

Dear Dr. Huthnance,

Thank you once again for your critical review of the manuscript. Your keen observations and insight always lead to improved clarity and presentation of the material, and I am delighted to continue the editorial cycle until your satisfaction is gained. Please find below a general response to the frequency-matching query, followed by responses to your other comments.

The physics of resonance is wonderfully expressed in the data of this paper, where the tidally-forced shelf-modes are hypothesized to provide the initial displacement and energy input that excite, and by virtue of their perpetual presence, sustain eigenmodes of the bay.

However, there still seems to be some confusion over the idea that a shelf-resonance can initiate and sustain the so-called bay-resonances. Specifically, that some form of frequency-matching is required. Please consider a piano or guitar string. When the string is displaced from equilibrium by the piano key it is not being forced at its resonant frequency. The fact that it resonates is a property of the system (length, tension, modulus, end-conditions). The middle-A string will always resonate at 440 Hz, even though the piano hammer or guitar pick is not exciting it at 440 Hz. This is true whether the string is struck with an impulse, or very slowly moved from its equilibrium position and released with zero impulse. Equally, one could consider a pendulum, or a mass suspended by a spring which are given an initial displacement (at zero frequency) and released. They will all oscillate at the resonant frequency determined by the physical parameters of the system.

Perhaps some confusion stems from the mathematical description of 2nd order physical systems and their expression of resonance. A common end-point of the ODE is to express amplification with a factor of the form $(\omega - \omega_0)^{-2}$ recognizing that driving a resonant system near the resonant frequency maximizes the response amplitude. A fundamental principle of communication and IT systems. However, this does not mean that a resonance can only be excited by driving it near the resonance frequency, as the pendulum, mass-spring, and vibrating string demonstrate, rather that frequency-matching maximizes the response.

An example of this is the well-known response of Hilo bay to tsunamis. As noted in the paper the tidally-forced shelf-mode at Hilo is 31 minutes. This mode is continually present and 'dominant' of water level variance at periods less than one hour under normal (tidal) conditions. When tsunamis with periods near 15 or 30 minutes strike Hilo, the amplification is impressive and is thought to account for the substantial damage and inundation observed at Hilo but not other Hawaiian locations, for example the 1946 and 1960 tsunamis. Thus the resonance is always present at small amplitude when forced by the tides alone, and excited to maximal amplitudes when forced near the resonant frequency.

As this point has been raised, and accepted, by one of the co-authors, by an independent reviewer prior to submission, by one of the reviewers, and yourself, it seems appropriate to include a brief section reviewing the essentials of resonance and its relevance to the hypothesis.

The two main matters are

a) Your response to reviewer 1 about models. I don't think numerical models are needed but I do think that the descriptions of modes, and the basis for identifying particular periods with particular mode forms, need to be clearer. There is a lot of terminology between which the distinctions are not clear. You refer (sometimes in reviewing others' work) to bay modes, harbour modes, shelf modes, shelf-resonances, transverse mode, edge-waves. How are these all distinct and to which ones are $\lambda/2$, $\lambda/4$ and $g\alpha/(\omega^2 - f^2)$ relevant? [I don't think there is any real difference between $\lambda/4$ and $g\alpha/(\omega^2 - f^2)$; only the form of cross-shelf depth profile assumed and whether f is taken into account.] I think some systematic introduction of terminology and description of forms would help to make this all clearer.

Thank you for this comment. The terms mode and resonance are used interchangeably, while the modifiers of shelf, bay, harbor are intended to qualify the spatial scales and domains. To clarify the nomenclature a paragraph has been added to the new Resonance section describing the usage.

Regarding the tidally-forced expression of Clarke & Battisti (shelf width $\approx g\alpha/(\omega^2 - f^2)$) in relation to Merians' formula, they are nearly equivalent however the Clarke & Battisti expression makes the implicit assumption that water velocities are modulated by f , that is they assume that the water velocities are driven by tidal forces that scale as f , whereas Merians' formula is purely geometric and invariant of latitude and assumed forcing. Nonetheless, as you point-out the effect of f is quite small and I agree with your assessment. Is there anything to be done about this in the manuscript? My preference would be to leave it as-is, as I find corroborating lines of inquiry satisfying and perhaps useful to the reader whom may not be aware of the different formulations.

b) This probably related to my (the editor's) confusion about what forms of oscillation are regarded as driving what other form. Another factor in this is the premature attribution to tidal forcing (in my opinion). There is a "progression" of expressions:

Lines 99-100 ". . . long period shelf-resonances driven by tides might constitute a more general excitation . . ." essentially as a hypothesis

Lines 148-9 "the tidally-forced shelf-resonances"

Lines 158-9 "the tidally-forced shelf-resonances"

Line 181 "with recognition of tidal energy as a driver of shelf-resonances"

Line 275 "shallow water tidal forcings"

Line 287 "tidally-forced shelf-modes are a plausible driver"

But I don't see evidence justifying all this until section 5, line 288 and after. Even then there has to be a question of how tides force motion at frequencies which don't seem to have any relation to tidal frequencies.

Thank you for this detailed observation. As above, frequency-matching serves to maximize the response of a resonant system. A pendulum illustrates that a zero-frequency displacement is sufficient to initiate a resonance. In order to clarify this, we have added a Resonance section outlining the essentials of mechanical resonance and our main hypothesis.

In relation to this, line 139 “These modulations are coherent with the tides” does not imply tidal forcing if the matter considered is (line 140) changing water depth etc. which occur during the semi-diurnal or diurnal tidal cycle. However, coherence with the spring-neap cycle does suggest tidal forcing. I think two separate ideas are being muddled here and should be separated.

Thank you for this observation. You are suggesting that the observed frequency modulations of these short-wavelength harbor resonances are not tidally-forced (energy supplied by the tides), whereas as the metamodes (long-wavelength shelf-resonances) are. I agree. We are simply attributing the observed modulations of these 1-5 minute period modes to changing water-depth/boundaries, which are tidally forced. At this point in the paper, the reader will not have encountered the material on metamodes and their spring/neap dependence, and we explicitly mention at the end of the paragraph that the driver is 'water levels'. That these short-period oscillations are not of direct tidal origin has been clarified in the text.

Other detailed suggested re-wording and comments (some related to the above. C.f. Referee 1, I think “seiche” in the singular is OK if used as an adjective, e.g. “seiche amplitudes”).

Line 18. “. . A thorough review of seiches is provided . .”

Line 22. “. . continuous seiches extend . .

Line 23. “. . 28-minute seiches in”

Line 32. Not “sustainment” but “sustenance” or “sustainability”

Line 36. “. . identified permanent seiching . .”

Line 44. “. . amplitudes increase to 10 cm typically . .”

Line 65. “. . seiche period at Colombo . .”

Line 76. “greatest” not “highest”.

Line 88. “. . excitation of seiches, . .”

Line 89. “. . seiches are indeed present . .”

Line 90. Maybe “at least widely occurring” rather than “global”.

Line 102. “bays/harbours”

Changed as suggested, except “bays/harbours” which I have left as “bays/harbors” to consistently use “harbor” throughout the paper. It can be changed everywhere to harbours if the editor prefers.

Lines 208-209. “shelf-modes appear to dominate the water level variance” – what is the basis for this statement? You don’t have data on the spatial form of oscillations.

The basis is the power spectra presented in figure 3. As noted in the caption “Triangles mark the tidally–forced shelf–resonance”, which is the connection to the spatial scale of the mode, and the maximal amplitudes at the shelf-resonances indicate that they “dominate the water level variance”. To be generally accurate this is not true when tidal oscillations are included, although they are not shown in the figure. Nonetheless it is useful to clarify this point and this has been done.

Line 215. “Monterey Bay seiching . .”

Changed.

Lines 218-220. What is the distinction between “transverse mode” and “shelf-resonance”? How long is “long period”? How is the 27.4 min mode to be considered as “a potential continuous forcing” if the forced oscillations are at a different frequency.

The shelf-resonance is a transverse mode in the sense that it is orthogonal to the shore, as is the fundamental transverse resonance of the bay itself. “Long period” indicates longer than wind-wave periods or harbor-modes, roughly 10 minutes in Monterey and have clarified this in the text. The frequency matching has been addressed in the new Resonance section.

Line 267. “shelf edge-wave” – do you mean an edge wave along the coast?

Yes, a prototypical alongshore edge-wave.

Line 279. The M2 tidal period is 745 minutes. Harmonics of this have periods $745/n$ minutes ~ 373 , 248, 186 . . minutes. Where do your periods 288, 199 minutes come from?

The 288 minute period is likely one, or a combination of, the MSK5, 3KM5, 2MK5 shallow-water constituents which have periods from 292 – 296 minutes. The 199 minute period would correspond to one of 3MN8, 2MN8, MSKO7 and/or 2MSO7 which span periods of 187 to 212 minutes. Keep in mind that the periods obtained from the power spectra (373, 288, 240, 199 min) are influenced by spectral leakage/averaging and the spectral binwidth which is not sufficient to resolve such closely spaced constituents such as MSK5, 3KM5, 2MK5.

Line 290. You label figure 5 as “Shelf-mode power spectral density” but I suppose in fact it is the spectral density of the measured elevation at a single point. Likewise line 398 and maybe elsewhere.

Yes, these are power spectral estimates of the water level oscillation amplitude measured at the tide gauges at the frequency of the corresponding shelf-mode. The text has been modified to make this clear.

Line 333. 35-50% (only) suggests that there is some other forcing as well.

Yes, the last paragraph of the conclusion section considers this in some detail. I like to end papers with open questions and would prefer to leave that material at the end rather than moving it here. If you feel it warranted a similar mention can be made here.

Line 340. The hypothesis in lines 100-101 was expressed rather differently (“shelf-resonances” not “the shelf-mode”) and referred to tidal forcing.

Thank for this notice. This has been changed to maintain consistency.

Line 342. Are the “transverse and shelf-modes” really different?

As above, the shelf-modes are transverse modes, but tidally-forced, whereas the transverse bay-modes are forced by the shelf-mode. Here, the “and shelf-modes” is redundant and has been removed.

Lines 362-374. There is the same question as for lines 220-221: how can a mode at one frequency (or period: 27.4 min) force oscillation at a different frequency. (This is not to doubt the energy argument but frequency match is another criterion).

Please see the new section Resonance, and keep in mind that we are considering a continuum system (water) not an independent set of uncoupled oscillators.

