several other figures have this problem.

R: Axis on Figs. 4 to 7 where edited as suggested by the referee. However, we note that on some occasions (e.g. SAR images) axis are labelled in order to optimize the area available for figure contents (usually small in printed versions), which should appear as large as possible (see e.g. Fig. 3).

9. The caption for Figure 12 should say what the different coloured curves are for and what the solid vs dashed lines are for. It would seem better to change the vertical scales for panels (c) and (d) to zoom in, although really I don't see that these figures are necessary. The values for the two dispersion parameters are discussed in the text and I don't see that these figures add anything. It would be helpful to indicate 1° and 3° N along the axis and perhaps distances.

R: A new version of Fig. 12 follows below, along with a new figure caption, according to the referee's suggestions. Additional labels were added to panels (c) and (d) in order to illustrate characteristic ocean values for dispersion as compared to typical values in the study region, and hence we prefer to keep these for consistency. Also, distances could not be added to Fig. 12 since three different wave paths are depicted in each panel (corresponding to different colours, see also Fig. 9), which would necessarily translate into three different distance axes (whereas latitude may be used as a common reference).



New Fig. 12. (a) Vertical extension of the waveguides calculated along the waves' propagation paths (in green, blue and black, according to dashed propagation paths in Fig. 9). (b) Same as panel (a) for the nonlinear parameter computed assuming an IT amplitude of 30 m. (c) Nonhydrostatic dispersion for $L = 100$ km. (d) Rotational dispersion for the semi-diurnal IT. Transitions to dashed lines in left panels indicate first SAR evidences of ISWs. Note broken vertical axes in right panels and typical ocean values showing comparatively weak dispersions in our study region. See text for more details.

10. It would help if distances were included in Figure 11 instead of just latitude. The difference in stratification between profiles P1 and P2 doesn't seem to be very large. Perhaps show some N profiles at some higher latitudes. What is the stratification like closer to shore or does figure go to the 500 depth contour?

R: Distances along the horizontal axis are now included in a new version of Fig. 11. Profiles P1 and P2 have been selected as representative of the waveguide at the shelf-break (depicted between the 200 and 1000m in the new version of Fig. 9) and just prior to ISW evidence seen from SAR (respectively). While the emphasis is on the waveguide narrowing, which is best seen in the density contours (in left panel), these profiles add additional information showing that the pycnocline (in yellow colours) is brought closer to the surface and slightly reinforced. We agree, however, that these differences are indeed not very large as the referee points out, which may be a consequence of their climatological (and hence averaged) nature, just as for the horizontal currents discussed in section 3.1. We also note that profiles to the north of P2 (thus farther offshore of the ISWs earliest evidences) would actually reveal an even stronger pycnocline, while the shallowing effect would be lost (cf. yellowish colours in left panel go deeper as they move further offshore). Nonetheless, the waveguide vertical decreasing is well depicted in Fig. 11a and in good agreement with SAR observations, whereas larger variations in stratification profiles may be appreciated from the in situ measurements shown in Fig. 3.2c in Vlasenko's et al. (2005). While slightly westwards of the ISWs propagation paths (considering site A, see our Fig. 1), their measurements reveal changes in the pycnocline depth of the order of 100 m as it shallows in the offshore direction. Finally, depth contours for 200, 500, and 1000 m are practically coincident and representative of a very narrow shelf-break (see new Fig. 9), where the IT waves are believed to be generated (see also our Fig. 10). The 500 m contour was therefore taken as representative of the shelf-break and used as an onshore limit for Fig. 11a – past which the climatological stratification used in this case would illustrate on-shelf conditions, which do not concern the offshore propagating ITs and associated ISWs observed in deep ocean.

11. Page 12, line 12. The current often looks to be more in the direction normal to that of wave propagation, not in the opposite direction.

R: We again agree with the referee, whereas currents in May flow approximately in the opposite direction of those seen in October, and not directly against the ISWs mean propagation path. This was changed in the text accordingly.

12. Does the Amazon river outflow make any contribution to the observed phenomena? For example surface currents, setting the stratification? What is the cause of the offshore thinning of the stratified layer? Is it simple geostrophic balance with the Brazilian current or does the Amazon river plume have anything to do with it?

R: We appreciated the referee's suggestion, but in this particular case the Amazon River outflow is unlikely to influence the ISWs dynamics, setting neither their near-surface currents

nor their stratification. Typically a NW Amazon plume is observed to develop and extend westward of 47° W (regardless of the wind regime, since it is the NBC the primary forcing term), and hence not necessarily related with the ISWs shown in our Fig. 1, which are eastwards of 46° W (see e.g. Figs. 10 and 13 in Nikiema et al., 2007, and references therein). As the referee suggests the offshore thinning of the waveguide is related to the NBC and the NECC, which are assumed to be in geostrophic and hydrostatic balance along the waves propagations paths (i.e. as they move increasingly far from the equator), and therefore in agreement with the thermal wind (shear) equation, which governs large-scale circulations in the ocean. When changing temperature for density, a relation between horizontal density gradients and vertical sheared flow may then be written as:

$$\frac{\partial \mathbf{u}_g}{\partial z} = -\frac{g}{f\rho_0} \hat{\mathbf{z}} \times \nabla \rho \quad (\rho),$$

where \mathbf{u}_g is the geostrophic velocity vector (in the zonal and meridional directions, and z is the vertical coordinate), g is the acceleration due to gravity, f the Coriolis parameter, ρ is the background density and ρ_0 some reference value. According to this relation, density changes in the horizontal are intertwined with geostrophic current variations in the vertical (i.e. shear). In particular, Fig. R3 shows how this translates to our study region where the NBC flows westwards, and the NECC flows either eastwards in October, or also westwards in May (but with much weaker currents). The density section (in light red) begins with a thicker waveguide at lower latitudes, owing to the negative vertical shear of the westward NBC (see also Johns et al., 1998) which holds at about 3°N (see also a surface map for the North Brazilian Current here http://oceancurrents.rsmas.miami.edu/atlantic/north-brazil_3.html). Between the NBC and the NECC a thinner waveguide is seen to develop and extend northwards, resulting from either a weaker westward NECC in May, or even a strong eastward NECC in October (in green and blue arrows in Fig. R3).

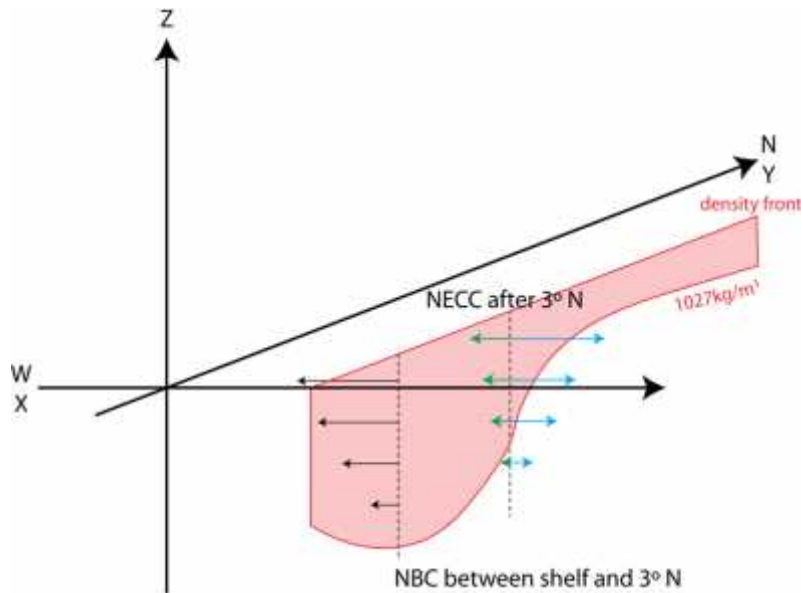


Figure R3. In the y - z plane a density front is seen towards north, representative of that seen in our Fig. 11, where isopycnals (see e.g. the 1027 kg/m³ red line) are first seen slightly falling and then raising substantially, meaning that at some depth z , the density is first decreasing and then increasing with increasing latitude. According to the thermal wind equation, the zonal velocity must therefore first decrease in the vertical as seen in the NBC (in black arrows), and then decrease to a less extent or even increase with z as seen in the NECC (for May in green and October in blue, respectively).

Additional references (not included in the main text):

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Effect of the North Equatorial Counter Current on the generation and propagation of Internal Solitary Waves off the Amazon shelf (SAR observations)

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Abstract:

Synthetic Aperture Radar (SAR) imagery from the Amazon shelf-break region in the tropical West Atlantic reveals for the first time the two-dimensional horizontal structure of an intense Internal Solitary Wave (ISW) field, whose first surface manifestations are detected several hundred kilometers away from the nearest forcing bathymetry. Composite maps and an energy budget analysis (provided from the Hybrid Coordinate Ocean Model – HYCOM) help to identify two major ISW pathways emanating from the steep slopes of a small promontory (or headland) near 44 °W and 0 °N, which are seen to extend for over 500 km into the open ocean. Further analysis in the SAR reveals propagation speeds above 3 m/s, which are amongst the fastest ever recorded. The main characteristics of the ISWs ~~main characteristics~~ are further discussed based on a statistical analysis, and seasonal variability is found for one of the ISW sources. This seasonal variability is discussed in light of the North Equatorial Counter Current. The remote appearance of the ISW sea surface manifestations is explained by a late disintegration of the Internal Tide (IT), which is further investigated based on the SAR data and climatological monthly means (for stratification and currents). Acknowledging the possibility of a late disintegration of the IT may help explain the remote sensing views of other ISWs in the world's oceans.

1. Introduction

Research efforts concerning Internal Waves (IW) are often motivated by satellite observations, which have the unique ability to render their two-dimensional horizontal structure (see e.g. Osborne and Burch, 1980; Alpers and Salusti, 1983; Apel et al., 1985; New and da Silva 2002; Sherwin et al., 2002; Ramp et al., 2004; Grisouard et al., 2011; Guo et al., 2012; Jackson et al., 2012; Mercier et al., 2012; Kozlov et al., 2014). Both new ~~ly~~ and unidentified IW hotspots, as well as previously studied regions, have ~~been~~ benefited from satellite views, frequently providing new and deeper insights into their generation, propagation and dissipation mechanisms (e.g. Zhao et al., 2004; Vlasenko and Alpers, 2005; Azevedo et al., 2006; Magalhaes et al., 2012) – see also <http://jmagalhae0.wix.com/internal-waves->. Data synergy, including in situ measurements and numerical modelling have also contributed to bridge IWs across multidisciplinary frameworks, spanning from fundamental oceanography to important applications (see e.g. review papers by Garret and Kunze, 2007; Lamb, 2014 and Alford et al., 2015). For instance, open questions still remain concerning ~~the~~ global tide energy dissipation, in particular that owing to IWs, which are important in ocean mixing and climate studies (see e.g. Zhao et al., 2012; Alford et al., 2015).

A regained interest has also come from acknowledging IWs as one of the available mechanisms by which mass and momentum are transported in the oceans, while recognizing that sources and sinks may be significantly far apart (see e.g. Moum et al., 2007; Ferrari and Wunsch, 2009; Shroyer et al., 2010; Zhang et al., 2015).

Satellite altimetry studies (see e.g. Ray and Cartwright, 2001) have indeed confirmed propagating distances of the order of 1000 km ~~Satellite altimetry studies (see e.g. Ray and Cartwright, 2001) have confirmed it~~ for the long Internal Tides

(ITs, i.e. IWs of tidal frequency), and shorter-scale Internal Solitary Waves (ISWs) have been shown by other ~~means of~~ remote sensing methods to propagate ~~along~~ considerable distances as well (e.g. da Silva et al., 2011; Guo et al., 2012).

