



## Abstract

The three-dimensional flow, temperature and salinity fields of the North Atlantic including the Arctic Ocean covering the time period 1992 to 2006 are simulated with the numerical ocean model CODE. The model reveals several new insights and previously unknown structures which help us to clarify open questions on the regional oceanography of Icelandic waters. These relate to the structure and geographical distribution of the coastal current, the primary forcing of the North Icelandic Irminger Current (NIIC), the path of the Atlantic Water south-east of Iceland and the structure of the North Icelandic Jet (NIJ). The model's adaptively refined computational mesh has a maximum resolution of 1 km horizontal and 2.5 m vertical in Icelandic waters. CTD profiles from this region and the river discharge of 46 Icelandic watersheds, computed by the hydrological model WaSiM, are assimilated into the simulation. The model realistically reproduces the established elements of the circulation around Iceland. However, analysis of the simulated mean flow field also provides further insights. It suggests a distinct freshwater-induced coastal current that only exists along the south-west and west coasts which is accompanied by a counter-directed undercurrent. The simulated transport of Atlantic Water over the Icelandic shelf takes place in a symmetrical system of two currents, with the established NIIC over the north-western and northern shelf, and a current over the southern and south-eastern shelf herein called the *South Icelandic Current* (SIC). Both currents are driven by topographically induced distortions of the Arctic Front's barotropic pressure field. The SIC is simulated to be an upstream precursor of the Faroe Current (FC). The recently discovered North Icelandic Jet (NIJ) also features in the model predictions and is found to be forced by the baroclinic pressure field of the Arctic Front, to originate east of the Kolbeinsey Ridge and to have a volume transport of around 1.5 Sv within northern Denmark Strait. The simulated multi-annual mean Atlantic Water transport of the NIIC increased by 85% during 1992 to 2006, whereas the corresponding NIJ transport decreased by 27%. Based on our model

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results we propose a new and further differentiated circulation scheme of Icelandic waters whose details may inspire future observational oceanography studies.

## 1 Introduction

The waters surrounding Iceland, flowing over the shelf and along the adjacent continental slope, form one of the most hydrographically complicated regions of the North Atlantic. The primary drivers of this complexity are topography and the interaction of four water masses. Iceland is located at the junction of the Mid-Atlantic Ridge and the Greenland–Scotland Ridge, which segments the continental shelf into four basins bounded by the Reykjanes Ridge to the south, the Kolbeinsey Ridge to the north, the Greenland–Iceland Sill (Denmark Strait) to the west and the Iceland–Faroe Ridge to the east (Fig. 1).

The water mass of primary importance for the Icelandic hydrography is the Atlantic Water which is sub-tropical in origin and therefore still comparatively warm (temperature  $T$  between 6 and 11 °C) and salty (salinity  $S$  between 35.0 and 35.2 psu) when reaching Iceland (Stefánsson, 1962). East of the Reykjanes Ridge this water mass flows northwards as part of the broad and sluggish North Atlantic Drift; a north-eastward continuation of the Gulf Stream. Along the western flank of the Reykjanes Ridge, however, the flow is more energetic. Here, the Irminger Current (IC), another Gulf Stream continuation, carries around 10 Sv of Atlantic and Subpolar Mode Water ( $1 \text{ Sv} = 10^6 \text{ m}^3 \text{ s}^{-1}$ ) northwards (McCartney and Talley, 1982; Bersch, 1995). The associated northward heat flux plays a crucial role for the marine and terrestrial climate of Iceland. South of Denmark Strait, the IC mostly recirculates towards the west and further southwards along the East Greenland continental slope. However, a small fraction (10–20 %) of the current branches off northwards through Denmark Strait and further eastwards over the North Icelandic shelf (Kristmannsson, 1998). This branch, called the North Icelandic Irminger Current (NIIC), is responsible for the mild climate north

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of Iceland and forms, to a certain extent, the lifeline of the local marine ecosystem (Vilhjálmsson, 1997).

In normal years, the Atlantic Water of the NIIC dominates most of the North Icelandic shelf area. However, on its eastward journey over the shelf the admixture of the second water mass, the Arctic Intermediate Water, becomes more and more important. This water mass, often also termed *Arctic waters*, is formed of Atlantic Water which moved into the Nordic Seas several years before and has been exposed to convective cooling and precipitation since that time. It is therefore colder ( $T$ :  $-1$  to  $4^\circ\text{C}$ ) and slightly fresher ( $S$ : 34.6 to 34.9 psu) than the Atlantic Water (Swift, 1986). The East Icelandic Current (EIC) carries Arctic Intermediate Water from the central Iceland Sea southwards along the eastern flank of the Kolbeinsey Ridge onto the north-eastern Icelandic shelf, causing the water here to be characteristically more Arctic than Atlantic. Thereafter, the EIC, whose volume flux was measured to be 2.5 Sv between June 1997 and June 1998 (Jónsson, 2007), flows eastwards along the northern flank of the Iceland–Faroe Ridge.