Line 378. “seiches” (twice).

Line 380. “seiching”

Line 382. “seiches”

Changed.

Continuous seiche in bays and harbors

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Abstract. Seiches are often considered a transitory phenomenon wherein large amplitude water level oscillations are excited by a geophysical event, eventually dissipating some time after the event. However, continuous small-amplitude seiches have been recognized presenting a question as to the origin of continuous forcing. We examine 6 bays around the Pacific where continuous seiches are evident, and based on spectral, modal and kinematic analysis suggest that tidally-forced shelf-resonances are a primary driver of continuous seiches.

1 Introduction

It is long recognized that coastal water levels resonate. Resonances span the ocean as tides (Darwin, 1899) and bays as seiches (Airy, 1877; Chrystal, 1906). Bays and harbors offer refuge from the open ocean by effectively decoupling wind waves and swell from an anchorage, although offshore waves are effective in driving resonant modes in the infragravity regime at periods of 30 s to 5 minutes (Okiihiro and Guza, 1996; Thotagamuwage and Pattiaratchi, 2014), and at periods between 5 minutes and 2 hours bays and harbors can act as efficient amplifiers (Miles and Munk, 1961).

Tides expressed on coasts are significantly altered by coastline and bathymetry, for example, continental shelves modulate tidal amplitudes and dissipate tidal energy (Taylor, 1919) such that tidally-driven standing waves are a persistent feature on continental shelves (Webb, 1976; Clarke and Battisti, 1981). While tides are perpetual, seiches are often associated with transitory forcings and considered equally transitory. A thorough review of ~~seiche~~-seiches is provided by Rabinovich (2009) wherein forcing mechanisms are known to include tsunamis, seismic ground waves, weather, non-linear interactions of wind waves or swell, jet-like currents, and internal waves. Excepting strong currents and internal tides, these forcings are episodic and consistent with a perception that

seiches are largely transitory phenomena. However, records of continuous ~~seiche~~ seiches extend back to at least Cartwright and Young (1974) who identified near-continuous ~~28-minute~~ 28-minute seiche in Baltasound, Unst and Lerwick in the Shetland Islands over a 16 week period in 1972. Their source hypothesis consisted of long waves from the North Sea trapped as edge-waves along the island shelf, and they noted large seiche amplitude modulations from fast-moving meteorological fronts.

Internal waves are known to influence seiches as demonstrated by Giese et al. (1990) who analyzed a 10 year time series of six minute data at Magueyes Island, Puerto Rico, noting distinct seasonal and fortnightly distributions of shelf-resonance and seiche amplitude suggesting that stratification and its influence on internal waves generated by barotropic tides are important components of the observed seiche variability. Subsequent work led Giese et al. (1993) to conclude that in locations where strong internal waves propagate to coastal regions, that seiche ~~sustainment~~ sustenance is possible, and to question “is this a general phenomenon to be expected wherever large coastal seiches are found, or is it specific to certain locations where large internal waves are known to occur?”

Golmen et al. (1994) studied a coastal embayment near the headland Stad in northwest Norway and identified ~~a-permanent-seiche~~ permanent seiching superimposed on the semi-diurnal tide concluding that tidal forcing was the only possible energy source of the observed oscillations. Continuing their exploration of internal waves and seiching, Giese et al. (1998) examined harbor seiches at Puerto Princesa in the Philippines finding that periods of enhanced seiche activity are produced by internal bores generated by arrival of internal wave soliton packets from the Sulu Sea. However, as one would expect from soliton excitation, their analysis suggests that these seiches are not continually present.

Woodworth et al. (2005) studied water levels at Port Stanley in the Falklands islands identifying continuously present 87 and 26 minute seiches with amplitudes of several centimeters. They also noted that amplitudes ~~rise-to-the~~ increase to 10 cm ~~level~~ typically once or twice each month with no obvious seasonal dependence, and that these seiches have their maximum amplitudes nearly concurrently. Tidal spring-neap dependence was not observed, leading them to reject barotropic or internal tides as a cause, instead, they concluded that rapid changes in air pressure and local winds associated with troughs and fronts are likely driving these seiches.

Persistent seiching around the islands of Mauritius and Rodrigues was observed by Lowry et al. (2008) with distinct fortnightly and seasonal amplitude variations along the east and southeast coasts, but no fortnightly or seasonal variations along the west coast. Breaker et al. (2008) noticed continuously present seiches in Monterey Bay, leading Breaker et al. (2010) to consider several possible forcing mechanisms (edge-waves, long period surface waves, sea breeze, internal waves, microseisms, and small-scale turbulence) and to question whether or not “the excitation is global in nature”. Subsequent analysis by Park and Sweet (2015) confirmed continuous oscillations in Monterey Bay over a 17.8 year record, and presented kinematic analysis discounting potential forcings of internal waves and microseisms while suggesting that a persistent mesoscale gyre situated outside the

Bay would be consistent with a jet-like forcing. However, jet-like currents are not a common feature along coastlines and could not be considered a global excitation of continuous oscillations.

Wijeratne et al. (2010) observed seiches with periods from 17 to 120 minutes persistent throughout the year at Trincomalee and Colombo, Sri Lanka, finding a strong fortnightly periodicity of seiche amplitude at Trincomalee on the east coast, but no fortnightly modulation on the west coast, rather a diurnal one attributed to weather forcing. The fortnightly modulations were strongest at neaps due to it taking a week for internal tides to travel from the Andaman Sea. It is notable that Wijeratne et al. (2010) attributed the 74 minute seiche period at Colombo and the 54 minute period at Galle to tidally-forced shelf-resonances.

Most recently, MacMahan (2015) analyzed 2 years of data (2011–2012) in Monterey Bay and Oil Platform Harvest, 270 km south of Monterey, concluding that low-frequency “oceanic white noise” within the seiche periods of 20 to 60 minutes might continuously force bay modes. The oceanic noise was hypothesized to consist of low-frequency, free, infragravity waves forced by short waves, and that this noise was of O(mm) in amplitude. So while the term “noise” applies in context of a low amplitude background signal, and the qualifier “white” expresses a uniform spatial and wide range of temporal distributions, the underlying processes are coherent low-frequency infragravity waves. Based on a linear system transfer function between Platform Harvest and Monterey Bay water levels, he concluded that the bay amplifies this noise by factors of 16–40 resulting in coherent seiching. It was also suggested that the highest-greatest amplification, a factor of 40, is associated with the 27.4 minute mode, however, as discussed below and in agreement with Lynch (1970), we find that this is not a bay-mode, but a tidally-forced shelf-mode, and find an amplification factor (Q) of 12.9. Further, as discussed below, we find low-frequency infragravity waves may not have sufficient energy to drive the observed oscillations, but are likely a contributor to observed seiche amplitude variability.

The foregoing suggests that coastal seiche are not to be considered solely a transitory phenomena, yet it seems that continuous seiche is not as widely known as its transitive cousin. For example, Bellotti et al. (2012) recognized the importance of shelf and bay-modes to tsunami amplification, but considered them to be independent processes, and the comprehensive review by Rabinovich (2009) falls short of continuous seiche recognition by noting that “in harbours and bays with high Q-factors, seiches are observed almost continuously.”

The focus of this paper is to present evidence in pursuit of the questions posed by Giese et al. (1993) and Breaker et al. (2010), namely, is there a continuous global excitation of seiche-seiches, and what is the source? We find that perpetual seiche-seiches are indeed present at six harbors examined around the Pacific, suggesting that there is a global-widely-occurring excitation. Spectral analysis of water levels identifies shelf-modes at each location as tidally-forced standing waves in agreement with Webb (1976) and Clarke and Battisti (1981), and identifies resonances down to harbor and pier scales.

Decomposition of shelf–resonance time series into intrinsic mode functions (IMFs) allows us to examine temporal characteristics of shelf–resonance in detail, leading to identification of fortnightly tidal signatures in the shelf–modes. An energy assessment of the available power from the shelf–modes as well as the power consumed to drive and sustain observed oscillations in Monterey Bay indicates that shelf–modes are indeed energetic enough to continually drive seiches. In terms of the query posed by Giese et al. (1993), we will suggest that their results are specific to locations where large internal waves are known to occur, and that ~~long-period~~ shelf–resonances driven by tides might constitute a more general excitation of continuous bay and harbor modes.

2 Resonance

Let us first recall some essentials of resonance as a general phenomenon. The simple harmonic oscillator consisting of a frictionless mass m coupled to a spring serves as a useful example. When the mass is displaced from equilibrium the spring imposes a restoring force $F = k_s x$ where k_s is the spring constant and x the displacement. Applying the equation of motion to this expression results in a 2nd order differential equation with solution $x(t) = A \cos(\omega_0 t + \phi)$ where A is the amplitude of the resulting oscillation, ω_0 the frequency of oscillation and ϕ a phase constant. The frequency of oscillation is $\omega_0 = \sqrt{k_s/m}$ and it is evident that this *natural* frequency is a property of the system governed by the stiffness of the spring and the mass, and is not influenced by the displacement or force which initiated the oscillation. When one adds damping to the system, or considers a different physical system such as a pendulum, vibrating string, electrical circuit, or in our case a semi–closed basin of water, the natural frequency remains a property of the system, not of the forcing.

A hallmark of resonance is that when a resonant system is driven at or near the resonant frequency, the resulting amplitude is maximized. This is conventionally expressed by an amplification factor of the form $(\omega - \omega_r)^{-2}$ where $\omega_r = \omega_0 \sqrt{1 - 2\zeta^2}$ is the *resonant* frequency of the system, ω the frequency of the driving force and ζ the damping ratio which in the case of a damped mass–spring system is $\zeta = \frac{c}{2\sqrt{k_s m}}$ where c is an integration constant. Again, the natural or resonant frequency is a property of the system, and as demonstrated by our simple harmonic oscillator it is not required to drive a resonant system at or near the resonant frequency in order to initiate and sustain small–amplitude resonant oscillations.

This is consistent with our hypothesis that tidally–forced shelf–resonances can initiate and sustain small–amplitude bay or harbor seiche even though the tidal forcing is not frequency–matched to the seiche. An energy analysis below based on the observed resonant properties of water levels in Monterey Bay illustrates this. Since the terms *mode* and *resonance* represent the same underlying property of the system, specifically a characteristic eigenmode, we use these two terms interchangeably. We apply modifiers of shelf, bay or harbor to qualify the spatial scale and domain over which the mode operates.