The Amazon shelf-break is an important source for intense ITs (see Fig. 01 for location) as was early recognized in the work of Baines (1982) – exceeded only by the classical Bay of Biscay and the more recently studied South China Sea (see his Figure 10). Similar results can be equally found in more recent models, which also feature ~~the Amazon shelf-break as an important hotspot for this study region as an important hotspot in~~ the conversion of barotropic to baroclinic energy (see e.g. Buijsman et al., 2015). The presence of ITs and ISWs has already been documented with in situ measurements (Ivanov et al., 1990; Brandt et al., 2002 and Vlasenko et al., 2005), and acknowledged in remote sensing data (see also Jackson, 2004), but no detailed description has yet emerged regarding the two-dimensional horizontal structure of the ISW field.

Previous studies include the work by Ivanov et al. (1990) documenting IT energetics across the North Equatorial Counter Current (NECC). According to these authors, a decrease of energy density in the semi-diurnal IT was found close to the core of the NECC, ~~admittedly-presumably~~ owing to energy transfers ~~into~~ smaller-scale IWs (reported visually by the onboard crew). More than a decade later, Brandt et al. (2002) ~~did~~-confirmed the existence of large-amplitude ISWs in this region, using high-resolution acoustical data (collected during November 2000). In their study, ship-mounted ADCP measurements along a meridional section (coincident with the NECC, see Fig. 01 for location) were used to measure the waves' main properties. Their findings included dominant directions close to 30 °T and inter-packets distances of the order of 70 km. More recently, Vlasenko et al. (2005) also reported

on data collected close to generation site A in Fig. 01 (during November 1980). According to these last authors, a density front was found in this region of the tropical West Atlantic (see Fig. 1 and their Figures 3.1 and 3.2), which is associated with the NECC. A very similar setting is also discussed in the South China Sea by Buijsman et al. (2010a), where the Kuroshio Current is also associated with a density front, with important implications for ISW development there. This feature is of particular interest to this study and its implications are further discussed in Section 3.

The specific location of the Amazon shelf-break region adds further motivations to the oceanographic framework. It is placed in the tropical West Atlantic together with an intricate current system, including two major ocean currents – the NECC and the North Brazilian Current (NBC, see Fig. 1 for locations). These are found very close to the Amazon River mouth, which accounts for up to one ~~third~~-fifth of the total fresh water input into the world's oceans (see e.g. Wisser et al., 2010). The NECC is highly variable, featuring a strong annual cycle (see e.g. Garzoli and Katz, 1983) and significant inter-annual variability in its mean location and strength (Hormann et al., 2012). In fact, the NECC includes a seasonal reverse cycle in the east-westward direction (see e.g. Garzoli and Katz, 1983 for further details), which will be shown to have a major influence in the propagation of ISWs.

The ISWs presented in this study will be shown to propagate ~~for~~-over considerable distances across major oceanic current systems. The Amazon shelf-break is particularly interesting because low-mode ITs generated at the shelf-break and propagating into the open ocean may interact with eddies and the meandering field of steady currents such as the NBC and the NECC. Dunphy and Lamb (2014) showed that, passing a mode-one IT through a mode-one baroclinic eddy results in

IW scattering into higher modes – with implications ~~to~~for the redistribution of energy into smaller-scales. Moreover, their proximity with a major fresh water source also conveys an additional interest with possible implications for climatology, biology and engineer-related fields of study – owing for instance to changes in near-surface stratification (i.e. mixing) or substantial discharges of suspended sediments (see Johns et al., 1998; Almeida-Filho et al., 2005).

The present study is therefore aimed at a first account of the full two-dimensional horizontal structure of the ISW field off the Amazon shelf. The remainder of the paper follows with Section 2 ~~by~~ describing the new satellite observations of short-period ISWs in the study region, together with preliminary interpretations of their main features and a statistical analysis of their horizontal structure. Section 3 follows with discussions on the generation and propagation. A summary and some concluding remarks are presented in the final section of the paper.

2. SAR Imagery Analysis

Synthetic Aperture Radars (SARs) have proven very useful amongst other means of satellite imagery, which are typically used to survey ISWs, owing mainly to their extensive field of view along with detailed spatial resolution. In particular, Wide-Swath acquisitions (WS) from the Envisat-ASAR have nominal spatial resolutions of 75 m, while covering large areas of approximately $400 \times 400 \text{ km}^2$. This makes them ideal to observe multiple packets of large ISWs, which are usually separated by semi-diurnal wavelengths (of the order of 100 km). Indeed, the ISW sea surface manifestations depicted in SAR imagery are now widely documented in the literature, and are essentially a result of hydrodynamic modulation of the sea surface

roughness and wave breaking, provided that the wind speeds are not excessively strong (see e.g. Alpers, 1985 and da Silva et al., 1998; Kudryavtsev et al., 2005).

A representative dataset was assembled based on 17 SAR images, all of which belonging to the Envisat-ASAR and acquired in WS mode ([all showing evidence of ISWs](#)). These range from 2004 to 2012 with no seasonal preferences being considered, since favorable [stratification conditions for ISWs occur year-round in equatorial regions year-round stratification conditions for ISWs are typical in equatorial regions](#) (see Table 1 for more details). The dataset was also chosen to provide a fair coverage of the semi-diurnal (M_2) and fortnightly tidal cycles. This is best seen in Fig. 2, which presents tidal heights for each acquisition (as listed in Table 1, according to Egbert and Erofeeva, 2002), and therefore offers an overall view of the dataset with ISWs being observed both in neap and spring-tides. Therefore, unlike other regions where high ISW activity is imaged mostly during spring tides (e.g. the Mascarene Ridge, see da Silva et al., 2011), these observations are not clustered around some particular phase of the flood-ebb tidal cycle (the latter effect being a consequence of sun-synchronous orbits; e.g. Valente and da Silva et al., 2009).

Fig. 1 presents a comprehensive view of the study region, together with a composite map resulting from the 17 acquisitions listed in Table 1, and whose envelope is shown by the black dashed line. For clarity, only the strongest ISW signatures were considered and only the leading wave is depicted for each ISW packet (in black solid lines, with a total of 59 packets being represented). A well-organized ISW field can be seen running along two major pathways apparently associated with distinct hotspots, which have been labeled A and B close to the 200 m isobath (representative tidal ellipses have been drawn for reference). The wave surface

manifestations are fairly aligned with the steep slopes on either side of a small promontory approximately at 44 °W and 0 °N, and therefore appear to emanate from localized, and yet distinctive, regions of the shelf-break (to be further discussed in Section 3).

Typical SAR views of the study region are given in Fig. 3. Both case studies reveal characteristic sea surface manifestations of intense ISWs appearing more than 500 km away from the nearest coastline (see insets for locations). Further inspection of the SAR signatures, in the same fashion as Thompson and Gasparovic (1986), indicates multiple wave packets propagating offshore and into the open ocean with a strong northeast component. The average inter-packet distances range from 121 to 140 km, which are typical wavelengths of long (semi-diurnal) ITs of the fundamental mode. Therefore, ISW packets have been labelled assuming generation in consecutive semi-diurnal cycles, and appear in sequence beginning from 0M₂ to 2M₂ in Fig. 3b, and to 4M₂ in Fig. 3a. These average inter-packet distances can be converted into mean propagation speeds, since a semi-diurnal generation is being assumed. In this case Fig. 3 reveals these waves to be ranked amongst the fastest ever recorded with mean values ranging from 2.7 to 3.1 m/s (for locations B and A, respectively).

Corresponding normalized backscatter intensities (i.e. $(I-I_0)/I_0$) taken across representative ISWs in Fig. 3 are presented in Fig. 4, which are characteristic of the study region. We note that, profiles were taken in the same fashion as [da Silva et al. \(2011\)](#), with backscatter intensities (I) being computed along a rectangular transect perpendicular to the wave packet, and normalized by some unperturbed mean value (I_0) usually taken ahead of the ISWs (see Fig. 3b for an example). We further note that, both case studies in Fig. 3 were selected to have low wind speed

components along the directions of wave propagation. According to Fig. 4, large-scale individual waves are seen propagating in the deep ocean, with horizontal scales (e.g. given by L in, see also Fig. 4a) of approximately 5 and 10–12 km, for A and B respectively.

To further investigate the main characteristics of the ISW field, we now present a statistical analysis, done for a sample of 59 different packets collected from the 17 images listed in Table 1. Fig. 5 begins with a distribution of the along-crest coherence lengths (or crestlengths), and it can be seen that it is slightly skewed towards the lower end, with values ranging up to more than 200 km for waves coming both from locations A and B (see also Fig. 1). Overall, the majority of the observations are characterized by crestlengths around 150 km, which make them comparable with other large-scale observations in the world's oceans (see e.g. New and da Silva, 2002; Ramp et al., 2004; da Silva et al., 2011; Magalhaes et al., 2012). Fig. 6 is also indicative of the presence of large-scale ISWs as it shows the distribution of the characteristic soliton widths (L). These, which were estimated as illustrated in Fig. 4 (see also Fig. 3), as is typically done for sea surface manifestations of ISWs in SAR images (see e.g. da Silva et al., 1998). The length scale L is considered here to be a proxy measure of the horizontal dimension (along the propagation direction) of the individual solitary waves (along their propagation direction) (just as in Figs. 3 and 4a). We note that for this particular study region the majority of the case studies presented ISWs with positive backscatter variations from the unperturbed I_0 (as is defined in da Silva et al., 1998, see their Figure 3), and hence L was generally defined as the length where normalized backscatter intensities are greater than I_0 (see dashed lines and L in Fig. 4). The characteristic soliton widths found in Fig. 6 ~~se-values~~ are of the order of several km, with mean

distribution values amounting to 5 and 6 km for sites A and B, respectively, and hence an order of magnitude higher when compared e.g. with the Bay of Biscay, Massachusetts Bay, or the western Iberian Shelf (Azevedo et al., 2006; da Silva and Helfrich, 2008; Magalhaes and da Silva, 2012). Consistently with case studies analyzed in Fig. 3, the distribution in Fig. 7 also reveals that few waves are found per wave packet. In fact, their distribution is strongly biased towards the lower end, with the majority of values ranging from one to two solitons per packet (for both sites A and B). Finally, Fig. 8 presents a distribution concerning the main directions of wave propagation found in the study region, for both locations A and B. According to Fig. 8a, the waves are generally traveling to the northeast, with waves from location A having a slightly more northern component than those from location B (as is also seen in Fig. 3). In addition, the pathway from location B is somewhat narrower (i.e. between 45 and 60 °T) when compared with the wider range of directions for ISWs coming from location A (i.e. between 30 and 60 °T). Fig. 8b further reveals that this wider range of directions comes in fact from a seasonal variation. It shows that waves emanating from A between February and May seem to propagate steadily towards 30 °T, but those between July and December have more northeasterly directions between 30 and 60 °T. We note that, no seasonal variability was found for waves associated with location B.

The SAR imagery highlights several important results concerning the two-dimensional horizontal structure of the ISW field in this region. In particular the composite map clearly reveals the existence of different pathways associated with two nearby and yet distinct locations (labelled A and B in Fig. 1). Moreover, the satellite data reveals a seasonal variability in the directions of propagation for location A. Another puzzling feature is that, unlike typical observations from other

marginal seas, the ISW sea surface manifestations appear quite far from the nearest continental shelves – as far as 500 km in comparison with distances of approximately 100 km in other regions (e.g. in the Mozambique Channel, see da Silva et al., 2009). These and other results are to be discussed in the next section.

3. Discussions

We now discuss in more detail some of the results of the SAR image analysis. In particular, seasonal differences are investigated with respect to the NECC's variability, and the remote appearance of the ISW sea surface manifestations is discussed in [the](#) light of ~~an-possible~~ unusually large disintegration time-scale of the IT.

3.1 Seasonal variability in the SAR data

Internal waves emanating from A (Fig. 1) can be further divided into two subsets according to their season. It was found that waves between February and May had a different and more straightforward path when compared with those between July and December – with January and June being somewhat of transition months.