East of Iceland the Arctic waters of the EIC are colliding with the northward flowing Atlantic Water, leading to the formation of the *Arctic Front* which is characterised by sharp temperature gradients (Hansen and Meincke, 1979; Orvik et al., 2001). The resulting density gradient leads to differences in dynamic sea level height, with higher values to the warmer and less dense southern side of the front. The Arctic Front continues south-eastwards along the Iceland–Faroe Ridge, to the region north of the Faroe Islands. Westwards it extends north of Iceland up to Denmark Strait, separating the NIIC Atlantic Water from the Arctic waters north of it (Fig. 1). Below the NIIC exists a deep undercurrent which carries Arctic waters westwards along the north Icelandic continental slope from east of the Kolbeinsey Ridge up to Denmark Strait. This current, discovered not more than nine years ago (Jónsson and Valdimarsson, 2004), is called the North Icelandic Jet (NIJ) and seems to make a crucial contribution to the Denmark Strait Overflow, a key element of the Atlantic meridional overturning circulation (Våge et al., 2011).

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accompanying ocean current, the Icelandic Coastal Current (ICC), is still unclear. Further, even the concept of the continuous circular clockwise flow seems to contradict drift observations at the south-east coast of Iceland (Valdimarsson and Malmberg, 1999).

The importance of the coastal water flow for the marine ecosystem is beyond dispute. Not only does it strongly influence the timing and spreading of the spring bloom (Pórðardóttir, 1986), it also acts as a dispersal vector for fish eggs and larvae transported away from spawning grounds to their nursery areas, and hence plays a crucial role in the recruitment process of several fish species in Icelandic waters (Ólafsson, 1985; Marteinsdóttir and Astþórsson, 2005).

The uncertainty over the structure of the ICC is a key motivation for the present study. We also explore the general forcing of the NIIC, a current flowing northwards against the prevailing wind direction (Fig. 14) and a subject of intensive research for more than 50 yr due to its exceptional hydrographical and ecological importance for North Icelandic waters (e.g. Stefánsson, 1962; Kristmannsson, 1998; Ólafsson, 1999; Jónsson and Valdimarsson, 2005; Halldórsdóttir, 2006; Logemann and Harms, 2006). Furthermore, we examine the structure of the relatively unexplored NIJ and the path of the Atlantic Water flow towards the south and south-east coast of Iceland, a controversial component of the regional hydrography (e.g. Valdimarsson and Malmberg, 1999; Orvik and Niiler, 2002; Hansen et al., 2003).

To address these objectives we need to explore and understand the three-dimensional flow, temperature and salinity fields of the waters surrounding Iceland and beyond. We use the tool of numerical ocean modelling, which offers the possibility to obtain the requested fields with high temporal and spatial resolution covering large areas and long time periods.

The most established numerical model of Icelandic waters is a two-dimensional application of the POM ocean model (Blumberg and Mellor, 1978). It was set up for Icelandic waters by Tómasson and Eliasson (1995) and further improved by Tómasson and Káradóttir (2005). The model is run on an operational basis at the Icelandic

Maritime Administration to predict tidal and atmospherically forced sea level elevations and currents.

The first three-dimensional model study on Icelandic waters was performed by Mortensen (2004). By using an application of the MIKE3 (Rasmussen, 1991) ocean model with a resolution of 20 km horizontal and 50 m vertical his study mainly dealt with the circulation in Denmark Strait, with volume, heat and salt fluxes of the EGC and the Denmark Strait Overflow.

In 2006 three further modelling studies on Icelandic waters were published. Ólason (2006) set up the MOM4 ocean model (Griffies et al., 2004) for the region with a resolution of around 15 km horizontal and 10 m vertical near the sea surface. Driven by climatological wind fields the model successfully reproduced the basic elements of the circulation. Sensitivity experiments regarding the role of the local wind stress in forcing the near surface circulation were carried out. Halldórsdóttir (2006) applied the same model whereas her numerical experiments examined the dynamic impact of the coastal freshwater and the sensitivity of the NIIC to wind stress variations. Eventually, Logemann and Harms (2006) published their work on the high resolution (1 km horizontal, 10 m vertical) simulation of the NIIC with the ocean model CODE. Time and space variability of the NIIC volume and heat fluxes for the years 1997–2003 were analysed and the origin and composition of NIIC water masses were estimated.

For the following years the development work on the CODE model with focus on Icelandic waters was carried on (Logemann et al., 2010, 2012) which finally led to the version whose output is presented here. This resolves the entire costal area with a grid spacing of 1 km horizontal and 2.5 m vertical. It uses coastal freshwater discharge values computed by a newly developed high resolution application of the hydrological model WaSiM (Schulla and Jasper, 2007; Einarsson and Jónsson, 2010) and it assimilates hydrographic measurements like CTD (conductivity, temperature, depth) profiles into the simulation.