3 Locations and Data

We examine tide gauge water levels from six bay/harbors shown in figure 1 with the tide gauge location denoted with a star. Monterey Bay is located on the central California coast of North America, Hawke and Poverty Bays on the eastern coast of the North Island of New Zealand, Hilo resides on the western shore of the island of Hawaii, Kahului on the northern shore of the island of Maui, and Honolulu on the southern shore of the island of Oahu. Tides at these stations can be characterized as mixed-semidiurnal with mean tidal ranges listed in table 1.

Three of the bays (Monterey, Hawke and Poverty) can be characterized as semi-elliptical open bays with length-to-width ratios of 1.9, 2.0 and 1.4 respectively. We therefore anticipate a degree of similarity between their resonance structures. Bays at Hilo and Kahului are also similar with a triangular or notched coastline, while Honolulu is an inland harbor of Mamala Bay.

Data for Hawke and Poverty bays at the Napier (NAPT) and Gisborne (GIST) tide gauges respectively are recorded at a sample interval of $T_s = 1$ min, and are publicly available from Land Information of New Zealand (LINZ) at http://apps.linz.govt.nz/ftp/sea_level_data/. Data for Monterey and Hilo at a sample interval of 6 min are available from the National Oceanic and Atmospheric Administration (NOAA) tide gauges at <http://tidesandcurrents.noaa.gov/stations.html?type=Water+Levels>. In addition to this publicly available data, we also analyze water level data from independent wave studies at Honolulu, Hilo and Monterey sampled at 1 s intervals. At Honolulu data was collected by Seabird 26+ wave and water level recorders using Paroscientific Digi-quartz pressure sensors at two locations, one collocated with the NOAA tide gauge inside the ~~harbour~~harbor, and the other at $157.865^\circ\text{W } 21.288^\circ\text{N}$ outside Honolulu harbor. At Monterey and Hilo data was recorded at 1 s intervals by WaterLog H-3611 microwave ranging sensors co-located with the NOAA tide gauges. Table 2 lists the approximate bay and harbor dimensions along with the periods of record and sampling intervals.

4 Continuous Modes

Continuous seiching throughout a 17.8 year period has been observed, but most studies are limited to periods less than a year, one or a few coastal locations, and a single geographic region. In figure 2 we present water level spectrograms at six locations around the Pacific basin where vertical bands are associated with seasonal or episodic wave energy, and horizontal bands indicate the presence of persistent oscillations. These oscillations appear to have essentially invariant amplitudes over an extended period of time suggesting that time varying processes such as weather or waves are not likely forcings. For example, inspection of the Kahului data at periods near 0.2 min (12 s) reveals time varying amplitudes from wind waves and swell, whereas the longer-period oscillations are essentially constant. We therefore have reason to suspect that there is a continuous forcing of bay and harbor oscillations.

A close examination of modes at Kahului with periods between 1 and 5 minutes does reveal a time-dependent *frequency modulation*. The 2 minute mode is a good example where a distinct sinusoidal oscillation in period is found throughout the record. This behavior is also observed at Monterey (Park and Sweet, 2015), Hilo and Honolulu where high-resolution (1 Hz) data was available, but is not shown in figure 2. These modulations are coherent with the tides, and are a manifestation of changing boundary conditions (water depth, exposed coastline and spatial resonance boundaries) as water levels change with the tide. It should be noted that these modes are not directly forced by tidal energy, rather by wind-waves, and that their frequency is changing due to changes in the system as the tide changes water levels.

5 Mode Identification

Spectrograms provide information regarding time dependence of energy, but are not well suited when detailed frequency resolution is desired. To identify resonances in the water level data we estimate power spectral densities with smoothed periodograms (Bloomfield, 1976) as shown in figure 3. The Monterey and Hilo estimates are composites of 6 min and 1 s data with periods longer than 12 min represented by spectra of the 6 min data. Horizontal arrows indicate the range of modes associated with their respective spatial domains as discussed below. Triangles mark the tidally-forced shelf-resonances, also discussed below.

To relate temporal modes with spatial scales we find solutions to the general dispersion relation $\omega^2 = gk \tanh(kd)$ where ω is the mode frequency obtained from power spectra in figure 3, k the wavenumber, and d the water depth which is a representative value from nautical charts over the horizontal dimensions of the respective mode wavelength in the bay or harbor. This provides estimates of the modal wavelength $\lambda = 2\pi/k$, which we list as $\lambda/2$ or $\lambda/4$ in table 3 for all prominent modes. $\lambda/2$ corresponds to spatial modes between two fixed boundaries, for example between two opposing coasts of a bay as found in the longitudinal direction of the semi-elliptical bays, while $\lambda/4$ corresponds to one fixed and one open (free) boundary condition as found in a transverse mode where one boundary is a coast and the other the open sea, as is the case for the tidally-forced shelf-resonances.

For example, the 55.9 minute mode at Monterey and the 160–170 minute modes at Hawke correspond to longitudinal modes between the ends of the bays and are therefore delineated as closed-boundary $\lambda/2$ modes. The majority of the open-boundary condition modes correspond to transverse bay and shelf modes, however there are exceptions such as the 1 minute mode at Monterey and the 32 second mode at Honolulu which are open-boundary waves supported by open basins near the tide gauges as evidenced on harbor maps. We cannot assure that all entries in table 3 are properly attributed as $\lambda/2$ or $\lambda/4$ modes, as we have not closely examined the physical boundary conditions of each mode.

5.1 Shelf Resonance

The period of a shallow water wave resonance supported by a fixed-free boundary condition is expressed in Merian's formula for an ideal open basin as $T = 4L/\sqrt{gd}$ where L is the shelf width corresponding to $\lambda/4$, and d the basin depth (Proudman, 1953). In addition to a shelf-mode standing wave based solely on geometric wave reinforcement, a shelf-resonance is dynamically supported when the shelf width is approximately equal to $g\alpha/(\omega^2 - f^2)$ where g is the gravitational acceleration, α the shelf slope, ω the frequency of oscillation and f the Coriolis parameter (Clarke and Battisti, 1981). Table 4 lists solutions for shelf-mode period (inverse of frequency) for each of the bays where the shelf slope is approximated as the depth of the shelf break divided by the shelf width, and where the basin depth is taken as one half the shelf break depth. Also listed are modal periods deemed to represent the shelf-resonances obtained from the power spectra in figure 3. The agreement is reasonable given the simplistic formulations and crude spatial representations, and when viewed from the perspective of the apparently time invariant modal energy evident in the spectrograms and with recognition of tidal energy as a driver of shelf-resonances, suggests that tidally-forced shelf-resonances are continually present.

5.2 Dynamic Similarities

Topological similarities between Monterey and Hawke bays are striking, each a semi-elliptical open bay with aspect ratios of 2.0 and 1.9 respectively, although a factor of 2 different in horizontal scale. One might expect that these similarities would lead to affine dynamical behavior in terms of modal structure, although not the specific modal resonance periods, and that indeed appears to be the case as seen in figure 3. Both bays exhibit highly tuned resonances evidenced by high quality factors (Q) in the bay modes. The shelf-resonances of both bays, 27.4 min at Monterey and 105.8 min at Hawke, indicated with the triangle symbol in each plot, are exceptional examples of this, while the longer period modes (56 min at Monterey and 165 min at Hawke) correspond to longitudinal bay oscillations. The semi-elliptical topology of these bays is such that boundaries of the longitudinal modes are not parallel as in an ideal rectangular basin, but are crudely represented as semi-circular boundaries. The range of spatial scales between these boundaries is reflected in the longitudinal spectral peaks with broad frequency spans at the base and evidence of a series of closely spaced modes corresponding to a range of wavelengths. This is contrasted to the shelf-modes where the resonances are remarkably narrow indicating the narrow-range of spatial scales reflected in the relatively uniform widths of the shelves at Hawke and Monterey Bays.

Poverty Bay is the other semi-elliptical open bay and it exhibits the same generic modal structure. Although here, the bay modes are shorter in period due to the significantly smaller size, and the shelf-mode is the longest period mode. It is also evident here that the shelf-mode is mixed with other modes as it does not have a high Q -factor as found at Monterey and Hawke, although part of

this difference could result from poorer trapping or more radiation or other energy loss associated with this mode.

Hilo and Kahului bays also share structural similarity, but lack the high degree of topological symmetry found in the semi-elliptical bays that support both longitudinal and transverse modes.

240 As is the case for the semi-elliptical bays, the power spectra of these two bays are conspicuously similar with the substantial difference being the precise frequencies of their associated modes. Here, shelf-modes appear to dominate the water level variance at periods less than one hour, but rather than a set of discrete, high-Q shelf-resonances as found at Monterey and Hawke, they are energetic over a broad range of frequencies and spatial scales. This suggests that the shelves here are not well
245 represented by a uniform width, but encompass a range of scales to the shelf break as evidenced in bathymetric data. In the following sections we examine specific resonance features at each of the bays.

5.3 Monterey

Monterey Bay ~~see~~seiche has been studied since at least the 1940's (Forston et al., 1949) with
250 a comprehensive review provided by Breaker et al. (2010). The primary bay modes at the Monterey tide gauge have periods of 55.9, 36.7, 27.4, 21.8, 18.4 and 16.5 min, where the 55.9 min mode represents the fundamental longitudinal mode, while the 36.7 min harmonic is attributed to the primary transverse mode. We identify the 27.4 min mode as a shelf-resonance, also recognized by Lynch (1970), and consider it to be a potential continuous forcing of ~~long period~~-water level oscillations
255 throughout the bay at periods longer than 10 minutes. The harbor modes (figure 3) have been associated with resonances between breakwaters, and are amplified by wave energy, whereas the bay modes are weakly-dependent on wave forcing (Park and Sweet, 2015).

5.4 Hawke

Hawke Bay is approximately 85 km long and 45 km wide with a rich set of modes at periods between
260 20 and 180 minutes. Modes at periods of 170.6, 167.1 and 160.1 min correspond to longitudinal oscillations, while the 105.8 min oscillation is identified as a shelf-resonance (table 3).