Fig. 9a highlights the seasonal variability between the internal waves. In green and blue colors are two representative case studies corresponding to generation site A, one dated 27 May 2009 (in green), and another dated 3 October 2011 (in blue). Note that, waves from location B are also shown for reference (those from Fig. 3b) – although they do not reveal significant seasonal variability and will no longer be discussed for that matter. These two case studies illustrate the differences in the ISW propagation characteristics, which were otherwise not apparent in the overall

view of Fig. 1. Apart from being refracted towards the east (also seen to less extent in Fig. 3a for November), the October case study shows an increase in the ISWs inter-packet separations, associated to an increase in their average propagation speeds (of more than 30% when comparing the 3.4 m/s in October with the 2.4 m/s in May), which also seems to enable them to penetrate further into the open ocean – note only leading waves are depicted (as in Fig. 01).

According to Fig. 09a, these differences appear to be related ~~to~~with the seasonal variability of the NECC, which is represented by the colored arrows marked along the waves' paths (in green and blue for May and October, respectively). These arrows are representative of the NECC within the bulk of thermocline (i.e. vertically averaged, see also Fig. 11), and were obtained from monthly mean climatological data (between 1980 and 2011) – provided from NOAA/OAR/ESRL PSD, Boulder, Colorado and available at www.esrl.noaa.html.

Between July and December (represented as the October data in blue) ISWs refract eastwards owing to the NECC, which acts not only to refract the waves, but also provides an additional current component in the along-ISW propagation direction, contributing to their observed increased propagation speeds and ultimately to their extended penetration farther to the northeast. At the same time the other regime in the NECC (i.e. from February to May and represented as the May data in green) is quite the opposite. The NECC does not flow east and refraction decreases as the currents weaken substantially during that period, to the point where currents reverse. Instead of contributing with an along-ISW current component, the NECC changes direction so that the flow ~~component~~ is now in the opposite direction, running nearly perpendicularly to the ~~of~~ ISWs direction of propagation, causing ~~ISWs~~ them to decelerate along their propagation path. Note also that, changes in the NECC are not

as strong when considering the ISWs associated with location B. That is probably because B generated waves only partially intersect the NECC ~~partially~~ (i.e. along its southern border) rather than running directly across its main core (see also Fig. 01).

Climatological data from other monthly means show similar results, except for January and December, which are transition periods between both seasons. However, the vertically averaged currents shown in Fig. 9a (of the order of 10 cm/s, see scaled arrow in black), are likely underestimated in the climatological dataset. This is important since we can compare the mean SAR-derived propagation speeds in Fig. 9, with those obtained from a standard boundary value problem (BVP) with appropriate boundary conditions of the form:

$$\frac{d^2\phi}{dz^2} + \left[\frac{N^2(z)}{(U-c)^2} - \frac{d^2U/dz^2}{U-c} \right] \phi = 0, \phi(0) = \phi(\theta - H) = 0 \quad (1).$$

Note that in BVP 1 (~~often termed~~ the Taylor-Goldstein equation), ~~is representing~~ the vertical velocity modal ~~vertical~~ structure, H is the local depth with z being the vertical coordinate (directed upwards from the surface where $z=0$), N is the Brunt-Väisälä frequency, and U is the current vertical profile along the direction of wave propagation, assumed to have phase speed c (see Miles, 1961; Smyth et al., 2011). Climatological profiles taken in appropriate locations along the waves' propagation paths were again used for currents and stratification (as shown in Fig. 9b and 9c and described in the previous section). Upon comparison, the 2.4 m/s estimated in May agrees reasonably well with the 2.3 m/s from the linear theory, but the 3.4 m/s in the October case study (see blue labels in Fig. 9) ~~are-is~~ well above the 2.5 m/s found from BVP 1. This is likely because in situ currents are stronger than their climatological mean values, which were used to compute U in BVP 1. In fact, that is precisely the case for the measurements reported in Brandt et al., (2002), which

reveal currents (in November) an order of magnitude higher than those in Fig. 9 (i.e. around 1 m/s, see their Figure 1). Solving BVP 1 with a proxy for the measurements in Brandt et al. (2002), results in phase speeds between 3 and 3.5 m/s, which are more consistent with the SAR. [Drifter observations along the NECC further confirm that typical values for April and November display the NECC reversing character, and also reinforce that corresponding monthly means for near surface currents are of the order of 1 m/s \(see Lumpkin and Garzoli, 2005\).](#) Therefore, despite displaying the NECC seasonal character (i.e. the reversing of the current's main direction between May and October), which seems essential in explaining the SAR refraction patterns, the climatological data may be under ~~estimating~~ ~~valuing~~ its true influence in the waves propagation speeds.

3.2 Semi-diurnal generation at the Amazon shelf-break

~~After discussing the seasonal variability in the SAR data, w~~We now turn to the sources of the ITs in the study region. The SAR dataset consistently reveals inter-packet distances comparable with mode-1 semi-diurnal ITs (i.e. around 130 km), regardless of the fortnightly tidal cycle, since ISWs are generated during neap and spring-tides (see Table 1). Despite the remote appearance of the ISW field (i.e. its unusually large distance from the shelf-break), its two-dimensional horizontal structure (best seen in Figs. 1 and 3) resembles that of the classical disintegration of an IT wave radiating from forcing bathymetry, just as in many other ISW hotspots like in the Bay of Biscay and the South China Sea (see e.g. New and da Silva., 2002; Buijsman et al., 2010b).

This hypothesis is further examined by quantifying how much of the surface tide is converted into internal motions at the tidal frequency by forcing bathymetry, and the corresponding propagation of the ITs. Conversion rates (C) and energy fluxes (F_E) are particularly useful parameters in numerical models accounting for ITs, and have been found to be good indicators in global circulation models and in several other independent studies (see e.g. Gerkema et al., 2004; Shriver et al., 2012; Jeon et al., 2014; Kang and Fringer, 2012). Despite detailed numerical simulations being beyond the scope of the present study, we note that preliminary results (following from Buijsman et al., 2015) proved very useful in identifying IT sources and their spatial distributions. These numerical simulations were useful for guiding satellite acquisitions, which were requested specifically to survey ISWs and their sources in an unusually large study region.

Modelling data corresponding to a series of simulations performed with the 3D global HYbrid Coordinate Ocean Model (HYCOM, see Bleck 2002) were then examined to assess the main characteristics of the IW field energetics—~~which will hereafter be referred to as the 6.1 solution.~~ HYCOM is a realistically forced ocean model, with both tidal and atmospheric forcing, and hence including the ocean's major current systems, such as the NBC or the NECC. The simulations used in this study refer to an annual period (from October 2011 to September 2012) and their tidal forcing includes the largest semi-diurnal and diurnal constituents (i.e. M_2 , S_2 , N_2 , K_2 , and K_1 , O_1 , P_1 , Q_1), where a nominal horizontal resolution is set to $1/12^\circ$ at the equator with 32 layers in the vertical direction. This particular dataset is computed in a similar fashion as ~~solution 18.5~~ in Buijsman et al. (2015), whose detailed description may be found in Metzger et al. (2010), and Shriver et al. (2012) and in Buijsman et al. (2015). ~~We note that the 6.1 results presented here refer to an~~

~~annual period (from October 2011 to September 2012) and its tidal forcing includes the largest semi-diurnal and diurnal constituents (i.e. M_2 , S_2 , N_2 , K_2 , and K_1 , O_1 , P_1 , Q_1).~~ According to the SAR data we seek *a priori* to establish a semi-diurnal generation for the IT, whose energy budget is therefore extracted from HYCOM simulations with a cut-off period between 9 and 15 hours. These band-pass filtered data are then used to compute the main IW field energetics, which will be discussed in the following paragraphs.

Fig. 10 displays the depth integrated and time-mean (i.e. over one full year) conversion rates (C , left panel) and energy fluxes (F_E , right panel) for the semi-diurnal tides in the ~~6.4~~ [HYCOM simulations](#) described above, as is usually done to obtain a two-dimensional horizontal view of the IT generation field (see e.g. Kang and Fringer, 2012 and Buijsman et al., 2015). The maps presented in Fig. 10 are consistent with the SAR data (i.e. composite map of ISW crests in black solid lines), since strong ITs are generated over the steep slopes of the Amazon shelf-break, which then propagate along two major pathways into the open ocean. As expected, the strongest conversion rates in the left panel are restricted to the shelf-break, running along the 200 m isobath, and decaying rapidly either in the on-shore or off-shore directions. Local maxima are coherent to each side of a small promontory (or headland – see reddish colors at locations A and B, near 44 °W and 0 °N) from which the ITs emanate, and where C values yield as high as 0.15 Wm^{-2} . Hence, this region is comparable to other known IT hotspots such as the Mascarene Ridge (see e.g. da Silva et al., 2015 and their Appendix A). We also note that, the conversion rates along the 200 m isobath (where C is strongest) are fairly uniform in the year-long period (considering mean values for different months, instead of the annual mean). In fact, tidal transports are invariantly (i.e. year-round) large precisely

at A and B, thus contributing to consistently large conversion rates (note that C is proportional to the square of tidal transports). This is in agreement with ISWs being observed in this region throughout the year. Finally, a significant reduction in C occurs between A and B~~decay may be found in between A and B~~, thus emphasizing the distinction between these two separate generation sites, where tidal ellipses are large and currents run mainly in the cross-shelf direction – therefore ideal for elevated conversion rates from the surface tide into vertical baroclinic modes.

The depth-integrated energy fluxes (F_E in Fig. 10b), which essentially convey where the semi-diurnal IT energy is being held, are consistent with two major IT pathways emanating from A and B. Note that location A has somewhat stronger energy fluxes than B, in agreement with slightly longer ISW crests (see Fig. 5). Moreover, the propagation of the IT energy from the steep slopes of the ~~amazon~~-Amazon shelf-break is directed precisely towards the ISWs sea surface manifestations seen in the composite map (Fig. 10b). It is interesting to see that the majority of the ISWs are first seen when F_E begins to decrease rapidly into the open ocean. This is probably a consequence of the filtered HYCOM data (which shows only IWs at the semi-diurnal frequency) – hence, higher frequency ISWs may become undetectable in HYCOM filtered data.

3.3 Evolution and disintegration of the IT

IT generation at the shelf-break has proven to be consistent with an energetics analysis, and therefore a likely possibility in the study region. However, it is yet to be explained why the ISW evidence is found more than several hundred km away from their forcing bathymetry (i.e. the shelf-break) into the open ocean (see Fig. 1).

We therefore proceed to some considerations concerning the evolution and disintegration of the IT.

As described in Gerkema and Zimmerman (1995), ITs evolve according to a balance between nonlinear and dispersive effects. There are two fundamentally different sources for dispersion, which may arise either from rotational or nonhydrostatic effects. Longer waves like the IT are essentially controlled via rotational dispersion

($\mu = \frac{f^2}{\omega^2}$, with f being the Coriolis parameter and ω the wave frequency), which is much more efficient at the larger wavelengths and lower frequencies (ITs are essentially hydrostatic). The shorter ISWs, on the other hand, are much more sensitive to nonhydrostatic dispersion, which basically scales with $\delta = \frac{H_t^2}{L^2}$. Note that, H_t is a measure of the vertical extension of the waveguide and L is representative of the wave's horizontal dimension (see also Helfrich and Grimshaw, 2008).

To balance dispersion, however, nonlinearity also needs to be accounted for, which is typically parameterized as $\alpha = \frac{A}{H_t}$, where A is the wave amplitude (see e.g. Helfrich and Melville, 1990; Gerkema and Zimmerman, 1995). Nonlinear effects may be adjusted by changing either the wave amplitude or the thickness of the waveguide (usually taken as the thermocline). The wave amplitude is the most common choice in several theoretical and modeling studies, but in the present case it is the ~~characteristics of the~~ thermocline (i.e. H_t) that seems to be the governing parameter. Note that, a sharp density front was reported in Vlasenko et al. (2005), running closely along the ISWs pathway associated with location A (see their Figure 3.1 and 3.2 and our Fig. 1). This means that the thermocline's vertical extension

(i.e. the waveguide) will decrease when moving off-shore, as is shown by those authors and their in situ measurements.