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Therefore, we propose that these model results could throw new light on the above mentioned questions and even enable us to propose previously not observed structures of the regional hydrography of Icelandic waters.

## 2 Model description

5 The numerical ocean model used for this study is CODE (Cartesian coordinates Ocean model with three-Dimensional adaptive mesh refinement and primitive Equations). First applied to the North Atlantic by Logemann and Harms (2006), the model was subsequently further developed with the focus on Icelandic waters. A detailed description of the current model version (9.221) with all physical equations, algorithms and numerical  
 10 techniques is given in Logemann et al. (2012). Here, we present the fundamentals of the model and outline recent improvements.

### 2.1 Equations

The primitive equations, based on the work of Bjerknes (1921), are used to approximate the oceanic flow in Cartesian coordinates  $(x, y, z)$ . The temporal change of the  
 15 horizontal flow vector  $(u, v)$  is:

$$\frac{du}{dt} = fv + \frac{\partial}{\partial x} \left( A_H \frac{\partial u}{\partial x} \right) + \frac{\partial}{\partial y} \left( A_H \frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial z} \left( A_V \frac{\partial u}{\partial z} \right) - \frac{1}{\rho_0} \frac{\partial p}{\partial x} - \frac{\partial \Phi_T}{\partial x}, \quad (1)$$

$$\frac{dv}{dt} = -fu + \frac{\partial}{\partial x} \left( A_H \frac{\partial v}{\partial x} \right) + \frac{\partial}{\partial y} \left( A_H \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial z} \left( A_V \frac{\partial v}{\partial z} \right) - \frac{1}{\rho_0} \frac{\partial p}{\partial y} - \frac{\partial \Phi_T}{\partial y}, \quad (2)$$

20 where  $f$  denotes the Coriolis parameter,  $A_H$ ,  $A_V$  turbulent exchange coefficients,  $\rho_0$  the ocean background density,  $p$  the hydrostatic pressure and  $\Phi_T$  the tidal potential, given by a first order approach (Apel, 1987). The vertical velocity component  $w$  results from

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the equation of continuity:

$$w = - \int_{z'=-D}^{z'=z} \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} \right) dz', \quad (3)$$

with  $D$  denoting the ocean depth. The vertical velocity at the sea surface ( $z = \zeta$ ) defines the temporal change of the sea surface elevation  $\zeta$

$$\frac{\partial \zeta}{\partial t} = w(z + \zeta). \quad (4)$$

The hydrostatic pressure  $p$  is computed with:

$$p = \int_{z'=z}^{z'=\zeta} \rho g dz' + p_{\text{AIR}} \quad (5)$$

in which  $p_{\text{AIR}}$  denotes the mean sea level air pressure,  $g$  the gravitational acceleration and

$$\rho = \rho(S, T, p) \approx \rho(S, T, -\rho_0 g z) \quad (6)$$

the density of seawater as a function of salinity  $S$ , temperature  $T$  and hydrostatic pressure  $p$  computed with the EOS-80 equations by Millero et al. (1980).

Temperature and salinity changes are computed with (e.g. Pedlosky, 1987)

$$\begin{aligned} \frac{\partial T}{\partial t} = & -u \frac{\partial T}{\partial x} - v \frac{\partial T}{\partial y} - w \left( \frac{\partial T}{\partial z} + \Gamma \right) + \frac{\partial}{\partial x} \left( K_{H,T} \frac{\partial T}{\partial x} \right) + \frac{\partial}{\partial y} \left( K_{H,T} \frac{\partial T}{\partial y} \right) \\ & + \frac{\partial}{\partial z} \left( K_{V,T} \left( \frac{\partial T}{\partial z} + \Gamma \right) \right) + Q_T, \end{aligned} \quad (7)$$

$$\frac{\partial S}{\partial t} = -u \frac{\partial S}{\partial x} - v \frac{\partial S}{\partial y} - w \frac{\partial S}{\partial z} + \frac{\partial}{\partial x} \left( K_{H,S} \frac{\partial S}{\partial x} \right) + \frac{\partial}{\partial y} \left( K_{H,S} \frac{\partial S}{\partial y} \right) + \frac{\partial}{\partial z} \left( K_{V,S} \frac{\partial S}{\partial z} \right) + Q_S, \quad (8)$$

in which  $\Gamma = \Gamma(T, S, \rho)$  is the adiabatic lapse rate, computed with the equation of Fofonoff and Millard (1983), whereas  $Q_T$  and  $Q_S$  denote the sum of surface heat and freshwater fluxes, respectively. These fluxes are derived by astronomic tables and by the atmospheric forcing (wind, air temperature, humidity, cloudiness) using the bulk formulas after Gill (1982).