5.5 Hilo

At Hilo we are afforded full spectral frequency coverage and find that pier modes have periods below 20 seconds corresponding to spatial scales less than 100 m. These modes are excited by waves and
265 swell just as the harbor modes at Monterey. Harbor modes at periods of 3, 4 and 5.9 minutes correspond to standing waves within the breakwater and spatial scales of 1, 1.3, and 1.9 km respectively. The shelf offshore Hilo is not a uniform width, but transitions from less than 2 km just south of the bay to roughly 18 km along the northern edge with the spectra revealing a corresponding plateau at periods between 10 and 30 minutes with a rather broad shelf-resonance centered on a period of 30.9

270 minutes, qualitatively different from the high-Q shelf-resonances at Monterey and Hawke bays. This well known 30.9 minute mode at Hilo corresponds to a shelf-resonance on a shelf width of approximately 17 km.

5.6 Kahului

Oscillations at Kahului follow the same general structure as Hilo with wave and swell excited pier
275 modes at periods less than 20 seconds, and within-harbor pier-breakwater modes at periods of 51 and 63 seconds. The primary harbor mode has peak energy at 188 seconds (3.1 minutes) corresponding to a $\lambda/2$ spatial scale of 1.1 km which is the dominant lateral dimension of the harbor.

An interesting feature of the Kahului power spectra is a low energy notch between periods of 120 and 160 seconds. This lack of energy corresponds to a lack of standing wave reflective boundaries
280 at scales of $\lambda/2$ from 650 to 1000 m. Such low energy features are present in all spectra indicating spatial scales where standing waves are not supported. The dominant shelf-mode at Kahului has a period of 35.5 minutes, similar to that of Hilo.

5.7 Honolulu

At Honolulu we have the benefit of both short sample times ($T_s = 1$ s) and two gauge locations, one
285 inside the harbor and one on the reef outside the harbor. The offshore power spectrum is shown in red in figure 3 exemplifying an open ocean or coastal location dominated by wind waves and swell. The rejection of wind wave energy inside the harbor is impressive, revealing a set of pier-modes in the 8 to 20 second band supported by rectangular basins around the gauge. Modes with periods of 82 and 88 seconds correspond to waves with $\lambda/2$ of approximately 500 m, which is the fundamental
290 dimension of the basin.

While the harbor is quite efficient in rejection of wind waves and swell, amplification of the shelf-mode and other long period resonances is a striking manifestation of the “harbor paradox” as noted by Miles and Munk (1961). Indeed, power spectra of the other harbors in figure 3 might suggest that they may be even more efficient amplifiers.

295 5.8 Poverty

Poverty Bay is a small-scale version of Hawke and Monterey bays with a similar resonance structure, however, the bay is small enough that the lowest frequency mode is not a longitudinal mode within the bay, but is the shelf-resonance at a period of 79 minutes. The 57.3 minute mode is not explicitly a Poverty Bay mode, but is a longitudinal mode of the open bay between Table Cape to the south
300 and Gable End to the north, inside which Poverty Bay is inset. We also note that the 42.1 minute mode is a shelf edge-wave evident in both Hawke and Poverty bays as discussed below. The reader is referred to Bellotti et al. (2012) for a detailed numerical evaluation of Poverty Bay shelf and bay modes.

5.9 Hawke and Poverty

Hawke and Poverty bays are located approximately 35 km apart along the southeast coast of northern New Zealand. Concurrent 7 month records allow examination of cross-spectral statistics between the two locations, with power spectra presented in the upper panel of figure 4 and coherence in the lower panel plotted as the upper and lower 95% confidence intervals. Power spectra reveal that the two bays share shallow water tidal forcings at periods longer than 180 minutes, but are essentially independent in terms of major oscillation frequencies between 20 and 180 minutes. There are coincident spectral peaks near periods of 42 and 58 minutes, however the coherence of the 58 minute energy is low indicating that is likely independent between the two locations.

Coherence at the shallow water tidal periods (373, 288, 240, 199 min) is quite high and as expected has near zero phase shift (not shown). Shelf-modes with periods from 100 to 160 minutes also share coherence in the 0.5 range, which is sensible since they have quarter wavelengths that are as long, or longer than, the 35 km separation distance. The only other energy with coherence reliably above the 0.5 range is the 42 minute mode. This mode has a phase shift of -160° from Napier (Hawke Bay) to Gisborne (Poverty Bay) indicating a traveling wave moving from south to north along the coast, empirically validating the shelf edge-wave explanation inferred numerically by Bellotti et al. (2012).

6 Shelf-metamodes

Since tidally-forced shelf-modes are a plausible driver of seiches, we expect that tidal amplitude variance should be reflected in seiche amplitudes, a view consistent with the strong fortnightly modulation of seiche amplitude reported by Giese et al. (1990) and Wijeratne et al. (2010). To examine such a dependence figure 5 plots time series of shelf-mode amplitude at each station. Amplitudes are ~~determined by shelf-mode power-spectral values versus time at~~ estimates of the water level oscillation measured at the tide gauge at the frequency of the corresponding shelf-mode for each station. To examine the mean temporal behavior of shelf-modes we decompose the time series into intrinsic mode functions (IMFs) by empirical mode decomposition (EMD, Huang and Wu (2008)). The temporal low-pass response of the shelf-modes is shown in figure 5 by the thick lines which are a superposition of the lowest frequency intrinsic mode functions (IMFs). We term these IMFs of shelf-mode amplitudes as *metamodes*.

It is clear from figure 5 that shelf-modes are continually present at all stations, albeit with significant temporal variability. The metamodes reveal annual modulations in the long period records of Monterey and Hilo, and fortnightly cycles at Kahului and Honolulu.

To assess the relative contribution of individual shelf-mode IMFs (metamodes) to the total shelf-mode variance, we list the mean period in days (T) of each metamode Hilbert instantaneous frequency vector, and the estimated partial variance of each metamode IMF with respect to the total variance in table 5. The partial variance is determined by excluding, one-by-one, IMFs from the re-

constructed signal (sum of all IMF's) and comparing that result to the variance of the original signal.

340 Since IMFs are not required to be linear or mathematically orthogonal, the sum of partial variances is not restricted to a value of one, although in practice IMFs are nearly orthogonal and the sum of these partial variances is within several percent of unity.

We note that the fortnightly astronomical tidal constituents, the lunisolar synodic fortnightly (M_{sf}) and lunisolar fortnightly (M_f), have periods of 14.76 and 13.66 days respectively with IMFs closest
345 to these periods highlighted in table 5 and shown in figure 6, where we find at Monterey and Hilo that the fortnightly variance in shelf-mode amplitude is the dominant contribution, while at Kahului, Hawke and Poverty bays it is the second strongest metamode. At Honolulu and Kahului the bulk of the metamode variance is on sub-daily time scales, however, the fortnightly mode is the strongest of the modes at diurnal and longer scales, a relation that holds at all stations except Poverty Bay. We
350 also note in figure 6 evidence of a seasonal dependence in metamode amplitude, and will correlate these IMFs with their corresponding fortnightly tidal IMFs below.

The foregoing indicates that fortnightly metamodes are present at all six stations suggesting that tidal forcing of shelf-modes is a likely driver. To assess an assumed linear dependence between fortnightly tidal forcing and metamodes, we compute IMFs on the tidal water level data and cross-
355 correlate the resulting fortnightly tidal IMFs with the fortnightly metamodes. Correlations are computed over a sliding window of length 20 days with results shown in figure 7 where dashed red lines indicate the 95% significance threshold. An interesting feature is that while all stations exhibit near perfect correlation at times, they also episodically transition to near zero or statistically insignificant correlation. This suggests that the fortnightly modulation of tidally-driven shelf-resonances is also
360 influenced by other factors, of which internal tide variability has previously been noted by Giese et al. (1990) and Wijeratne et al. (2010).

While there is significant temporal variability in the fortnightly tidal metamode correlations, it appears that the majority of the time correlations are quite high and significant above the 95% level. The IMF mode numbers and mean correlations statistics are listed in table 6 where $T_R\%$ is the per-
365 centage of time that the 95% significance threshold is exceeded, \bar{R} is the mean correlation of values above the 95% significance threshold, and \bar{Lag} the mean lag value of 95% significant correlations. Overall, these data suggest that correlations significant above the 95% level are present 76–87% of the time, and from a linear model perspective that fortnightly tidal oscillations account for 35–50% of the metamode variance.

370 7 Mode Energy

Knowledge of a mode's temporospatial characteristics allows estimation of the total energy sustained by the mode. Figure 3 indicates that Hilo, Kahului, Honolulu and Poverty bays are dominated by energy of the shelf-mode, while at Hawke and Monterey bays the shelf-mode is the second largest

amplitude. We are therefore motivated to investigate modal energy in Monterey Bay to test our
375 hypothesis that the ~~shelf-mode~~ tidally-forced shelf-resonance is a potential driver of bay oscillations
from a kinematic perspective.

We estimate the total potential energy to support a mode by assuming a raised-cosine profile of
amplitude h either orthogonal to the shore for the transverse ~~and shelf-modes~~ modes, or parallel for
the longitudinal modes. Multiplying this profile area (A_M) by the alongshore extent of the mode
380 (L_A) gives the volume of water displaced: $V_M = L_A A_M$ where we have neglected the influence of
shoaling on the transverse modes as the wavelength is much longer than the shelf width. (This as-
sumption is supported by the agreement between the shelf-mode spatial scales based on the observed
shallow-water frequencies in table 4.) The energy to move this volume is equivalent to the work per-
formed to change the potential energy of the mass in the gravitational field $E_M = \rho V_M h_M g$, at an
385 average power output of $P_{out} = E_M/T$ where T is the modal period. This leading-order value does
not incorporate dissipation and momentum, terms that we ignore in subsequent energy estimates.

The ratio of energy stored in the mode resonance to energy supplied driving the resonance is
the Q factor. If Q is large (the resonance signal-to-noise ratio is high, as is the case for the shelf-
mode at Monterey), it may be estimated from the power spectrum: $Q = f_M/\Delta f$, where f_M is the
390 mode resonant frequency and Δf the -3dB (half power) bandwidth of the mode. This allows one to
estimate the power required to drive the mode $P_{in} = E_M/(QT) = P_{out}/Q$.

Modal length scales (λ) are taken from table 3, amplitudes (h) are from bandpass filtering the
17.8 year water level record at the NOAA tide gauge, and Q from 1 Hz water level power spectra
(95% CI 2.6 dB) modal means over 120 hour windows over 63 days (Park and Sweet, 2015). The
395 alongshore extent of the modes, L_A , are estimated from a regional ocean modeling system (ROMS)
implementation in Monterey Bay (Shchepetkin and McWilliams, 2005) as reported in Breaker et al.
(2010).