In order to assess the effects of this environmental constraint ~~in~~on the disintegration process of the IT, we now present Figs. 11 and 12, which refer to a vertical density distribution running along the waves' propagation path. Fig. 11a refers to ISWs associated with location A during the May season (i.e. between February and May). It shows a vertical density section corresponding to the green dashed line in Fig. 9, and calculated from a monthly mean as discussed in Section 3.1 for the NECC.

The overall view in Fig. 11a is that of a marked density front with an associated narrowing of the thermocline, which decreases more than 100 m between the shelf-break and the open ocean. This result is actually quite similar to that provided by the in situ measurements in Vlasenko et al. (2005), and consequently builds on the consistency of our interpretations. Two ~~additional~~ vertical profiles (labeled P1 and P2 in Fig. 11b) help to characterize the waveguide vertical structure as it runs along the waves' propagation path, with P1 close to the ~~200-500~~ m isobath, and P2 just prior to the first SAR signatures seen in Fig. 1 (approximately at 3 °N). It is clear that the first evidence of the ISWs is closely related to the narrowing of the thermocline, which is simultaneously strengthened and brought closer to the surface (see P1 and P2 in Fig. 11b).

To further understand how this will influence the evolution of the IT, we now take H_t to be between the bolded isopycnals (in Fig. 11a), which were selected to be representative of the bulk of the thermocline seen in Fig. 11b. Fig. 12a presents this parameter and again confirms how the waveguide is significantly decreased in height as the waves evolve into the open ocean, whereas the remaining panels in

Fig. 12 illustrate its effect in the nonlinear and dispersion parameters. We note that a semi-diurnal IT was assumed in the calculations (meaning ~~is set for~~corresponds to a period of 12.42 h) with a nominal amplitude of 30 m (based on Figure 3.3 in Vlasenko et al., 2005 and needed to compute) and a horizontal dimension of 100 km (i.e. $L=100$ km in).

According to the results shown in Fig. 12, the decrease in the thermocline seen in Fig. 11 forces the nonlinear parameter to increase to higher values (green lines in panels a and b). Also shown in Fig. 12b is a disintegration envelope for the IT based on the numerical results presented in Helfrich and Grimshaw (2008), which show that the disintegration process increases monotonically within those limits. Note that in this case reaches this disintegration envelope ~~precisely~~approximately when the first ISWs appear in the SAR– i.e. between 3 and 4 °N (marked as transitions between solid and dashed lines). At the same time, changes in the thermocline vertical extension are inconsequent to the nonhydrostatic dispersion seen in Fig. 12c, which is otherwise expected close to zero for the long wavelengths of the IT (as already mentioned). The same applies to the rotational (or Coriolis) dispersion, but for different reasons, since it is the proximity to the equator that dictates the low values of μ .

Altogether, this means that increasing nonlinearity (with no dispersion to compensate) will force the long IT to steepen and seek balance at the smaller wavelengths, up to the point where it may disintegrate into short-scale waves (Gerkema and Zimmerman, 1995). During the course of disintegration the nonhydrostatic dispersion will increase quite substantially (since L drops by an order of magnitude) until balance may eventually lead to solitary-like waves. Note that, lines in the dispersion panels (in Fig. 12) are therefore limited to IT disintegration

point (taken as 4°N), beyond which both and would have to accommodate the scales of the ISWs.

The climatological conditions met for ISWs originating from location A, between July and December (i.e. concerning the October season in Fig. 9), are quite similar to the previous case. A monthly mean for the October density field (this time running along the blue dashed line in Fig. 9) also reveals a similar decrease in the vertical structure of the waveguide with Ht ~~performing-decreasing~~ accordingly (see blue line in Fig. 12a). All the remaining parameters also behave in the same fashion as discussed earlier. Note for instance that dispersion still remains ~~week~~weak, being indistinguishable from May, and therefore shown in black for both cases. Nonlinearity in Fig. 12b (in a blue line) is also indicative of IT disintegration and runs very closely with that of May (in a green line). In fact, regardless of the seasonal character in the NECC, the vertical density structure along the ITs emanating from A seems to be consistent with its disintegration somewhere between 3 and 4°N , and therefore in agreement with the earliest ISW signatures seen in the SAR.

The case for ITs emanating from B (see Fig. 1) is again very much alike. Following the same methodology as in Figs. 11 and 12, climatological data reveal a narrowing in the waveguide when moving along the waves' propagation path (see black dashed line in Fig. 9). The overall effect in the thermocline is as in the previous cases, with the waveguide becoming narrower, shallower, and slightly stronger (see black line in Fig. 12a). This means that, the nonlinear parameter will increase beyond the disintegration threshold (note dispersion is again very weak) with reasonable agreement with the SAR observations seen to the North of 2°N (see Fig. 12b with black dashed line again marking transition from the IT to ISWs).

In this study region the dispersion parameters are indeed very small. Note that, μ is $O(10^{-2})$ and β is $O(10^{-4})$, whereas mid-latitudes values for μ are of the order of 10^{-1} and β typically ranges from 0.2 to 30×10^{-3} (assuming typical length-scales for the IT) – see also Fig. 12c and 12d for comparisons. This means weak dispersion will be at play (see also Figure 6 in Gerkema and Zimmerman, 1995). Therefore, the evolution of an IT assumed to radiate from the Amazon steep slopes ~~will be~~ bound to break as soon as the nonlinear effects come into play. According to Helfrich and Grimshaw (2008) the disintegration envelope is set here to be between 0.12 and 0.16 (in Fig. 12), but as these authors suggest, it is not clear that this will always be the case. However, their trial run in the South China Sea is fairly close to our parameters, and we further note that shifting the disintegration boundaries does not change to the good agreement between the SAR evidence and the waveguide influence on the IT disintegration process.

Finally, we comment on the measurements reported in Brandt et al. (2002), who had recognized the possibility of IT generation at the Amazon shelf-break, but whose measurements were sparse and did not convey the tidal nature the ISW packets. In fact, their inter-packet distances were reported around 70 km whereas SAR imagery shows typical values to be over 100 km (see Fig. 3). Therefore, rather than assuming a classical disintegration of the IT into ISWs (as is proposed here), they reasoned instead that different generation mechanisms might be at work, which were based on a large ageostrophic component of the NECC centered around 4.5°N . We note however, that the SAR clearly illustrates the semi-diurnal nature of the ISWs, which was not available to Brandt et al., (2002). Furthermore, ISWs are seen in the SAR much farther south of 4.5°N , particularly in site B, and thus not necessarily related with the ageostrophic component of the NECC.

In light of the available evidence, a late disintegration of the IT provides a more consistent hypothesis, assuming a significant amount of energy is transferred from the surface tide to the IT. This was indeed confirmed by independent modelling provided by the HYCOM data, whose vertically integrated conversion rates are computed and presented in Table 2 for selected regions with large IT energy fluxes and known ISW activity (see e.g. Baines, 1982; Arbic et al., 2012; Jackson et al., 2012). Representative spatial averages were made for meaningful comparisons amongst the different regions – i.e. values noted under $\{C\}$. According to these results, the Amazon shelf shows large amounts of tidal energy being converted into the IT, which are of the same order of other known major hotspots of large ITs and elevated ISW activity, such as the Bay of Biscay, the Luzon Strait, or the Mascarene Ridge (Indian Ocean).

4. Summary and Conclusions

The Amazon shelf-break in the tropical West Atlantic is a powerful hotspot for intense ISW sea surface manifestations. SAR revealed, for the first time, their two-dimensional horizontal structure and yielded important results concerning their generation and propagation characteristics. Two distinct generation sites were identified off the slopes of a small promontory, each associated with a different pathway of ISWs (see Fig. 1), but both consistent with an energetics analysis exhibiting high IT conversion rates. SAR images revealed some unusual characteristics of these large-scale waves, such as their elevated propagation speeds and remote appearance several hundred kilometers away from the nearest forcing bathymetry. These large distances were explained in light of a late disintegration of

the IT, based on standard parameters governing the balance between nonlinear and dispersion effects, and the decrease of the waveguide (i.e. thermocline) thickness along a pronounced density front [\(which ultimately relates with the geostrophic NBC and NECC\)](#). [Finally, contributions from the NECC were also shown to be a likely explanation for the ISW seasonal propagation characteristics, using climatological data and SAR evidence, but a more substantiated approach would be welcome in order to confirm these results.](#) Despite the first account of the two-dimensional horizontal structure of the ISW field given in this paper, there are several important questions that remain [callusive](#), and would likely benefit from high resolution modelling and/or detailed in situ measurements.

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Table 1. List of all Envisat-ASAR images (in Wide-Swath mode) used in this study and ~~used~~depicted in Fig. 1, along with dates and times of acquisition and the respective generation site (i.e. A or B). Times after high water (TAHW – taken accordingly at sites A and B) are shown together with T~~idal heights and ranges are also shown~~, indicating each image's ~~their~~ tidal phase within the semi-diurnal and fortnightly tidal cycles, respectively.

Date	Time (UTC)	Site	Tidal height & range (m)
2012Feb22	12h42m	A	-0.83 & 2.25
2012Jan26	12h32m	B	+0.13 & 2.00
2011Dec16	12h35m	B	+0.78 & 1.74
2011Dec02	12h48m	A	+0.059 & 1.34
2011Nov02 – in Fig. 3a	12h48m	A	+0.61 & 1.40
2011Oct25	12h41m	A	-1.15 & 2.46
2011Oct14	14h44m	A	-0.88 & 1.88
2011Oct06	12h37m	B	-0.11 & 1.35
2011Oct03 – in Fig. 9	12h47m	A	+0.65 & 1.48
2011Aug10	12h26m	B	-0.85 & 1.74
2009May27 – in Fig. 9	12h41m	A	+0.01 & 2.53
2009May25	01h12m	A	-1.00 & 2.22
2009May11	12h44m	A	-0.36 & 2.04
2005Jul25	01h12m	B	-0.01 & 2.19
2005Apr14	12h10m	B	+0.60 & 1.32
2004Sep06 – in Fig. 3b	12h25m	B	+0.35 & 0.91
2004Jul25	00h44m	B	+0.80 & 1.60

<u>Date</u>	<u>Time(UTC)</u>	<u>Site</u>	<u>TAHW(h)</u>	<u>Height & Range(m)</u>
<u>2012Feb22</u>	<u>12h42m</u>	<u>A</u>	<u>+4.5</u>	<u>-0.83 & 2.25</u>
<u>2012Jan26</u>	<u>12h32m</u>	<u>B</u>	<u>+2.7</u>	<u>+0.13 & 2.00</u>
<u>2011Dec16</u>	<u>12h35m</u>	<u>B</u>	<u>+0.8</u>	<u>+0.78 & 1.74</u>
<u>2011Dec02</u>	<u>12h48m</u>	<u>A</u>	<u>-0.6</u>	<u>+0.59 & 1.34</u>
<u>2011Nov02 - in Fig.3a</u>	<u>12h48m</u>	<u>A</u>	<u>-0.1</u>	<u>+0.61 & 1.40</u>
<u>2011Oct25</u>	<u>12h41m</u>	<u>A</u>	<u>+5.8</u>	<u>-1.15 & 2.46</u>
<u>2011Oct14</u>	<u>14h44m</u>	<u>A</u>	<u>+5.8</u>	<u>-0.88 & 1.88</u>
<u>2011Oct06</u>	<u>12h37m</u>	<u>B</u>	<u>-3.3</u>	<u>-0.11 & 1.35</u>
<u>2011Oct03 - in Fig. 9</u>	<u>12h47m</u>	<u>A</u>	<u>+0.7</u>	<u>+0.65 & 1.48</u>
<u>2011Aug10</u>	<u>12h26m</u>	<u>B</u>	<u>+5.2</u>	<u>-0.85 & 1.74</u>
<u>2009May27 - in Fig. 9</u>	<u>12h41m</u>	<u>A</u>	<u>+3.1</u>	<u>+0.01 & 2.53</u>
<u>2009May25</u>	<u>01h12m</u>	<u>A</u>	<u>+5.8</u>	<u>-1.00 & 2.22</u>
<u>2009May11</u>	<u>12h44m</u>	<u>A</u>	<u>+3.9</u>	<u>-0.36 & 2.04</u>
<u>2005Jul25</u>	<u>01h12m</u>	<u>B</u>	<u>+2.9</u>	<u>-0.01 & 2.19</u>
<u>2005Apr14</u>	<u>12h10m</u>	<u>B</u>	<u>+1.1</u>	<u>+0.60 & 1.32</u>
<u>2004Sep06 - in Fig. 3b</u>	<u>12h25m</u>	<u>B</u>	<u>-0.3</u>	<u>+0.35 & 0.91</u>
<u>2004Jul25</u>	<u>00h44m</u>	<u>B</u>	<u>-0.3</u>	<u>+0.80 & 1.60</u>

Table 2. Representative spatial averages for the depth-integrated and time mean conversion rates {C} given in the [6.4 HYCOM simulation solution](#) for known hotspots with elevated ISW activity. See text for more details.