The coefficients of horizontal turbulent exchange  $A_H$  (in Eqs. 1 and 2),  $K_{H,T}$  and  $K_{H,S}$  (Eqs. 7 and 8) are estimated using the approach of Smagorinsky (1963) assuming a Prandtl number  $Pr = A_H/K_{H,T} \approx A_H/K_{H,S}$  of 10. The coefficients of vertical turbulent exchange  $A_V$ ,  $K_{V,T}$  and  $K_{V,S}$  are computed after Pohlmann (1996) based on the approach of Kochergin (1987). Herein, the Schmidt number  $S_M = A_V/K_{V,T} \approx A_V/K_{V,S}$  is computed after Mellor and Durbin (1975).

## 2.2 Numerics

The model equations are numerically solved with the technique of finite differences in Cartesian coordinates. A three-dimensional staggered Arakawa-C-grid (Mesinger and Arakawa, 1976) with a spatially variable resolution is constructed. The equations' numerical equivalents are formulated centred in space and mostly implicitly in time. In order to avoid numerical diffusion of the advection terms a flux limiter function (van Leer, 1979) is used, which ensures the abundance of the total variation diminishing (TVD) condition.

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## 2.2.1 Adaptive mesh refinement and model domain

CODE uses a technique of adaptive mesh refinement which is oriented at the “tree-algorithm” of Khokhlov (1998). This algorithm starts with a model domain being divided by a regular three-dimensional computational mesh of basic cells. If there is an area which demands a higher resolution, each basic cell of this area is split into eight “children” with halved side lengths. Some of these children may be split further, each of them into eight “grandchildren”, those perhaps into “grand-grandchildren” and so on, until the area of interest is sufficiently resolved. The model equations are only solved for “childless” cells, but the “parent” cells are not removed from the computer memory. At each time step, they obtain the average properties of their children instead. These values may be used for numerical operations at coarser parts of the mesh.

The actual form of adaptive mesh refinement is static, i.e. it does not vary in time, and just follows topographical criteria. By using a mix of five different stereographic projections, with their projection points along the 40° W meridian, a Cartesian coordinates model domain containing the entire North Atlantic including the Arctic Ocean was constructed (Fig. 2). This domain is resolved by a basic mesh with a spacing of 128 km horizontal and 160 m vertical. First the cell thickness is refined up to 2.5 m close to the sea surface then the horizontal mesh structure is modified. The refinement begins in the Nordic Seas, the Irminger and Iceland Basin, the Canadian Archipelago and along the northern Mid-Atlantic Ridge, continues with further refinement along the Greenland–Iceland–Scotland Ridge and finally leads to a 1 km mesh along the Icelandic coast (Fig. 2).

## 2.2.2 Data assimilation

The simulated temperatures and salinities of the far field, i.e. the area south of 60° N, north of 70° N, west of 30° W and east of 5° W, are restored to the climatologic fields of the PHC (Polar Science Center Hydrographic Climatology) data set (Steele et al.,

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2001). The restoring consists of a 365-days Newtonian scheme towards the 12 monthly fields of the PHC.

However, within the highly resolved area around Iceland this restoring to climatological means, which would have led to an underestimated temporal and spatial variability, was discarded. Instead, we used the NISE (Nilsen et al., 2006) and VEINS data set (ICES, 2000) and extracted 16 802 CTD (conductivity, temperature, depth) profiles from the period 1992 to 2006 recorded between 60° N and 70° N and between 30° W and 5° W. This meant that 93 profiles per simulated month were available on average.

In order to adjust the model towards these observations we used the data assimilation technique of IAU (incremental analysis updating) processes (Bloom et al., 1996). The model performing a “free forecast” simulation was stopped when having reached the 15th of a month. The CTD data of this month was bundled and compared with the simulated fields. Assuming the CTD-derived profiles to be errorless the resulting model error profiles were horizontally interpolated, in order to create three-dimensional temperature and salinity error fields. The model was jumped one month back in time and the simulation re-started, but now with the correction terms  $\Delta u$ ,  $\Delta v$ ,  $\Delta w$ ,  $\Delta K_{H,T}$ ,  $\Delta K_{H,S}$ ,  $\Delta K_{V,T}$ ,  $\Delta K_{V,S}$ ,  $\Delta Q_T$ ,  $\Delta Q_S$  determined for every grid cell at every time step in order to correct the flow field, mixing rates or surface fluxes.

This way, Eq. (7) becomes

$$\begin{aligned}
 \frac{\partial T}{\partial t} = & - (u + \Delta u) \frac{\partial T}{\partial x} - (v + \Delta v) \frac{\partial T}{\partial y} - (w + \Delta w) \left( \frac{\partial T}{\partial z} + \Gamma \right) + \frac{\partial}{\partial x} \left( (K_{H,T} + \Delta K_{H,T}) \frac{\partial T}{\partial x} \right) \\
 & + \frac{\partial}{\partial y} \left( (K_{H,T} + \Delta K_{H,T}) \frac{\partial T}{\partial y} \right) + \frac{\partial}{\partial z} \left( (K_{V,T} + \Delta K_{V,T}) \left( \frac{\partial T}{\partial z} + \Gamma \right) \right) \\
 & + Q_T + \Delta Q_T + \Delta Q_T^{\text{NUM}}, \tag{9}
 \end{aligned}$$

whereas the correction terms are zero, with the exception of the one related to the dominant heat flux; assumed to be the cause of the error. Salinity (Eq. 8) is treated analogously. Once the 15th of the month is reached again, new error fields are computed and the corresponding correction terms are added to the previous terms before

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the model jumps back in time again and repeats the simulation. The current model version uses three of these iterations. The correction term  $\Delta Q^{\text{NUM}}$  denotes additional corrections of the simulated temperature or salinity, being activated during the last two iterations, having the function of “un-physically” correct numerical errors like numerical diffusion or erroneous initial or boundary conditions.