Results of these estimates are shown in table 7 where we find seiche amplitudes averaged over
the 17.8 year period of 0.9, 1.4 and 1.6 cm for the 55.9, 27.4 and 36.7 minute modes respectively,
400 although amplitudes of 4 cm in the 27.4 minute mode are common during seasonal maximums. The
27.4 min shelf-mode is estimated to produce a total power of 998 kW, which is more than sufficient
to supply the required input power of both the primary longitudinal (55.9 min, $P_{in} = 23$ kW) and
transverse (36.7 min, $P_{in} = 169$ kW) bay modes. So even though the primary longitudinal bay-modes
in both Monterey and Hawke bays have larger amplitudes than the respective shelf-modes, the shelf-
405 mode in Monterey supplies a continuous source of energy capable of sustaining the fundamental
bay-modes, which by virtue of their resonant amplification can exceed the amplitude of the shelf-
resonance itself. Regarding the O(mm) low-frequency infragravity waves suggested by MacMahan
(2015), we note that a 27.4 minute mode with an amplitude of 3 mm would produce an estimated
 P_{out} of 41.6 kW (not shown in table 7), which would be insufficient to drive the observed 27.4 min
410 mode as it requires a power of $P_{in} = 77$ kW.

8 Discussion

Resonant modes are a fundamental physical characteristic of bounded physical systems expressed in bodies of water as seiches. As such, they can be excited to large amplitudes by transitory phenomena such as weather and tsunamis, and since large amplitude ~~seiche~~seiches are easily observable ~~seiche~~seiches are often viewed as transitory given that they dissipate after cessation of the driving force. Moving from transient to persistent behavior, seasonal weather patterns are known to sustain nearly continuous ~~seiche~~seiching for extended periods (Woodworth et al., 2005; Wijeratne et al., 2010), as are internal waves (Giese et al., 1990). On the other hand, small amplitude, temporally continuous ~~seiche~~seiches were recognized by Cartwright and Young (1974) and Golmen et al. (1994), with Giese et al. (1993) and Breaker et al. (2010) posing questions as to the possibility of global excitations. Motivated by these questions we have examined tide gauge water levels around the Pacific basin looking for continuous seiching and forcings.

In the process of analyzing the resonant structure of these bays and harbors, we have quantified resonant periods and estimated spatial scales corresponding to each mode (table 3). In some cases, we have identified the physical attributes of a bay or harbor associated with specific temporospatial resonances. In a more general sense, we have also illustrated broad dynamical similarities between bays with affine topologies, such as the clearly defined modes of the semi-elliptical bays when compared to the less structured, shelf dominated bays such as Hilo and Kahului. This analysis also provides empirical verification of the numerically inferred edge-wave by Bellotti et al. (2012) near a period of 42 min along Hawke and Poverty bays.

Simple geometric and dynamical estimates of tidally-forced shelf-modes are consistent with modes observed in the power spectra at all stations, and their continual presence in water level spectrograms and mode amplitude time series indicates that tidally-forced shelf-modes are continuously present at each location. Decomposition of shelf-mode amplitude time series identifies metamodes reflecting dynamic behavior of the shelf-modes, and we find that fortnightly metamodes are the dominant mode at periods longer than diurnal. Assuming that these fortnightly modulations are of tidal origin, cross-correlation of fortnightly IMFs of tidal data with the fortnightly metamodes leads to the conclusion that within the bounds of a linear system model from one-third to one-half of the fortnightly metamode variance is coherent with tidal forcing. We therefore suspect that tidally-forced shelf-modes are a continuous energy source in harbors and bays adjacent to continental or island shelves.

A natural question is: does the proposed source contain sufficient energy to sustain the observed resonant oscillation? Power estimates of the most energetic modes at Monterey suggest that the shelf-mode is fully capable as a primary driver of continuous seiche, while the low-frequency infragravity waves suggested by MacMahan (2015) may not have sufficient energy.

9 Conclusions

Examination of 6 coastal locations around the Pacific with diverse shelf conditions finds that tidally-forced shelf-resonances are continually present. An energy assessment of the shelf-mode and primary seiche in Monterey Bay indicates that the shelf-resonance is fully capable of supplying the power input required to drive the primary bay oscillations even though the grave mode produces more output power than the shelf-mode, a consequence of the resonance structure of the Bay. Hawke Bay is dynamically similar to Monterey and we suspect that a similar relation holds there, while at the other locations the shelf-mode is the dominant energy source. Our conclusion is that tidally-forced shelf-modes constitute a global candidate for continuous seiche excitation, a view consistent with that of Lynch (1970), Golmen et al. (1994) and Wijeratne et al. (2010) who identified tidally-forced shelf-resonances as specific seiche modes. In locations where tidally-forced shelf-resonance is a primary seiche generator, we suspect that internal waves and weather, which clearly can be a primary forcing in their own right, serve to modulate seiche amplitudes.

Specific to Monterey Bay, these results offer a simpler explanation for continuous seiche generation than the mesoscale gyre hypothesis proposed by Park and Sweet (2015) which lacked a physical mechanism to transfer energy into the Bay, and is more energetically reasonable than the infragravity waves suggested by MacMahan (2015).

In the course of attributing tidal forcing as the driver of the observed shelf-resonances, we introduced the idea of metamodes, dynamical modes of shelf-mode amplitude determined by empirical mode decomposition. The metamodes exhibited fortnightly modulation, and it is likely that examination of other metamode components may be useful towards understanding the dynamic behavior of modal structure in coastal environments.

However, it is also clear that we do not understand the cyclic nature of fortnightly tidal and meta-mode correlation. One possibility is that there is a time varying phase-lag between the two such that destructive superposition episodically creates nulls. A linear spectral analysis might use a coherency statistic to identify this, but such an option is not available for IMFs with variable instantaneous frequencies. It is evident that internal tides play a role, and it may be that episodic changes in stratification as noted by Giese et al. (1990) lead to modulation of the metamodes and contribute to the observed decorrelation, and it is deemed likely that the free, long-frequency infragravity waves suggested by MacMahan (2015) also contribute.

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Table 1. Tidal ranges at the six tide gauges. GT is the great diurnal range (difference between mean higher high water and mean lower low water) and MN the mean range of tide (difference between mean high water and mean low water).

Location	GT (m)	MN (m)
Monterey	1.63	1.08
Hawke	1.78	1.06
Hilo	0.73	0.53
Kahului	0.69	0.48
Honolulu	0.59	0.39
Poverty	1.58	1.06

Table 2. Approximate shelf widths and dimensions of Bays and Harbors, data period of record and sampling interval T_s . Note that data from Hilo and Monterey include both long-period data recorded at $T_s = 6$ min and short-period data recorded at $T_s = 1$ s.

Location	Harbor (m)	Bay (km)	Shelf (km)	Period of Record	T_s
Monterey Bay and Harbor	600 x 500	40 x 20	15	August 25 1996 – June 23 2014	6 min
				September 14 2013 – November 29 2013	1 s
Hawke Bay, Napier Harbor	650 x 360	85 x 45	60	July 18 2012 – August 9 2013	1 min
Hilo Bay and Harbor	1950 x 1000	13 x 8	17	August 7 1994 – February 15 2010	6 min
				February 18 2014 – March 4 2014	1 s
Kahului Bay and Harbor	1100 x 950	23 x 11	20	February 14 2013 – June 4 2013	1 s
Mamala Bay, Honolulu Harbor	1000 x 500	19 x 5	15	June 30 2012 – September 27 2012	1 s
Poverty Bay, Gisborne Harbor	500 x 300	10 x 7	45	April 19 2009 – August 11 2010	1 min

Table 3. Temporospatial scales according to the dispersion relation $\omega^2 = gk \tanh(kd)$ where ω is frequency, k the wavenumber, d the water depth corresponding to horizontal dimensions of the respective mode, λ the wavelength and period is $2\pi/\omega$. Periods are in min and lengths in km unless otherwise noted. Periods are obtained from the peak modal energy represented in the power spectra shown in figure 3. Depths for each bay are taken as representative values from nautical charts, depths for shelf–resonances are assumed to be 150 m, one half a nominal shelf–break depth of 300 m. Spatial scales are listed as $\lambda/2$ for modes assumed to be fixed–fixed boundary standing waves, and $\lambda/4$ for fixed–open boundaries.

Monterey				Hawke				Hilo				Kahului				Honolulu				Poverty			
Period	Depth	$\lambda/2$	$\lambda/4$	Period	Depth	$\lambda/2$	$\lambda/4$	Period	Depth	$\lambda/2$	$\lambda/4$	Period	Depth	$\lambda/2$	$\lambda/4$	Period	Depth	$\lambda/2$	$\lambda/4$	Period	Depth	$\lambda/2$	$\lambda/4$
55.9	60	40.7		170.6	30	87.8		30.9	150		17.9	35.5	150		20.4	45.5	150		26.2	86	150		49.5
36.7	150		21.1	167.1	30	86		19.5	150		11.1	25.8	150		14.8	27.1	150		15.6	79	150		45.5
27.4	150		15.8	160.1	30	82.4		14.6	150		8.4	22.3	150		12.8	20.9	150		12	57.3	15		10.4
21.8	60	15.9		105.8	150		60.9	12.7	150		7.3	18.1	37	10.2		11.2	150		6.4	50	15		9.1
18.4	60	13.4		78.4	30	40.3		5.9	12	1.9		15.8	37	9		4.3	14	1.5		42.1	15		7.6
16.5	60	12		65.2	30	33.6		4	12	1.3		10.2	37	5.8		3.9	14	1.4		28	15	10.2	
10.1	60	7.4		56	30	28.8		3	12	995 m		5.1	13	1.7		2.9	14	1		23.2	15	8.4	
9	60	6.6		49.3	30	25.4		1.7	12	566 m		3.1	13	1.1		2.2	14	768 m		19.6	15	7.1	
4.2	60	3.1		47.1	30	24.2		1.3	12	423 m		1.9	13	645 m		1.7	14	587 m		15.7	15	5.7	
1.87	8	480 m		40.5	30	20.8		39 s	12	209 m		1.5	13	503 m		1.5	14	513 m		14.4	15	5.2	
1	8		128 m	36.7	30	18.9		29 s	12	155 m		1.3	13	437 m		1.4	14	475 m		11.8	10	3.5	
41 s	8	175 m		35.1	30	18.1		23 s	12	125 m		1.04	13	349 m		1.2	14	427 m		10.2	10	3	
31 s	8	132 m		33.1	30	17		17 s	12	90 m		51 s	13	287 m		45 s	14	261 m		5.2	10	1.5	
16 s	8	67 m		29.6	30	15.2		15 s	12	78 m		38 s	13	212 m		32 s	14		91 m				
12 s	8	50 m		27.8	30	14.3		14 s	12	74 m		20 s	13	110 m		29 s	14	165 m					
				25.1	30	12.9		13 s	12	68 m		16 s	13	84 m		26 s	14	150 m					
				22.9	30	11.8		12 s	12	61 m		11 s	13	56 m		16 s	14	90 m					
				21.3	30	11		11 s	12	53 m		9 s	13	48 m		14 s	14	77 m					
				19.7	30	10.1		9 s	12	44 m						12 s	14	63 m					
				18.4	30	9.5		7 s	12	34 m						10 s	14	51 m					
				17.2	30	8.9		6 s	12	26 m						8 s	14	40 m					
				14	30	7.2																	
				10.6	30	5.5																	
				7.4	30	3.8																	
				5.4	30	2.8																	
				3.2	20	1.3																	
				2.2	20	0.9																	