Known IT and ISW hotspots	{C} in (W/m ²)
Amazon west flank	≈0.1
Amazon east flank	≈0.1
Andaman Sea	≈0.1
Hawaiian Ridge	≈0.1
Estremadura Promontory	≈0.1
Bay of Biscay	≈0.2
Luzon Strait	≈0.2
Mascarene Ridge	≈0.3

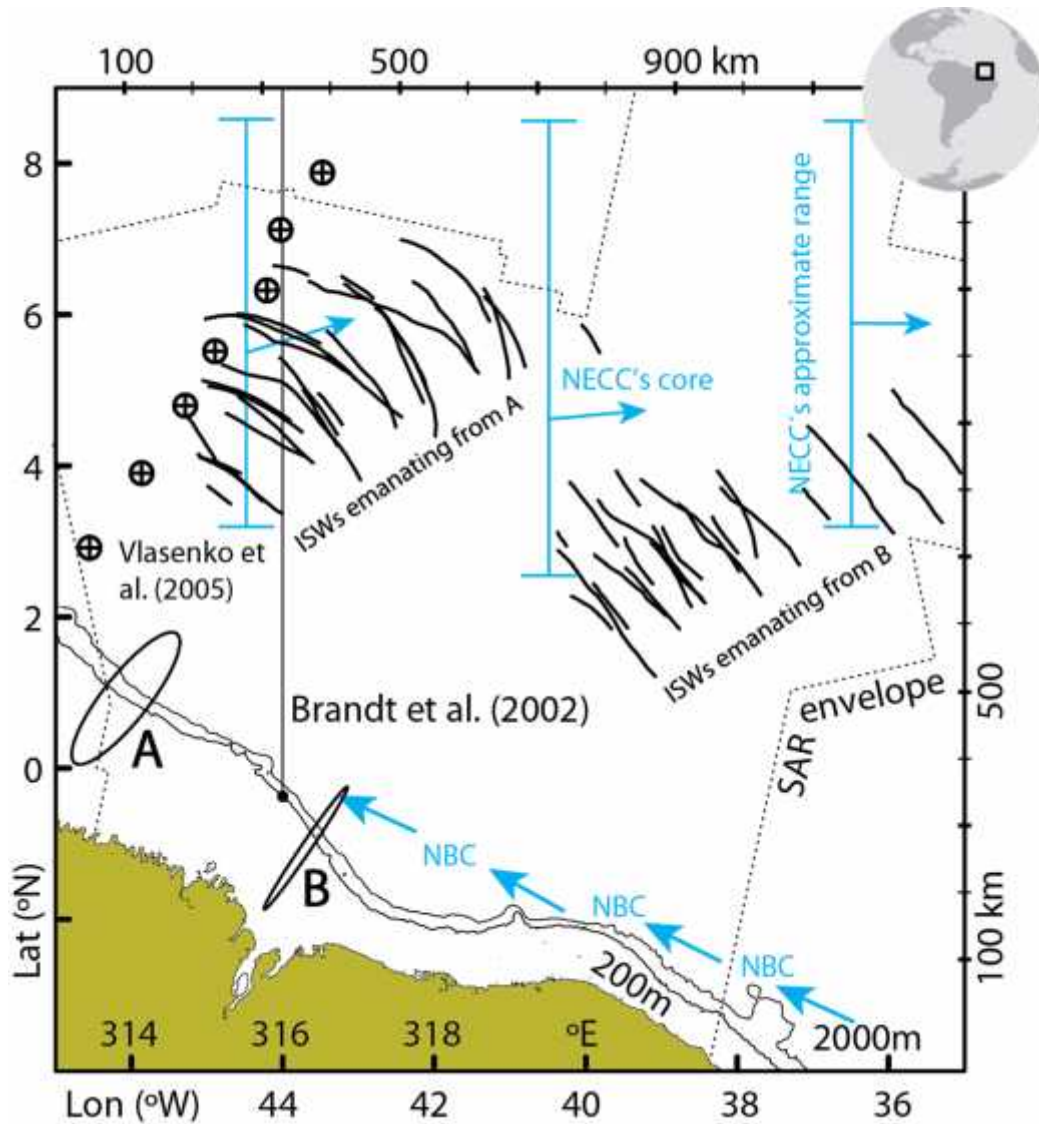


Figure 01:

ISW composite map (see inset for location in the tropical West Atlantic) with SAR observations limited to the dashed black envelope. Two major current systems are depicted in blue arrows: the North Brazilian Current (NBC), and the North Equatorial Counter Current (NECC) with vertical solid lines indicating its approximate range across the study region. Tidal ellipses in black are representative of sites A and B (scaled to [a maximum current of 0.5 m/s](#) in the across-shelf direction), which represent key areas of ISW generation on either side of a small promontory (near 44 and 0 °N). The shelf-break is highlighted by the 200 and 2000 m isobaths (in black lines). A meridional section along 44 °W used in Brandt et al. (2002) is shown for reference, together with the approximate location of a series of CTD stations (see circles with black crosses) discussed in Vlasenko et al. (2005).

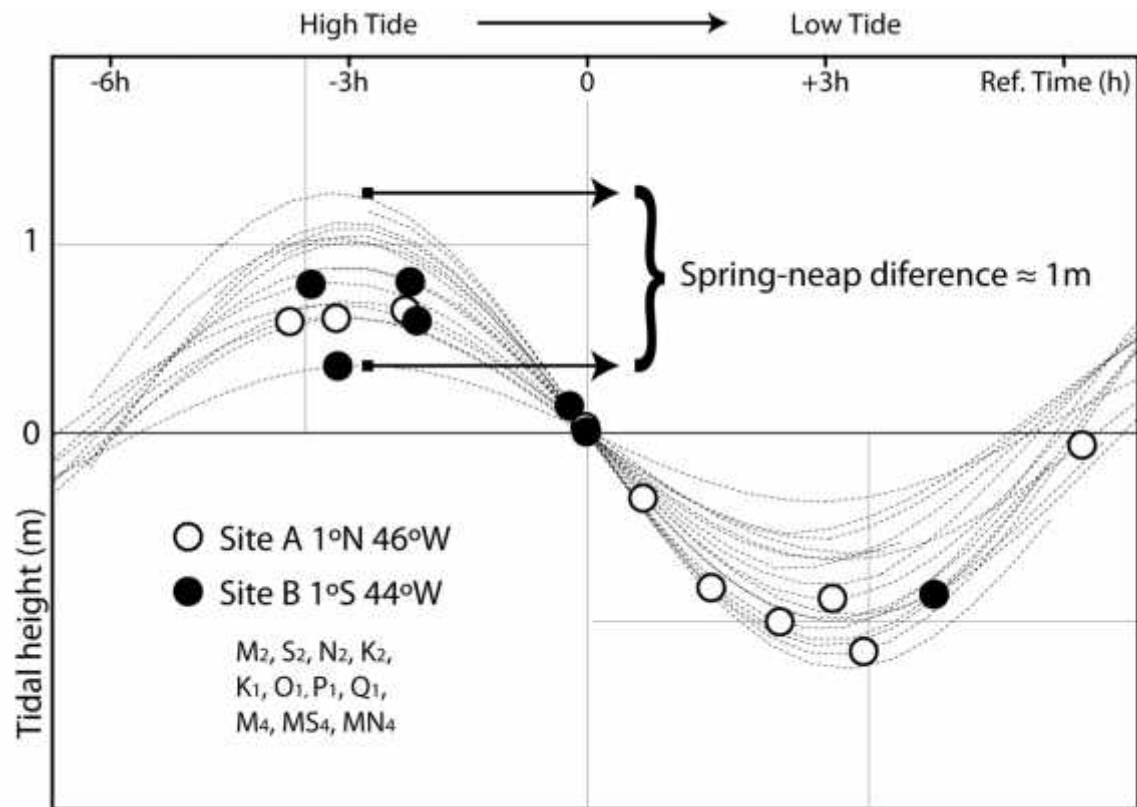


Figure 02:

Semi-diurnal and fortnightly tidal phases for all dates and times listed in Table 1 (each corresponding to an image acquisition), using the tidal constituents listed on the bottom-left corner, and computed correspondingly at locations A and B (represented by white and black circles, respectively). Note that, for a meaningful comparison, the time running in the horizontal [direction](#) is with reference to [the](#) transition ~~from-between~~ high tide ~~and to~~ low tide (i.e. when tidal heights change sign), and differences in the vertical range concern different phases of the spring-neap cycle.

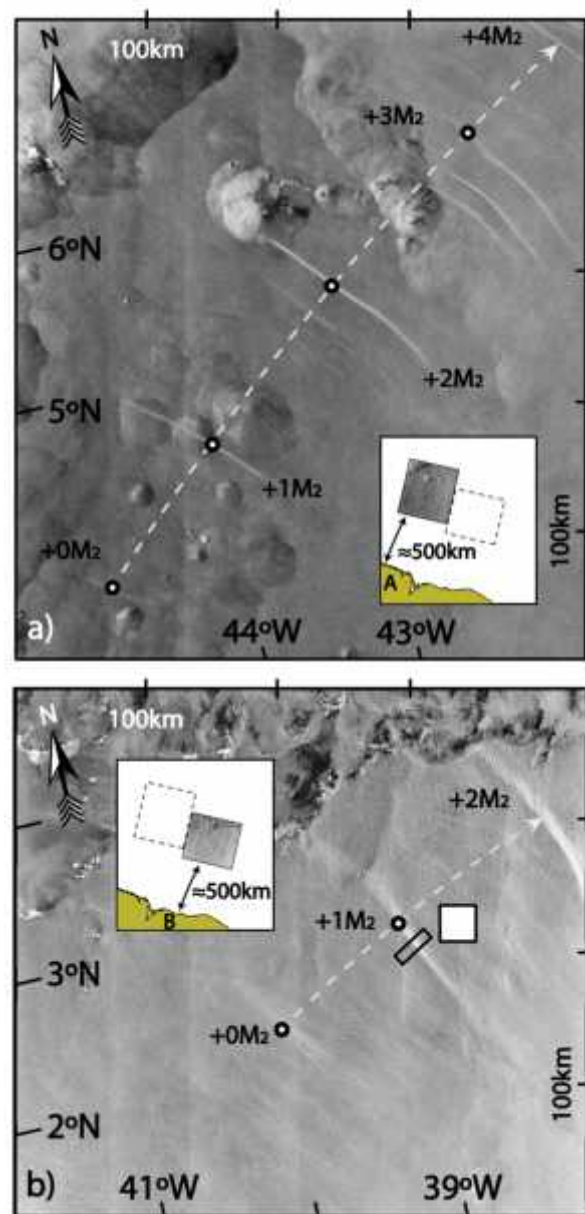


Figure 03:

a) Subset of a SAR image dated 20111102 (acquired at 12h 47m) showing ISW packets emanating from location A, and labeled sequentially from 0M₂ to 4M₂ (with 4M₂ being only partially imaged). An inset shows the image location with respect to the coastline, and the dashed line shows the location of Fig. 3b for comparison. b) As in top panel for a SAR image dated 20040906 (acquired at 12h 25m) and showing ISWs packets emanating from location B. The black rectangle and white filled square are examples of a radar backscatter transect and mean unperturbed background. A north pointing arrow and distance scales are added on the top and right-hand sides of both images.