### 3 Simulation of the period 1992–2006

#### 3.1 Setup

The two oceanic boundaries of the model domain – slightly south of the equator between South-America and West-Africa and across Bering Strait in the Arctic – are treated as closed boundaries. Because of the far field restoring towards climatological values, the hydrodynamic implications of these boundary conditions are assumed to be negligible for Icelandic waters. Initial model data, describing the summer 1991, were taken from a model run performed by a previous model version (Logemann et al., 2010).

The atmospheric forcing of the model consists of the 6-hourly NCEP/NCAR re-analysis fields (Kalnay et al., 1996). The model reads in the following seven parameters: precipitation rate, specific humidity (2 m), sea level pressure, air temperature (2 m), total cloud cover, zonal and meridional wind speed (10 m).

During the simulation, three-hourly means of the physical ocean state, including sea ice properties, were stored. The averaging period of three hours was chosen to resolve tidal dynamics.

##### 3.1.1 Icelandic river runoff

In order to simulate the hydrodynamic impact of river runoff along the Icelandic coast, the output of the hydrological model WaSiM (Schulla and Jasper, 2007), operated by the Icelandic Meteorological Office, was used. WaSiM, a grid-based water flow and

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balance **simulation model**, was first set up and calibrated for the Icelandic environment by Jóhannesson et al. (2007). With a horizontal resolution of 1 km, geologic structures like topography, rivers, watersheds, rock types or glaciers, relevant for the hydrological cycle, are captured by the model. WaSiM is able to simulate the water cycle above and below the land surface. The latter, i.e. the groundwater flow, is regionally important in Iceland because of young porous postglacial lava fields and high hydraulic conductivity through tectonic faults and fissure swarms (Einarsson and Jónsson, 2010).

The meteorological input data, i.e. precipitation, evaporation and air temperature fields, was provided by the PSU/NCAR MM5 numerical weather model (Grell et al., 1994; Rögnvaldsson et al., 2007).

Hence, the hydrological input data for our ocean model consisted of the daily coastal freshwater discharge of 46 watersheds (Fig. 3). Because this data set does not distinguish between river and groundwater flow, we allocated the different discharge values to 58 model river mouths, i.e. we treated the discharge as being carried out completely by rivers, assuming these small-scale structures of the freshwater flow through the ocean-land interface to be not relevant for the scales of ocean dynamics examined here. The available WaSiM data covered the period 1992–2006 and thus provided the temporal range of the ocean simulation.

Figure 3 shows the seasonal variation of the discharge and its spatial variation. Along the west coast, several watersheds show higher mean winter values compared with summer values due to higher precipitation in the winter months. However, most watersheds, and especially those being fed by glacier melt, e.g. at the south-east coast, show maximum values during late spring or summer. Jónsdóttir (2008) estimated Iceland's mean overall discharge for the time period 1961 to 1990 to be  $4800 \text{ m}^3 \text{ s}^{-1}$ . She divided this value into a non-glacial ( $3900 \text{ m}^3 \text{ s}^{-1}$ ) and a glacial component ( $900 \text{ m}^3 \text{ s}^{-1}$ ). Whereas the glacial component, being rather small during most of the year, reaches its maximum at the time of highest air temperatures (July–August), the more evenly distributed and stronger non-glacial component shows a maximum, caused by snow melt, during late spring. This also dominates the seasonal pattern of the 1992–2006

discharge data we used (Fig. 3). Here, the mean discharge is  $5145\text{m}^3\text{s}^{-1}$ . This increase in comparison to the 1961–1990 mean is consistent with the local climate trend (Jónsdóttir, 2008).

## 3.2 Results and validation

5 In general, the model confirms the classical image of the circulation discussed above. The three-dimensional hydrography of Icelandic waters from 1992 to 2006 is well reproduced, including temporal anomalies, like the collapse of the NIIC during spring 1995 or its maximum in July 2003 (Jónsson and Valdimarsson, 2005). In accordance with observations (Jónsson and Valdimarsson, 2004; Våge et al., 2011) the model shows  
10 the NIJ as a deep undercurrent along the North Icelandic continental slope dominating the deep southward transport in northern Denmark Strait.