Table 4. Estimates of shelf–resonance periods. T_R is a solution to $L = g\alpha/(\omega^2 - f^2)$ where L is the shelf width, g the gravitational acceleration, α the shelf slope, ω the frequency of oscillation and f the Coriolis parameter. The shelf slope is estimated as break depth / width where we assume a break depth of 300 m. T_M is from Merian’s formula $T_M = 4L/\sqrt{gd}$ for an open basin where d is the basin depth which we assume to be one–half the shelf break depth. T_{PSD} are values from the power spectral density estimates from shelf–mode frequencies marked with triangles in figure 3.

Location	Latitude (deg)	Width (km)	T_R (min)	T_M (min)	T_{PSD} (min)
Monterey	36.6	15	28.9	26.1	27.4
Hawke	39.5	60	115.2	104.3	105.8
Hilo	19.7	17	32.8	29.5	30.9
Kahului	20.9	20	38.6	34.8	35.5
Mamala	21.3	15	29.0	26.1	27.1
Poverty	38.7	45	86.6	78.2	79.0

Table 5. Mean period in days (T) of Hilbert instantaneous frequencies and percent variance of shelf–resonance power spectral density IMFs (metamodes). Modes with mean periods close to fortnightly tidal constituents with periods of 14.76 (M_{sf}) and 13.66 days (M_f) are highlighted.

IMF	Monterey		Hawke		Hilo		Kahului		Honolulu		Poverty	
	T	% Var	T	% Var	T	% Var	T	% Var	T	% Var	T	% Var
1	14.70	33.5	3.34	2.0	14.47	35.2	0.30	42.6	0.29	40.2	2.42	6.2
2	28.31	22.2	6.27	5.1	29.99	22.3	0.60	17.7	0.51	19.5	5.52	19.6
3	55.48	15.1	9.02	5.9	60.83	14.6	1.26	14.2	0.98	18.5	8.46	13.7
4	105.06	8.1	13.48	13.9	137.37	18.2	2.41	5.3	1.77	4.7	14.21	17.9
5	219.06	10.0	17.01	4.4	273.70	11.4	4.65	2.3	3.48	6.2	22.60	10.9
6	387.16	8.4	24.48	7.4	476.62	0.8	14.58	18.1	6.07	3.6	34.31	8.4
7	726.21	1.8	35.73	6.8	938.04	0.5	28.41	4.5	13.86	7.0	53.76	11.4
8	1662.14	1.3	45.45	4.8	1790.81	0.2	56.98	2.2	22.55	0.1	68.84	7.1
9	2244.04	0.2	65.67	9.7	4009.13	1.4			46.91	0.8	111.13	4.3
10			162.08	35.9							164.22	5.6

Table 6. Cross-correlation of fortnightly tide and shelf-resonance metamode IMFs. $T_R\%$ is the percentage of time that the 95% confidence level is exceeded, \bar{R} the mean value of 95% significant correlation over the record, and \bar{Lag} the mean lag value of 95% significant correlation over the record.

Location	Tide	Shelf	$T_R\%$	\bar{R}	\bar{R}^2	\bar{Lag} days
	IMF	IMF	>95%	>95%	>95%	>95%
Monterey	5	1	87	0.67	0.45	-0.35
Hawke	7	4	84	0.71	0.50	-0.26
Hilo	5	1	82	0.71	0.50	0.53
Kahului	7	6	86	0.59	0.35	0.01
Honolulu	7	7	76	0.65	0.42	0.64
Poverty	6	4	82	0.69	0.48	-0.17

Table 7. Estimates of total energy and power generated by resonances in Monterey Bay. Modal amplitudes (h) are mean values from bandpass filtering the 17.8 year record of water levels at the NOAA tide gauge. T is the mode period, W_{FIR} is the filter bandpass, $\lambda/2$ the mode half wavelength, L_A the alongshore extent of the mode in the bay, V the volume of water displaced, E_M the potential energy, Q the mode amplification, $P_{in} = E_M/(QT)$ the input driving power of the mode, and $P_{out} = E_M/T$ the modal power.

T (min)	W_{FIR} (min)	h (cm)	$\lambda/2$ (km)	L_A (km)	V (Mm^3)	E_M (GJ)	Q	P_{in} (kW)	P_{out} (kW)
27.4	25–30	1.4	31.6	38	33.30	1.64	12.9	77.6	998.1
36.7	35–39	1.6	42.2	40	52.55	2.91	7.8	168.6	1319.7
55.9	53–59	0.9	40.7	18	13.27	0.43	5.6	22.7	127.3

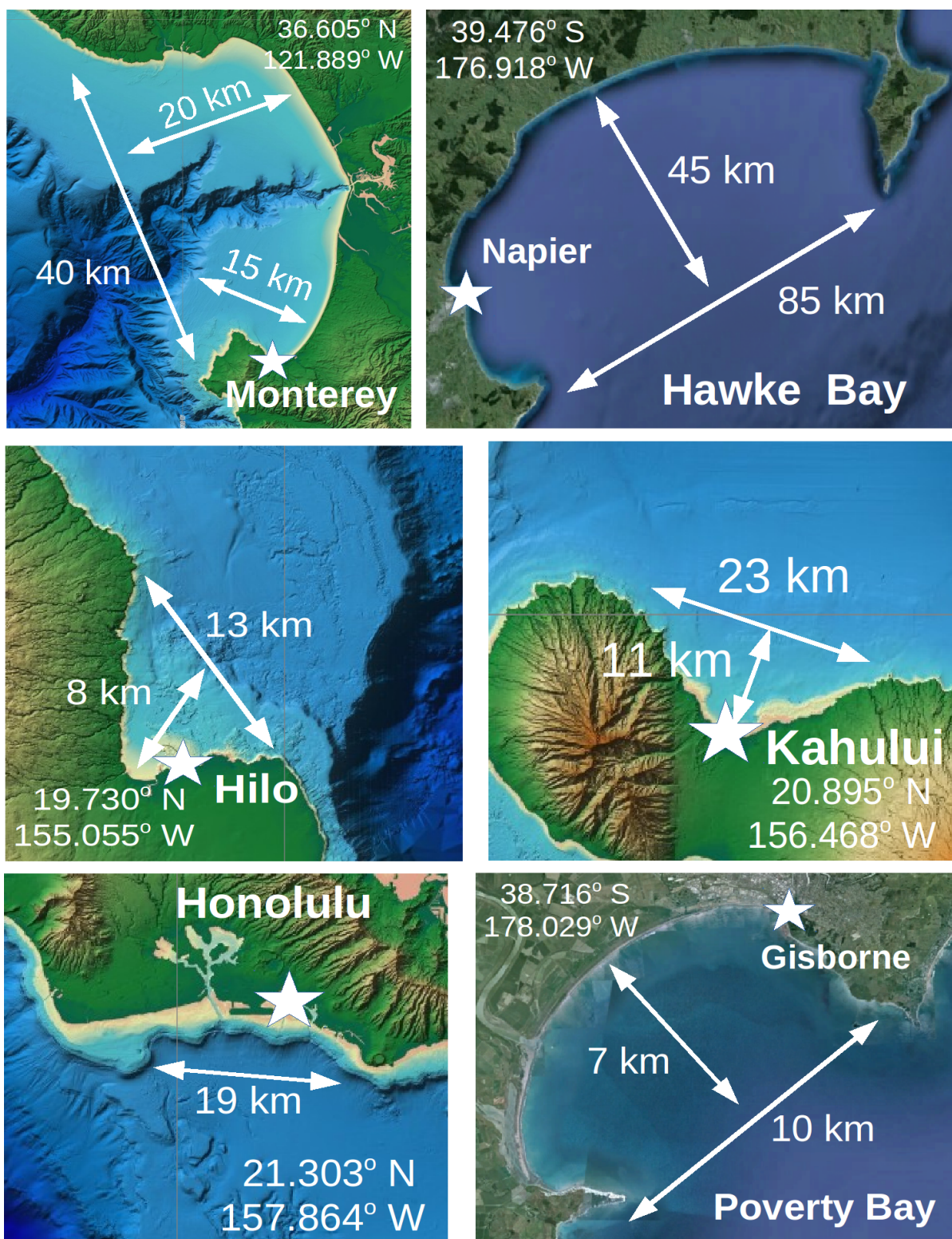


Figure 1. Location and approximate dimensions of bays. Tide gauge locations are marked with a star and denoted by latitude and longitude.