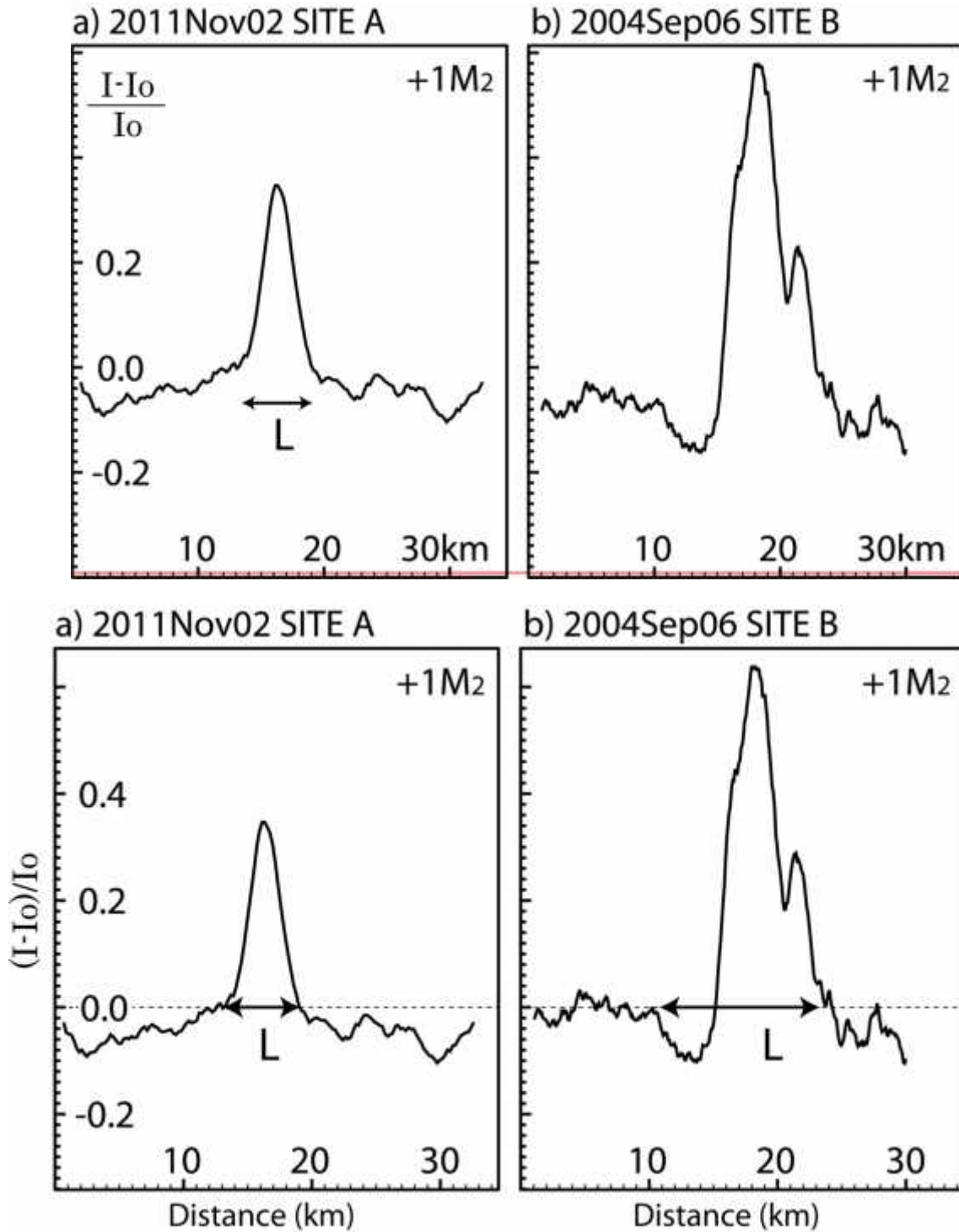


Figure 04:

Normalized backscatter intensity profiles corresponding to representative wave packets in Fig. 3a (left panel) and 3b (right panel). See Fig. 3b for a representative radar backscatter transect (used to compute I) and corresponding mean unperturbed background (used to compute I_0). A measure of the waves' horizontal dimension s is given in the left both panels (i.e. L , see text for more details).

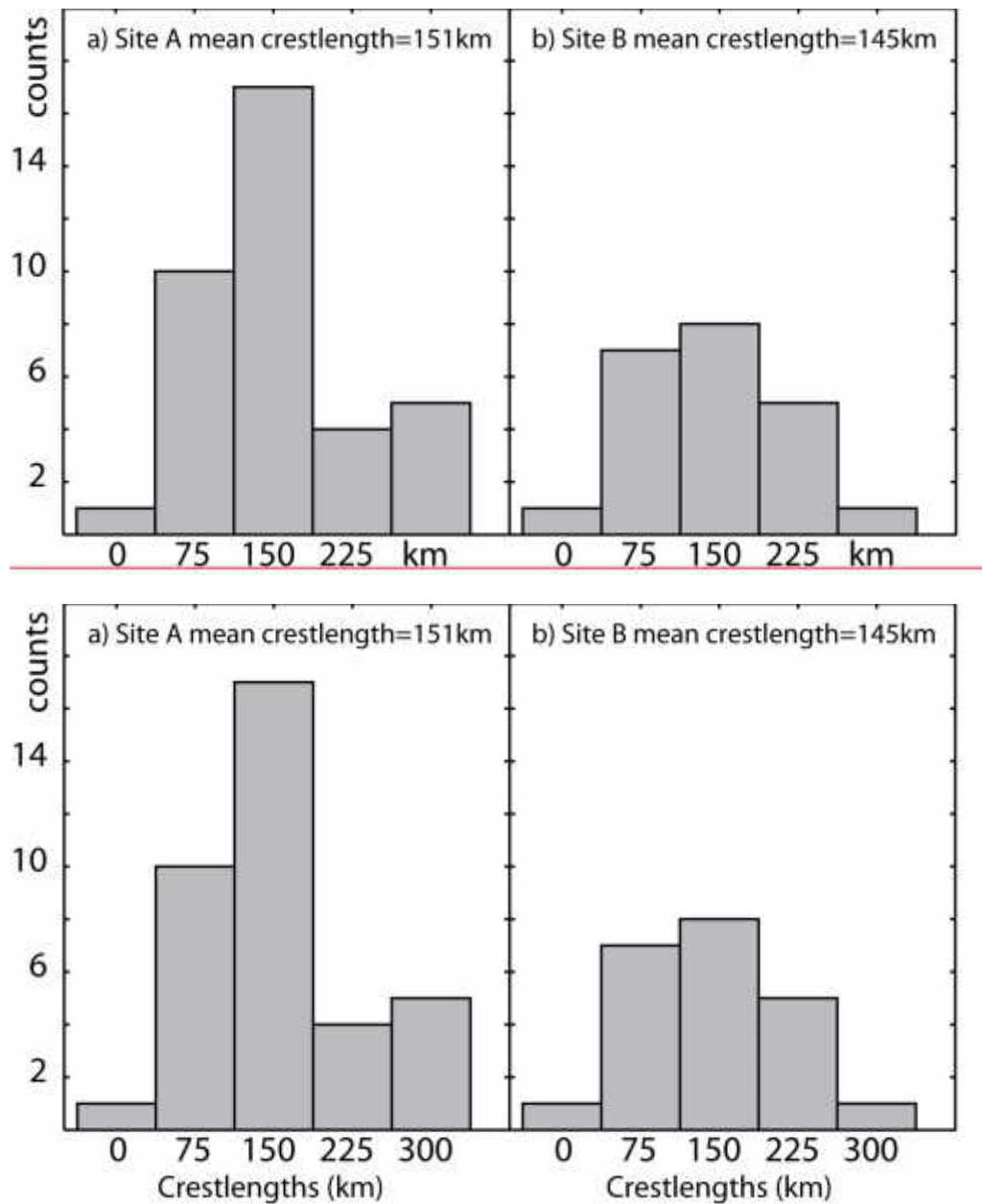


Figure 05:

Crestlength distribution for the leading ISWs in each packet.

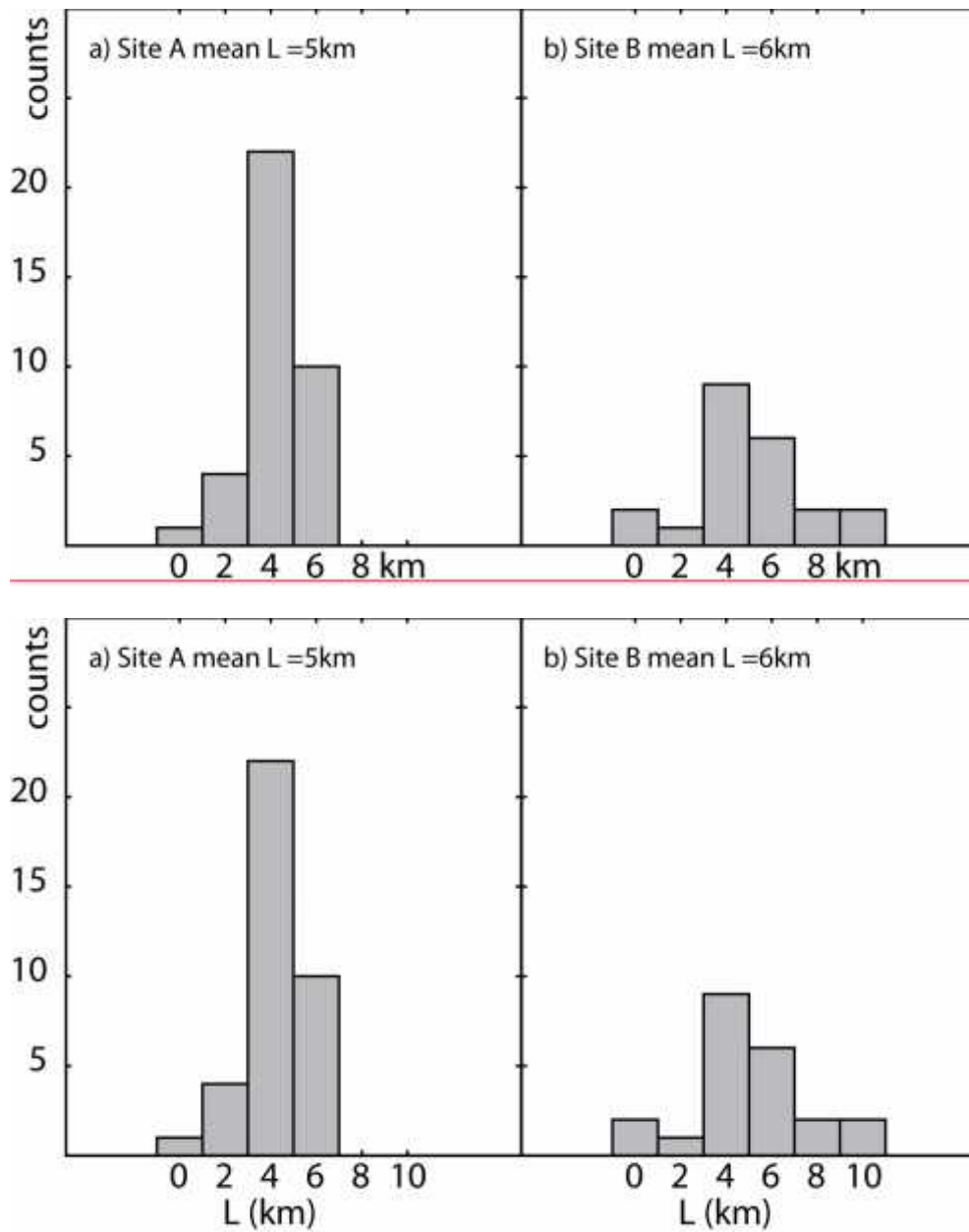


Figure 06:

Soliton width (L) distribution for the leading ISWs in each packet.