The simulated NIIC volume flux is realistic, but it has been under-estimated by previous model versions, which led to several model experiments incorporating a manipulated wind field over Denmark Strait (Logemann et al., 2010). However, not wind  
15 stress changes but the assimilation of CTD profiles finally caused the decisive jump of the simulated NIIC volume flux. This was surprising considering our numerical experiments that investigated the role of local density gradients in Denmark Strait in forcing the NIIC did not show clear results (see Sect. 4).

The simulated temperature and salinity fields of Icelandic waters are close to observations (Fig. 4), which is not surprising considering the assimilation of CTD data.  
20 However, there are still deviations between the measured and the modelled data which are primarily caused by the sparse temporal resolution of the data assimilation routine, which was called only once per simulated month; i.e. the simulated fields describing the 15th of each month were corrected towards estimations based on all measurements  
25 made during this month. The model errors at the time and location of the CTD profiles are given in Table 1.

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The simulated ocean currents are also in general agreement with observations. We compared the modelled flow field at the depth of 15 m with observations from a series of surface drifter experiments performed by Valdimarsson and Malmberg (1999). These include 19 GPS tracks of drift at the depth of around 15 m in Icelandic waters between 5 May 1998 and December 1999. By using a low-pass filter to remove tidal and shorter periods, i.e. by computing the mean drift over time intervals of 60 h, 607 drift vectors were derived. These vectors were compared with their modelled counterparts (Fig. 5).

This comparison of the flow velocity resulted in a median model error of  $-0.64 \text{ cm s}^{-1}$  with a standard deviation of  $6.54 \text{ cm s}^{-1}$ , whereas the median error of the modelled flow direction was  $4^\circ$  to the right with a standard deviation of  $67^\circ$ . A former model version without CTD assimilation showed a median velocity error of  $-2.8 \text{ cm s}^{-1}$  (Logemann et al., 2010). Considering that we did not assimilate flow measurements into the model, the improvement is remarkable.

Figure 6 shows the simulated mean flow field around Iceland at a depth of 15 m, averaged over the period 1992 to 2006. The striking features are the general eastward flow north and south-east of the island and the contrasting area of sluggish north-westerly flow in the south-west. Figure 7 gives a schematic overview of the simulated three-dimensional circulation pattern, denotes different currents and defines 16 analysis sections. The current's mean properties across these sections – volume flux, temperature and salinity – are listed in Table 2.

The definitions of the currents revealed in this study (Fig. 7) are based upon the 1992–2006 mean flow field, i.e. we refer to the long-term mean dynamic structures and do not consider the water mass composition of the flow. These definitions, comprised of positions and directions, were applied to the  $12 \times 15$  monthly mean flow, temperature and salinity fields in order to obtain the values listed in Table 2. Here, the temperature and salinity values refer to the mean properties of the fictional water body which would have been filled by the current during the simulated period of 15 yr.

With this technique we identified the following currents in Icelandic waters:



### 3.2.1 Icelandic Coastal Current (ICC) and Icelandic Coastal Undercurrent (ICUC)

We define the ICC as a near-shore ocean current being driven by a runoff induced density front, therefore directed clockwise around the island. In order to analyse the spread of the coastal freshwater over the Icelandic waters, we computed the seasonal mean freshwater thickness fields. The freshwater thickness  $h_{FW}$  is defined as the hypothetical thickness the layer of freshwater would form if it would be separated from the seawater with which it is mixed. By constraining to the upper 300 m of the water column we used:

$$h_{FW} = \int_{-300}^0 \frac{S_{REF} - S}{S_{REF}} dz \quad (10)$$

with the reference salinity  $S_{REF} = 35.2$  psu, which is assumed to be the salinity of pure Atlantic Water. Figure 8 shows the resulting simulated mean winter and summer freshwater thickness fields around Iceland.

Given the seasonality of the discharge (Fig. 3) we find surprisingly little seasonal variation of the offshore freshwater thickness. Furthermore, only along the south-west and west coast a clear riverine freshwater signal can be detected, whose northern parts are stronger in winter than in summer. This seasonality can be explained by increased wintery advection of freshwater fallen as rain or snow over the ocean south of Iceland and by the coastal freshwater discharge of this area being slightly increased during winter. Along the south-east coast, despite the great glacial discharge there, hardly any freshwater is found, not even during summer, and along the north coast we see an  $h_{FW}$  minimum in contrast to the high values of the Arctic waters of the Iceland Sea north of it.

What explains the missing freshwater along the south-east and north coast? When looking at the circulation pattern (Figs. 6 and 7) it appears that the coastal freshwater signature is missing where there is a strong direct flushing of Atlantic Water. Hence, we conclude that in these areas the freshwater is more efficiently removed by mixing and

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advection compared to the south-western and western areas which are more shielded from an exchange with the ambient ocean.

Therefore, within the 1992–2006 mean flow and salinity fields, we detected a river runoff-induced freshwater zone near the coast, a necessary precondition of the ICC, only along the south-west and west coasts. Consequently, apart from several small-scale occurrences in bays and fjords, the simulated ICC mainly exists in these areas.