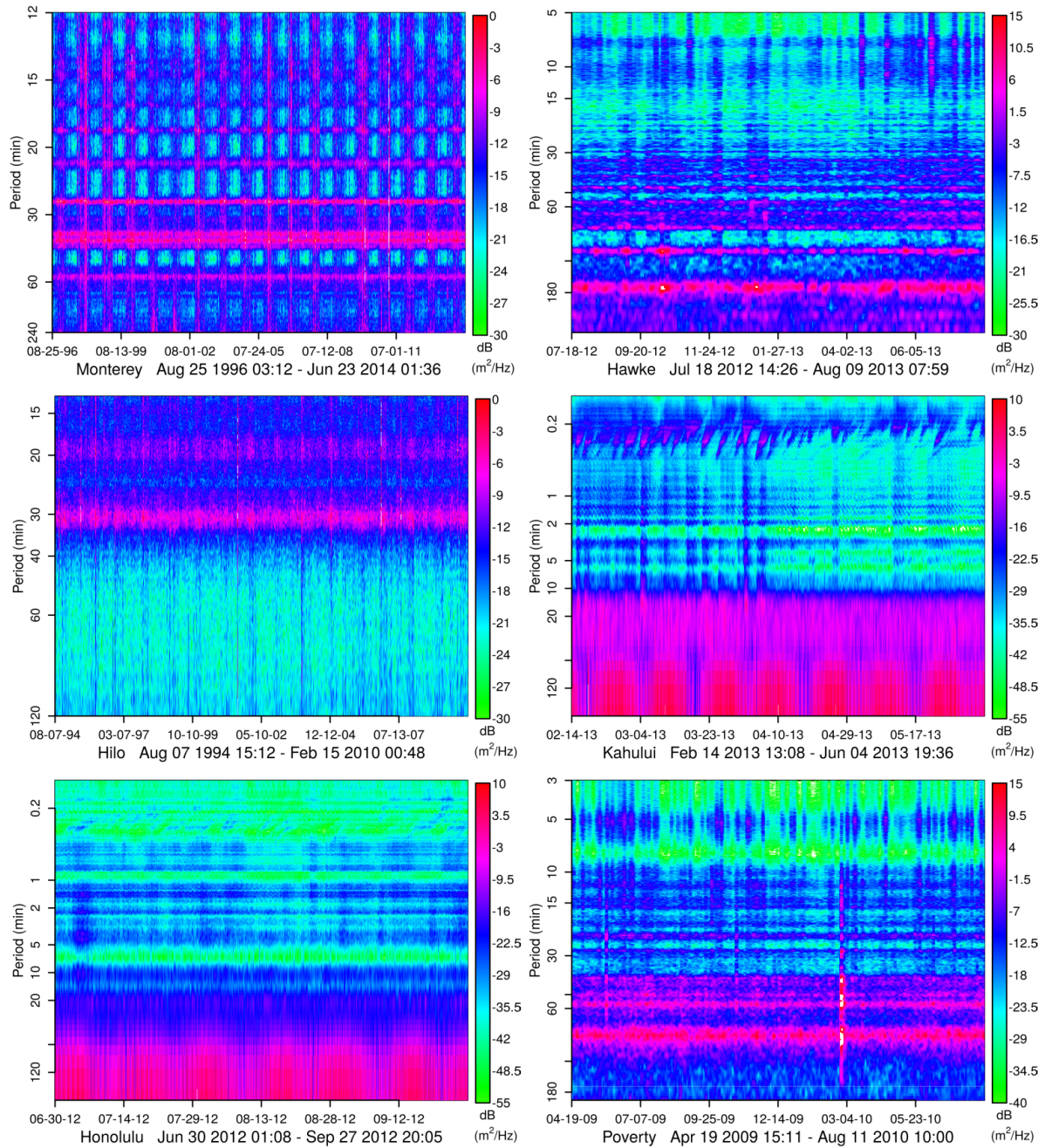


Figure 2. Spectrograms of water level data at each tide gauge. Horizontal bands indicate continuous oscillations, vertical bands are associated with periods of increased wave energy.

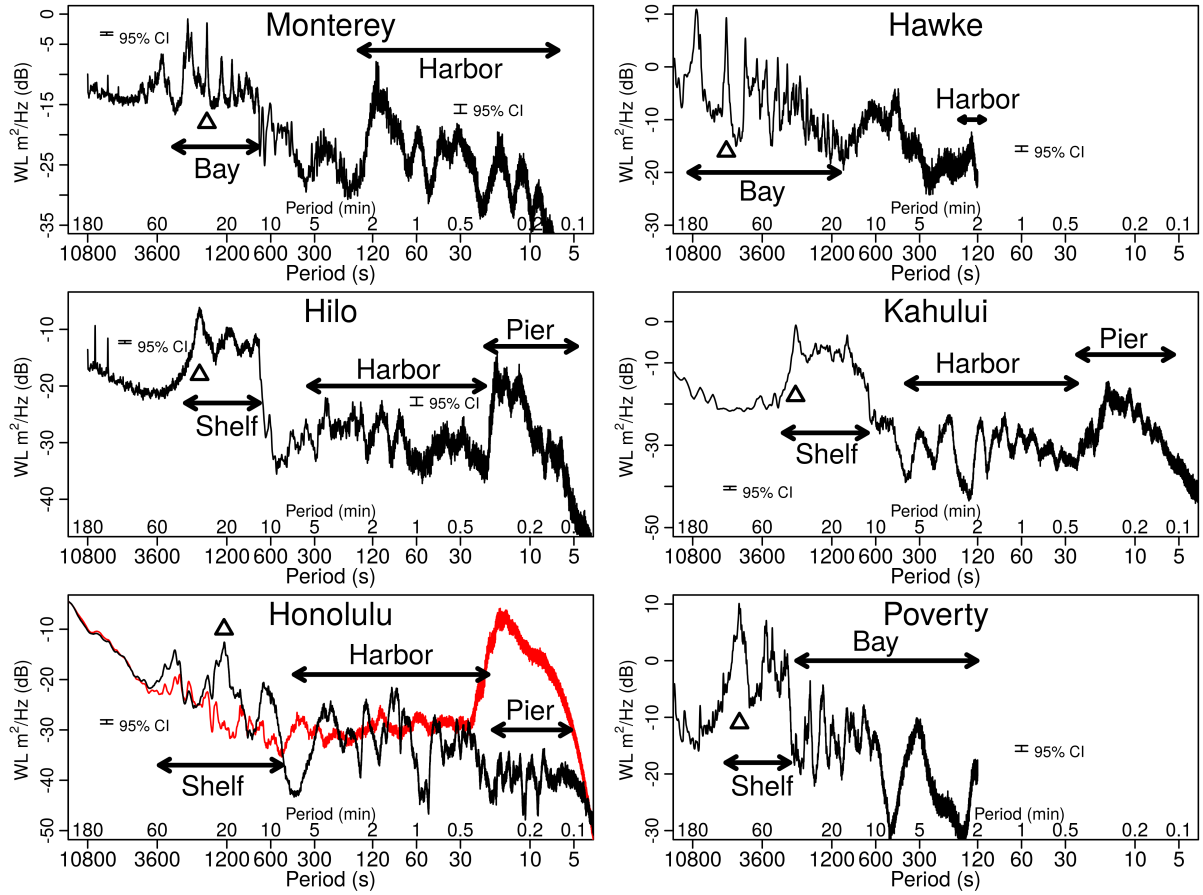


Figure 3. Power spectral density (PSD) estimates of water level (WL) at each tide gauge. Horizontal arrows indicate the frequency span of resonant modes associated with spatial scales. Triangles mark the tidally-forced shelf-resonance. The red curve at Honolulu plots data from outside the harbor.

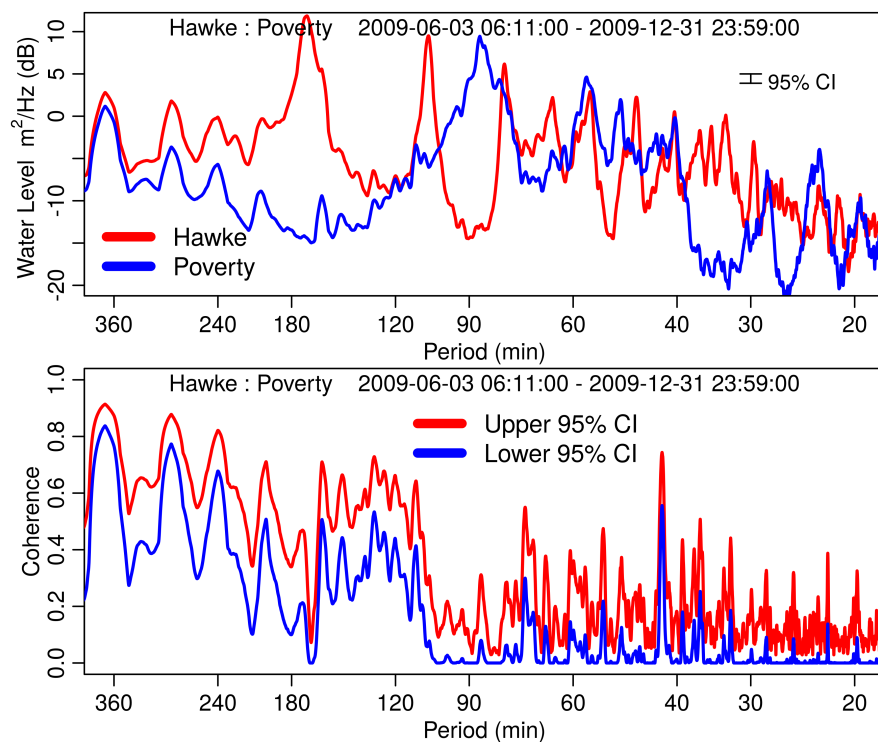


Figure 4. Power spectral density (top) of concurrent water levels at Napier in Hawke Bay, and Gisborne in Poverty Bay. Bottom: coherence of the power spectra shown as the upper and lower 95 percent confidence interval values.