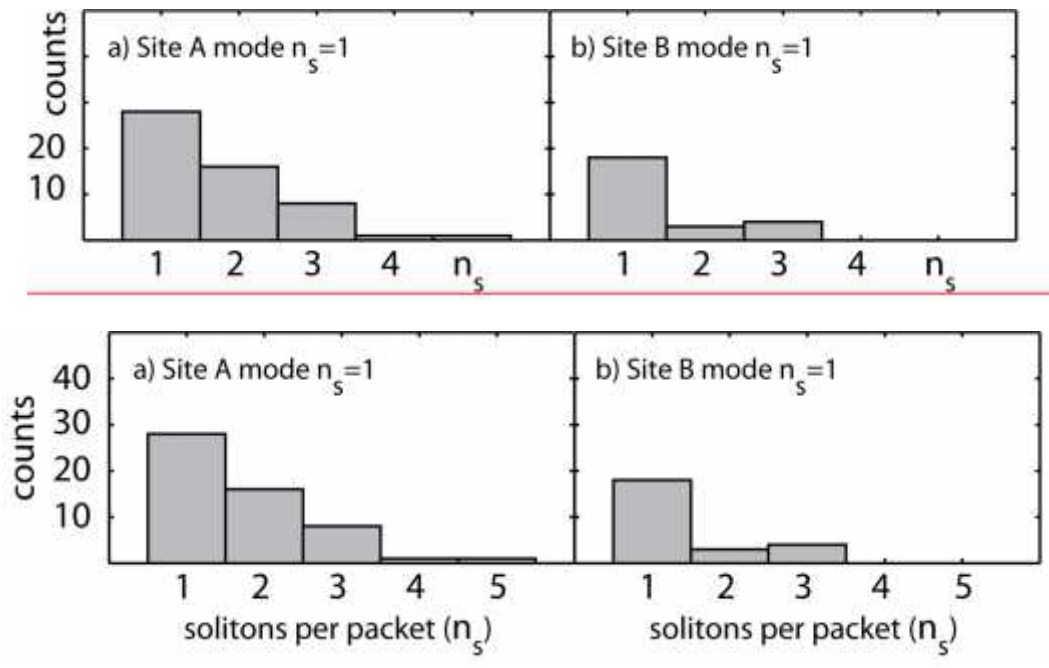


Figure 07:

Distribution of number of solitary waves per packet (n_s).

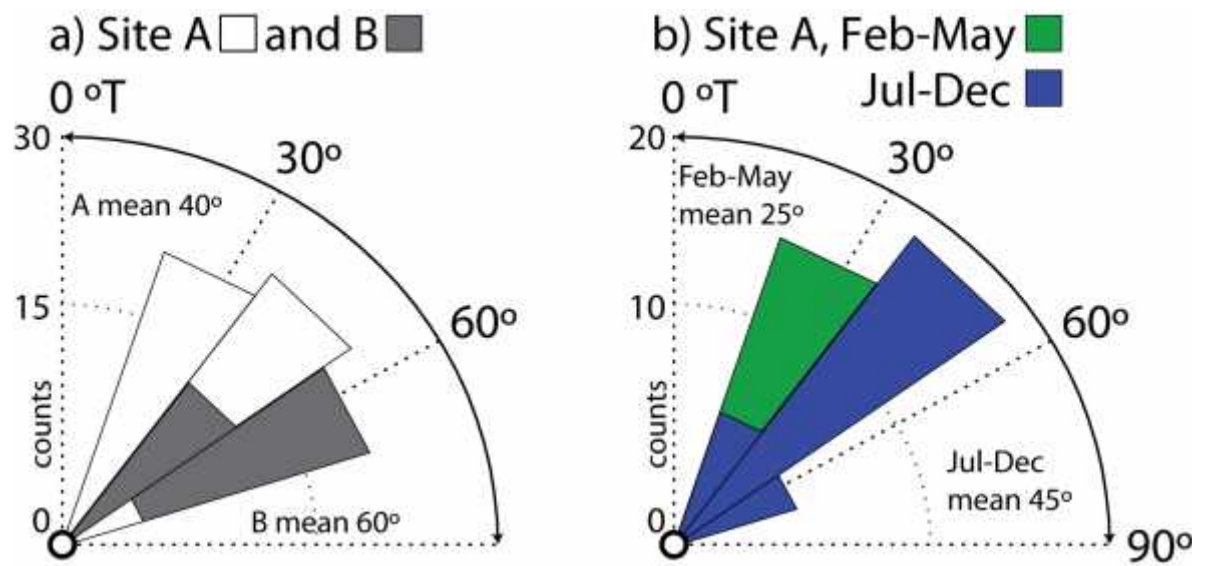
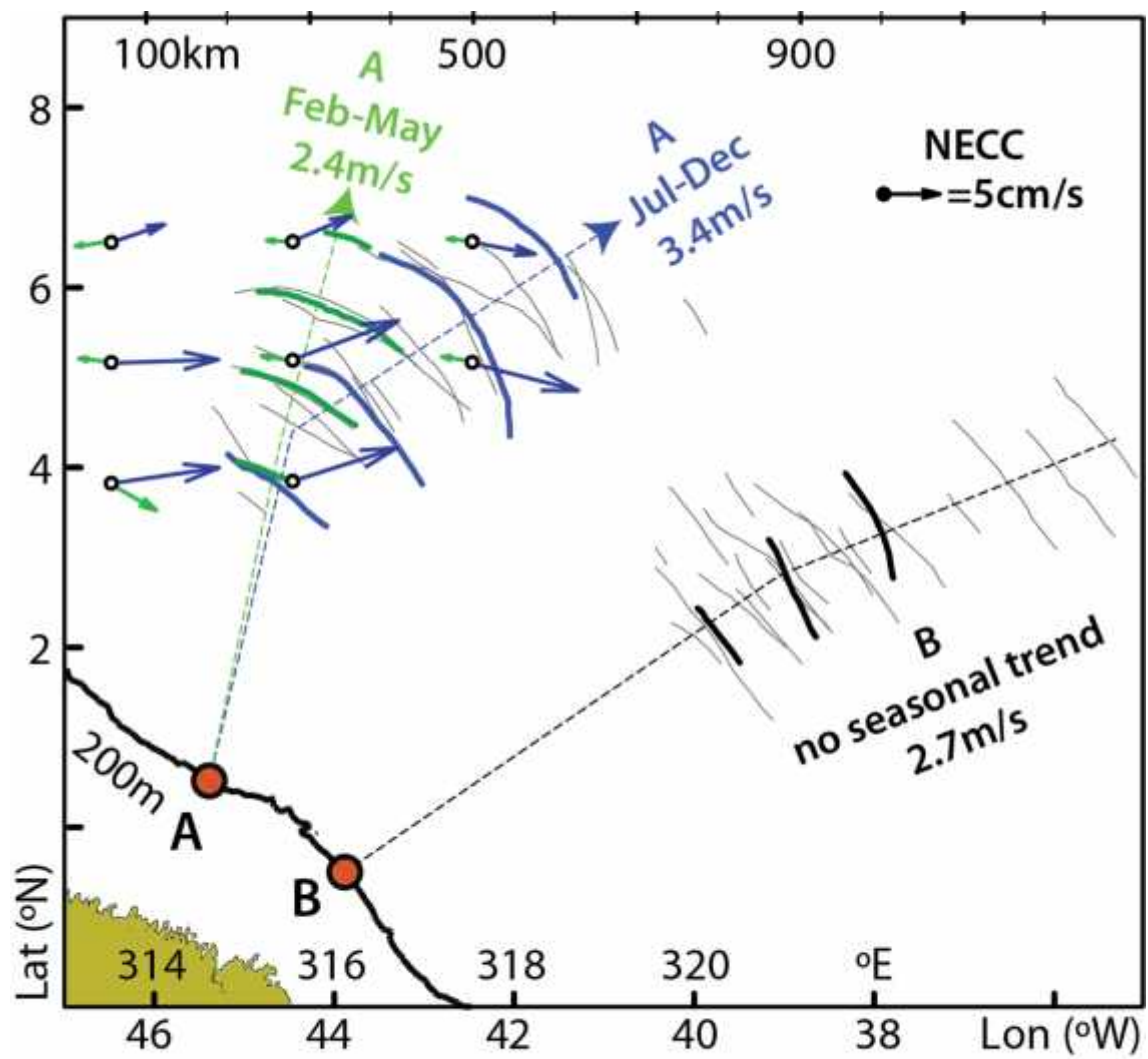


Figure 08:

Propagation direction distribution for each packet of ISWs.



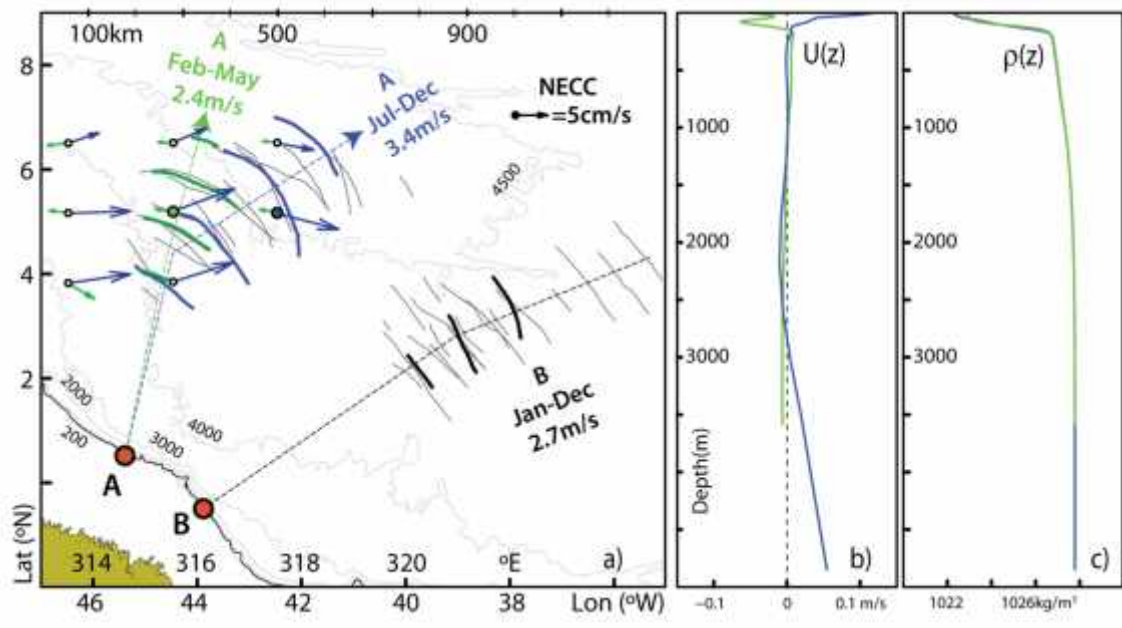


Figure 09:

a) Same composite map as in Fig. 1 is shown in thin black lines, with two case studies highlighted in thick colored lines, corresponding to location A and representative of two different seasons: from February to May (in green and dated 27 May 2009) and from July to December (in blue and dated 3 October 2011). Their averaged propagations speeds are also indicated along with idealized propagation paths (in thin dashed lines). Corresponding NECC monthly means (in green and blue for May and October respectively) are shown to depict its seasonal character along the ISWs propagation paths (see also scaled arrow in black). For reference, the case study shown in Fig. 3b is also highlighted in thick black lines, along with its mean propagation speed. Depth contours for 200, 500, 1000, 2000, 3000, 4000, and 4500 m are given in thin grey lines. Note the 200 and 1000 m contours are hardly distinguishable from the 500 m isobath (highlighted in a black thick line), and thus representative of a steep shelf-break. (b) Climatological monthly means for horizontal velocity vertical profiles projected along the ISWs idealized propagation path, taken at green and blue filled circles in panel (a) (for May and October, respectively). (c) Same as panel (b) for potential density vertical profiles. See text for more details. ~~As Fig. 1 but with two case studies highlighted, corresponding to location A and representative of two different seasons: from February to May (in green and dated 27 May 2009) and from July to December (in blue and dated 3 October 2011). Their averaged propagations speeds are also labelled along with idealized propagation paths (in dashed lines). Corresponding monthly means of the NECC (in green and blue for May and October respectively) are shown to depict its seasonal character along the ISWs propagation paths (see also scaled arrow in black). For reference, the case study shown in Fig. 3b is also highlighted in black, along with its mean propagation speed. See text for more details.~~