Originating north-east of the Westman Islands near the mouth of the Markarfljót River, the ICC is amplified around 100 km downstream by the discharge of the rivers Þjórsá and Ölfusá. With a volume flux usually between 0.01 and 0.03 Sv the current follows the coastline in a generally north-westerly direction towards Denmark Strait where it finally mixes into the NIIC (Fig. 7). Around the Snæfellsnes peninsula (eastern end of Sect. 2) the ICC is exceptionally strong (0.08 Sv), broad and deep, pumping large amounts of freshened Faxaflói Bay water over the very steep topography to the north.

The general ICC structure is found in our model as a narrow (around 10 km) along-shore current, reaching from the sea surface down to the depth between 10 and 30 m, which is associated with a sharp horizontal salinity gradient (Fig. 9).

The effect of the Atlantic Water flushing or mixing can be seen when comparing the near-shore flow and salinity fields at Sect. 1 (Fig. 9) and Sect. 5 (Fig. 10). Across Sect. 1 we see the ICC, associated with a sharp salinity increase from below 33 psu close to the coast to values above 34 psu 20 km further offshore. However, along Sect. 5 this coastal salinity gradient is smaller by one order of magnitude (from 34.8 psu at the coast to 34.9 psu 20 km offshore) and the near-shore, wind-driven current is even directed westward, i.e. to the opposite direction of a potential freshwater driven coastal current.

With the exception of Sect. 3 where the coastal pressure field is probably already dominated by the NIIC, we find the ICC being accompanied with a counter-directed undercurrent which we call the *Icelandic Coastal Undercurrent* (ICUC) (Figs. 7 and 9). This current has a volume flux comparable to that of the ICC but has a distinctly higher salinity. Though there is no observational evidence of the ICUC, its existence would fit to observations from other coastal density fronts (e.g. Pickard, 2000) and to the

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dynamical explanation by Chapman and Lentz (1994). In brief, the down-slope Ekman transport of the ICC bottom boundary layer leads to a tilting and down-slope movement of the density front. This causes an up-slope pressure gradient force below the front and hence a geostrophic, counter-directed undercurrent. The accompanied up-slope bottom Ekman transport finally stops the down-slope movement of the front and the system becomes stationary.

### 3.2.2 Irminger Current (IC) and West Icelandic Irminger Current (WIIC)

The IC is simulated to be the significantly strongest ocean current in Icelandic waters, flowing along the continental slope west of Iceland (Figs. 6 and 7). Originating along the western flank of the Reykjanes Ridge, the current transports 10.6 Sv of Atlantic and Subpolar Mode Water with a mean temperature of 5.9°C and a mean salinity of 34.94 psu, which is in good accordance with observations (Bersch, 1995; Reynaud et al., 1995). Between the continental slope and the Icelandic coast, over the West Icelandic shelf, we find a third northward current which is rather sluggish and broad and herein called the *West Icelandic Irminger Current* (WIIC) (Fig. 7). Figure 9 shows this current flowing across Sect. 1 with its core close to the surface between kilometre 420 and 445. The WIIC originates over the continental slope north of the Reykjanes Ridge and flows northward over the western shelf until it finally joins the NIIC in the Denmark Strait. The mean volume flux is 0.2 Sv, the temperature varies seasonally between 6 and 9°C and the salinity is slightly above 35 psu.

### 3.2.3 North Icelandic Irminger Current (NIIC), North Icelandic Jet (NIJ) and East Icelandic Current (EIC)

Having reached the southern Denmark Strait the IC is deflected to the west by the Greenland–Iceland Ridge and finally recirculates southward along the Greenland continental slope. A fraction of around 1.4 Sv branches off in Denmark Strait, forms the NIIC and flows northward along the Icelandic shelf edge, which again is in agreement

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with observations (Kristmannsson, 1998; Jónsson and Valdimarsson, 2005). This current absorbs the WIIC ( $\approx 0.2$  Sv) in southern Denmark Strait. Shortly after crossing the Denmark Strait Sill, having lost around 0.2 Sv to the southwards flowing EGC, the NIIC splits into an inner (iNIIC  $\approx 0.3$  Sv) and an outer branch (oNIIC  $\approx 1.1$  Sv).

Whereas the iNIIC flows eastward along the North Icelandic coast, the oNIIC takes an outer eastward route along the North Icelandic continental slope. The iNIIC can be traced downstream to the east coast of Iceland. The oNIIC, however, broadens and increases its volume flux by entrainment of Arctic waters (1.1 Sv at Sect. 4, 1.7 Sv at Sect. 5, 2.0 Sv at Sect. 7). Before reaching the Kolbeinsey Ridge the oNIIC divides into three branches whereas the northernmost branch ( $\approx 0.9$  Sv) with a mean temperature below  $1^\circ\text{C}$  and a salinity close to 34.8 psu already shows more Arctic than Atlantic Water characteristics which may cast into doubt its denotation as an NIIC branch.