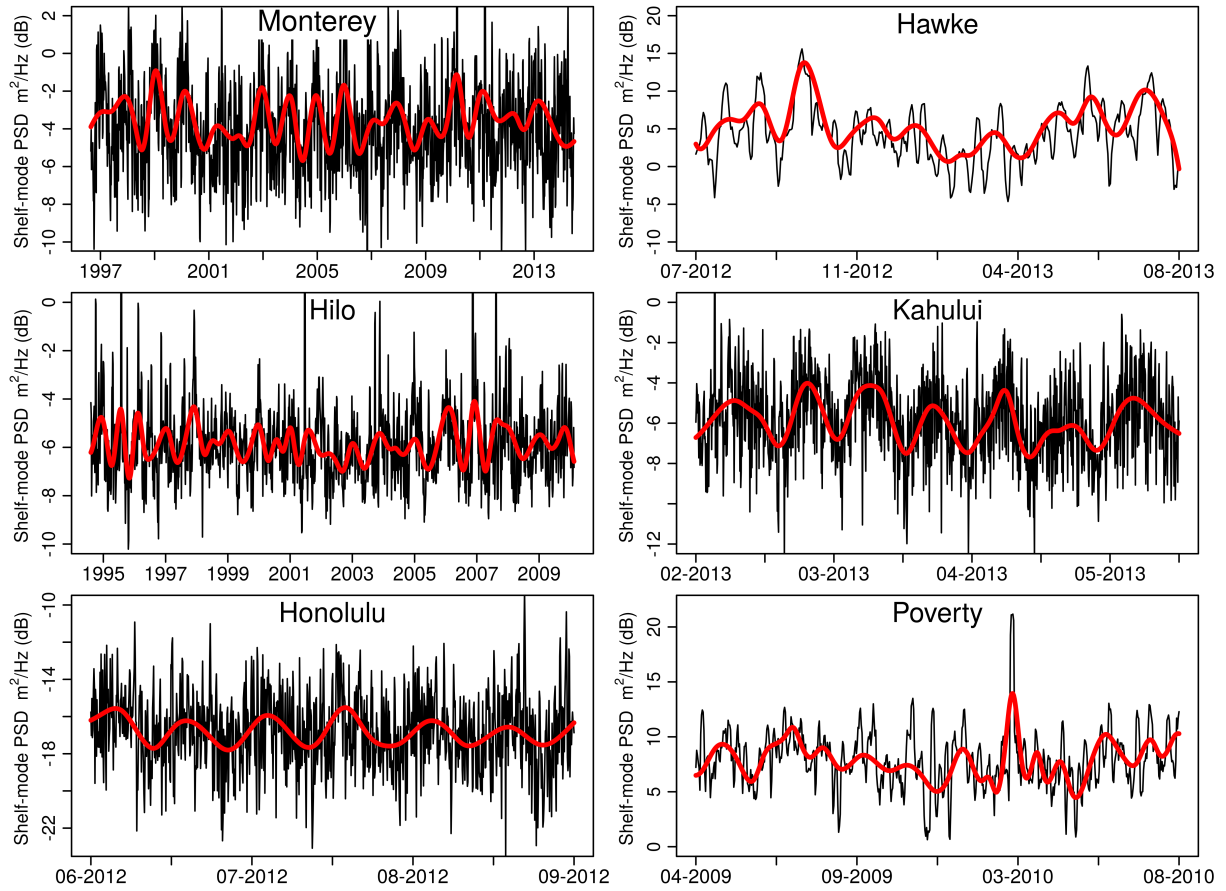


Figure 5. Shelf–resonance power spectral density (PSD) amplitudes (black) with low–frequency IMFs (meta–modes) in red. The large amplitude in Poverty Bay is a result of the February 27, 2010 Chile 8.8 M earthquake and tsunami.

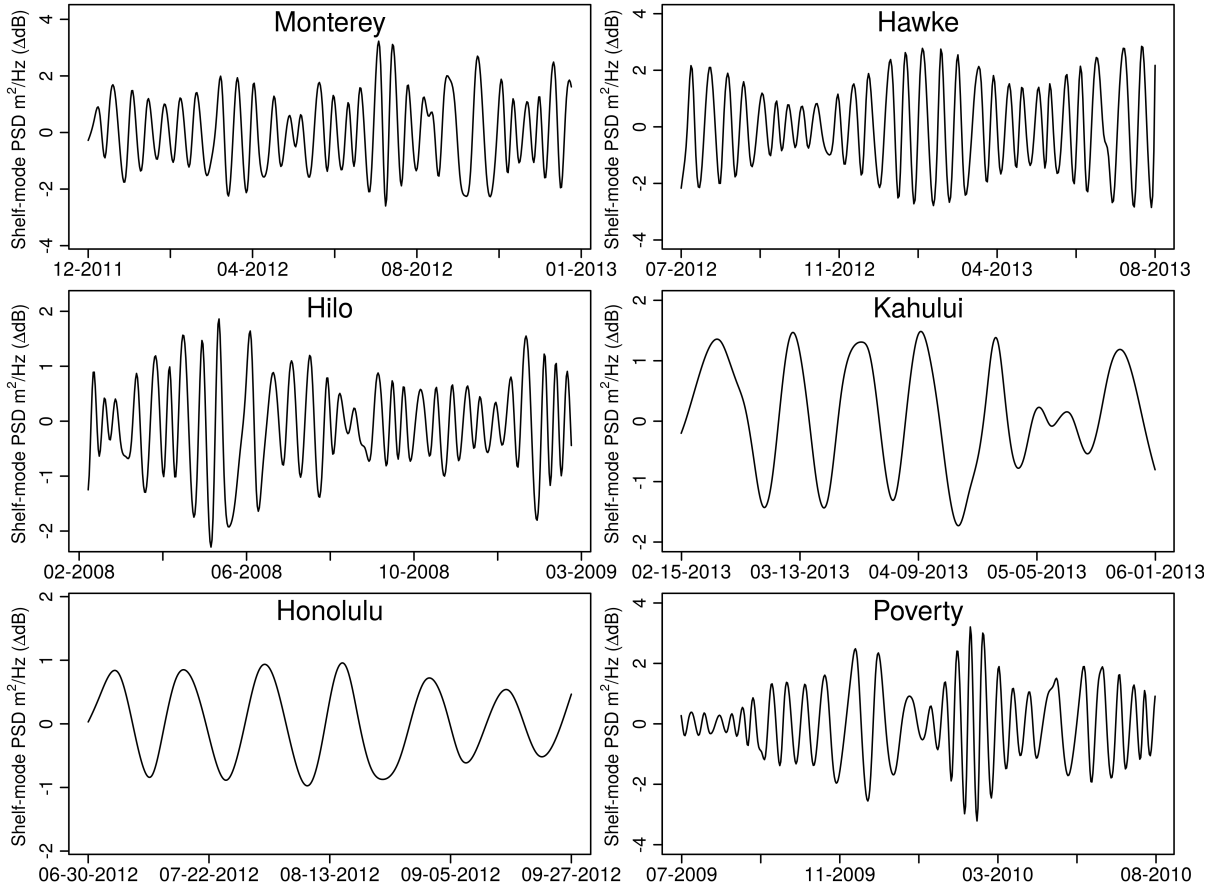


Figure 6. Intrinsic mode functions (IMF) of shelf-mode amplitude variance (metamodes) with mean Hilbert instantaneous frequencies corresponding to fortnightly periods (highlighted in table 5). Amplitudes are with respect to the mean values shown in figure 5. Records at Honolulu and Kahului are limited to 3 months, while the other stations show excerpts of approximately 13 months.

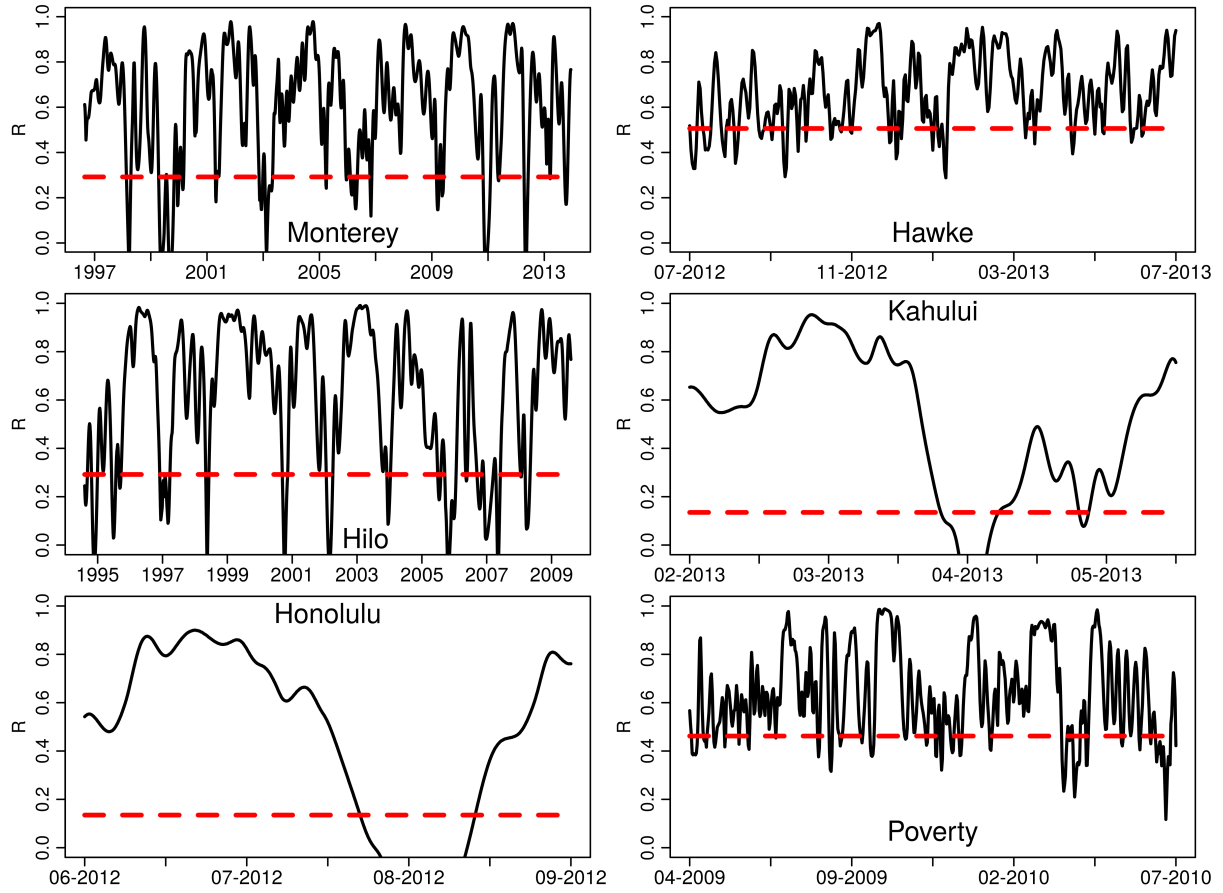


Figure 7. Correlation coefficients between tide and shelf–resonance metamode IMFs with fortnightly periods. The dashed red lines indicate the 95% confidence levels.