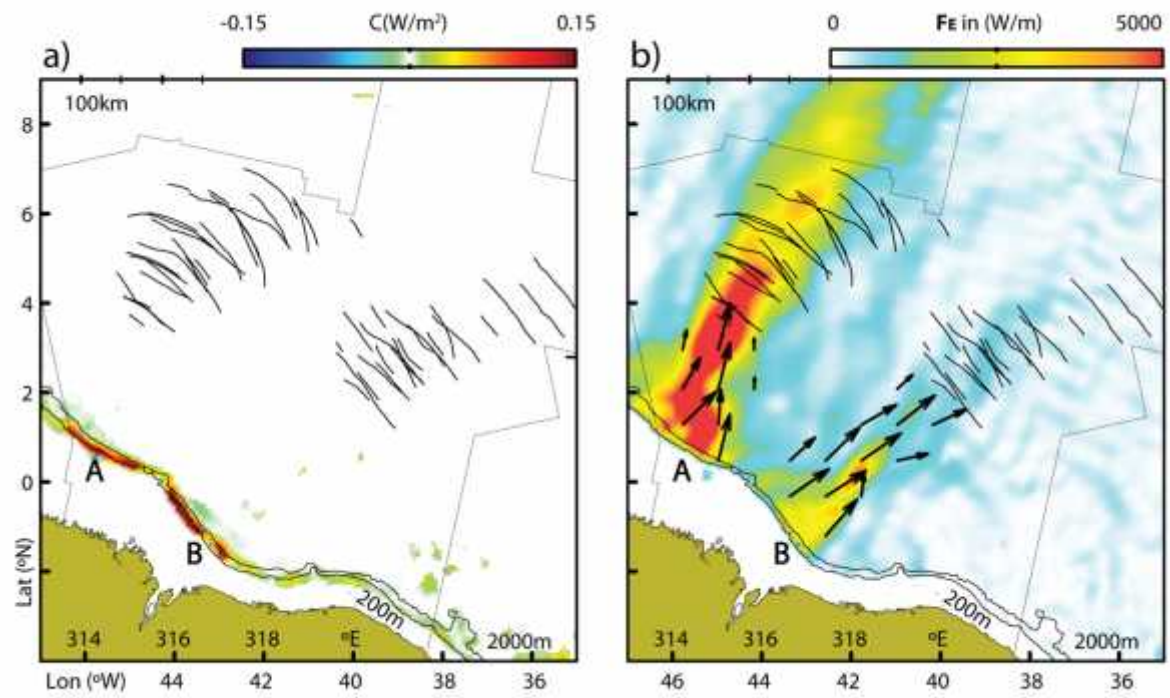


Figure 10:

As in Fig. 1 with depth-integrated and time averaged conversion rates (C , in left panel) and baroclinic pressure fluxes (F_E , in right panel) for the [6-HYCOM simulation](#) presented in the text (color scales on top of each panel). Two major IT hotspots are seen in A and B where the surface tide is converted to baroclinic vertical modes, which then propagate in two different pathways to the open ocean.

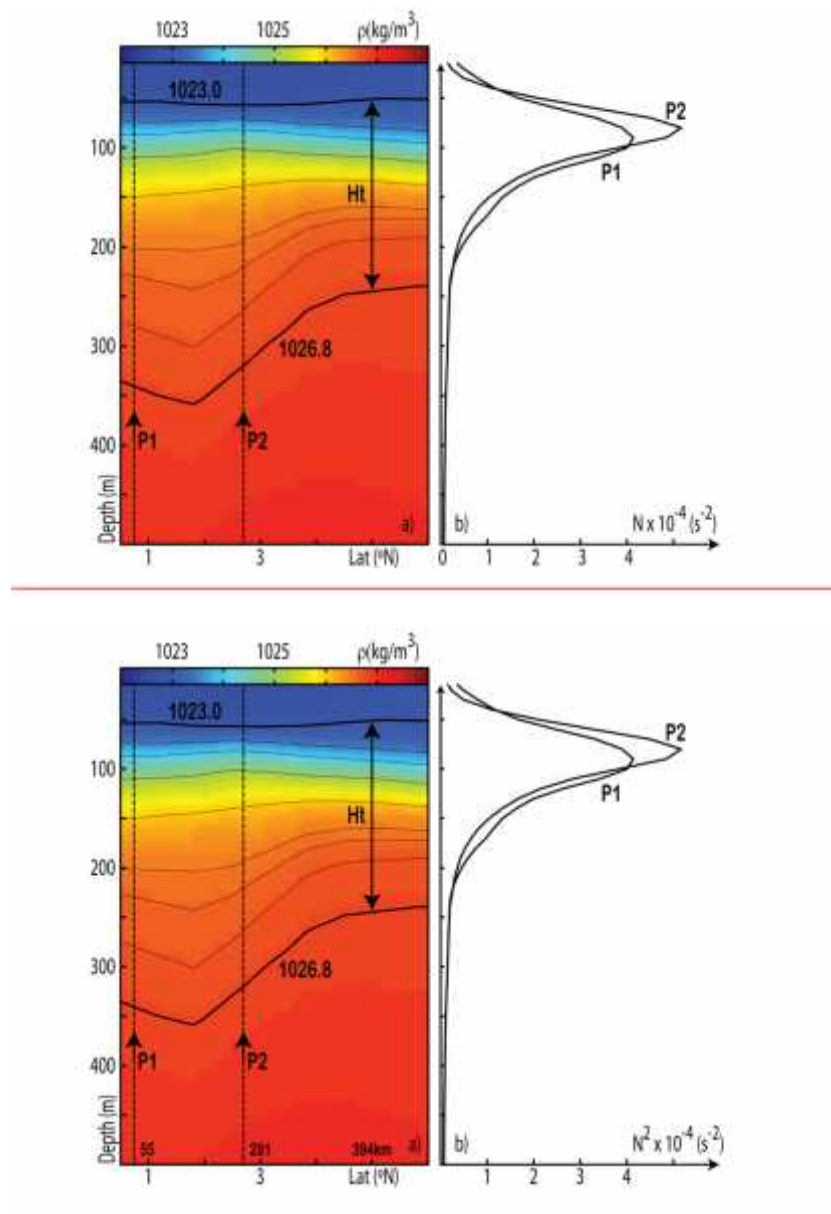


Figure 11:

a) May climatic vertical density section running along the green dashed line defined in Fig. 9 for generation site A. Selected isopycnals are also shown to highlight a density front. b) Brunt-Väissälä vertical profiles taken at P1 and P2 ~~in panel (a), representative of pre and post ISW sea surface signatures seen in the SAR (respectively).~~ Note that P1 is representative of the shelf-break where ITs are being generated, while profile P2 is set just prior to the SAR ISWs observations, where thermocline thickness decreases substantially.

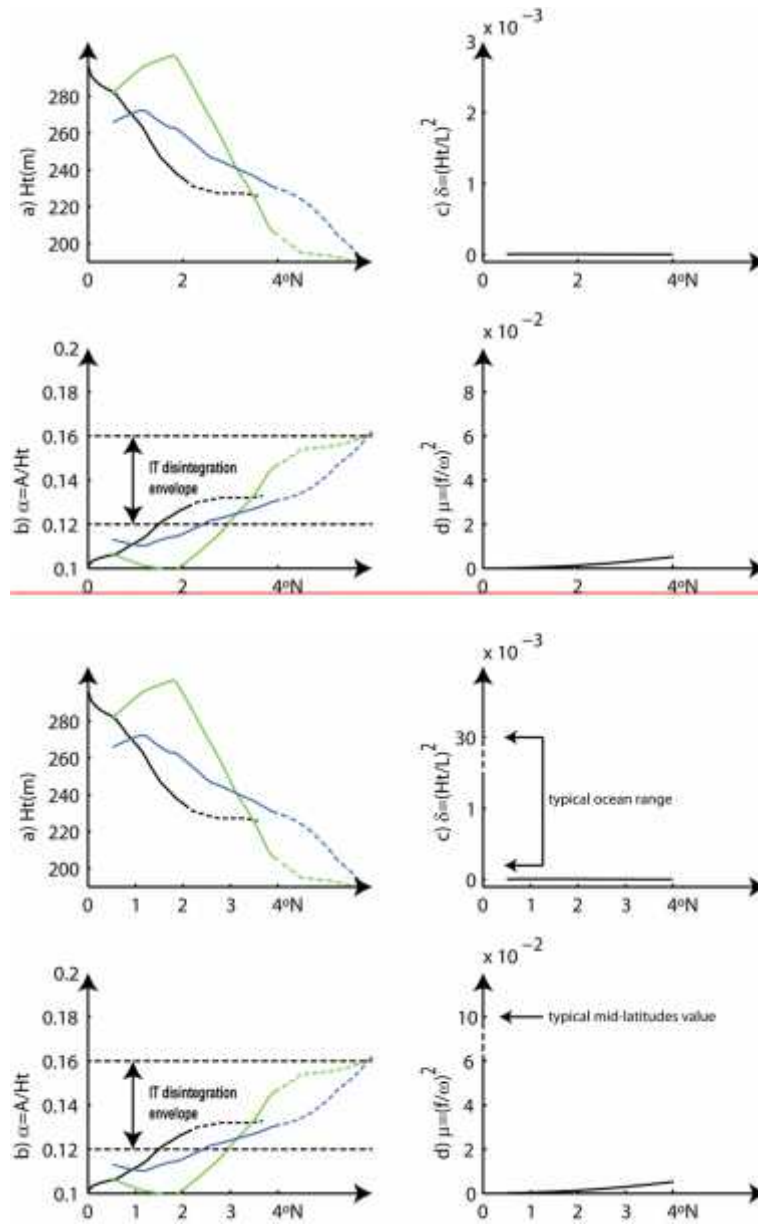


Figure 12:

(a) Vertical extension of the waveguides calculated along the waves' propagation paths (in green, blue and black, according to dashed propagation paths in Fig. 9). (b) Same as panel (a) for the nonlinear parameter computed assuming an IT amplitude of 30 m. (c) Nonhydrostatic dispersion for $L=100$ km. (d) Rotational dispersion for the semi-diurnal IT. Transitions to dashed lines in left panels indicate first SAR evidences of ISWs. Note broken vertical axes in right panels and typical ocean values showing comparatively weak dispersions in our study region. See text for more details.