East of the Kolbeinsey Ridge the three oNIIC branches, together with Arctic waters flowing southward along the ridge, form the origin of the EIC. With a volume flux of around 1 Sv the current follows the continental slope to the east and then the northern flank of the Iceland–Faroe Ridge.

Below the EIC we find a counter-directed, cold ( $-0.5$  to  $0.4^\circ\text{C}$ ) and salty (34.876 to 34.889 psu) undercurrent; the NIJ (Figs. 7 and 11). Flowing westward along the continental slope at a depth of between 200 and 1000 m, the current reaches a volume transport above 2 Sv east of the Kolbeinsey Ridge (Sect. 10). After crossing the ridge the volume transport is reduced to 1.4 Sv (Sect. 8) and continues to decrease as the flow is approaching northern Denmark Strait. However, through Sect. 5 we still see an NIJ of 0.96 Sv with a temperature of  $0.2^\circ\text{C}$  and a salinity of 34.881 psu. Further downstream, across Sect. 4, the NIJ is simulated to swell up to 1.53 Sv, which is in agreement with the value of  $1.5 \pm 0.2$  Sv measured north-west of Iceland by Våge et al. (2011).

The Arctic characteristics of the simulated eastern NIJ ( $T = -0.46^\circ\text{C}$ ,  $S = 34.889$  at Sect. 10) are very close to those of the densest part of the Denmark Strait Overflow:  $-0.48^\circ\text{C} < T < -0.23^\circ\text{C}$ ,  $34.90 < S < 34.91$  as observed by Våge et al. (2011). Cer-

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tainly, at least a slight entrainment of the warmer and more saline NIIC water from above will modify the NIJ water. However, within our simulation this vertical heat flux was clearly over-estimated, causing the NIJ with a temperature of  $0.43^{\circ}\text{C}$  (Sect. 4) to be too warm when reaching Denmark Strait.

### 3.2.4 South Icelandic Current (SIC) and Icelandic Slope Current (ISC)

Over the southern and south-eastern Icelandic shelf the model shows an intense flow of Atlantic Water ( $7.0\text{--}7.6^{\circ}\text{C}$ ,  $35.12\text{--}35.17$  psu) towards the east and north-east, respectively. This boundary current, herein after called the *South Icelandic Current (SIC)*, has highest current speeds over the narrow shelf at the southernmost tip of Iceland at around  $19^{\circ}\text{W}$ : more than  $20\text{ cm s}^{-1}$  averaged over the 1992–2006 period (see Fig. 7). Here, the near surface core of the current is found less than 5 km south of the coastline. Further downstream this distance increases and the current broadens as the shelf broadens. Thereby additional Atlantic Water is entrained leading to an increasing SIC volume transport towards the east: 0.3 Sv at Sect. 14, 0.7 Sv at Sect. 13 and 1.7 Sv at Sect. 12. The current is nearly unaffected by horizontal density gradients and therefore shows a homogeneous vertical velocity profile from the surface down to the sea floor (Fig. 12). Finally, having reached the Iceland–Faroe Ridge, the SIC turns to a south-easterly direction, follows the ridge and opens out into the Faroe Current (FC). The FC flux of Atlantic Water north of the Faroe Islands was found to be  $3.5 \pm 0.5$  Sv (Hansen et al., 2003). Hence, our results indicate that a major part of this water could stem from the south-east Icelandic shelf or, in other words, that the FC could originate off the south coast of Iceland.

Along the south-eastern continental slope of Iceland, at the depth between 500 and 1100 m, our model shows a topographically steered deep counter-current, herein called the Icelandic Slope Current (ISC) (Fig. 7). The ISC contains the overflow of the Iceland–Faroe–Scotland Ridge but also entrains re-circulating deeper Atlantic Water which explains the increase of its volume flux between Sect. 13 (0.32 Sv) and Sect. 14 (1.13 Sv).

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### 3.2.5 Inter-annual variability of the NIIC, NIJ and SIC

Besides tidal dynamics, the temporal variability of the NIIC is mainly caused by the wind field over Denmark Strait (Ólafsson, 1999; Logemann and Harms, 2006). Logemann and Harms (2006) showed that a reduced southward wind stress north of the strait caused an increase of the NIIC volume transport since the late 1990s, reaching a maximum in 2003. This trend was accompanied by an increase of the NIIC temperature and salinity which is attributed to a weakened Subpolar Gyre circulation south of Iceland (Häkkinen and Rhines, 2004; Hátún et al., 2005). Our results confirm this structure. Regarding the NIIC volume flux they show that, in terms of the 13 month moving average, the absolute maximum of the period from 1992 to 2006 occurred in 2003. We obtain the same result when expanding the period's end from 2006 to 2010 by taking into account of the observational records of Jónsson and Valdimarsson (2012). A comparison of the modelled and observed NIIC is given in Fig. 13. Here, the mean relative model error is  $-1\%$  with a standard deviation of  $18\%$ . Note that in Fig. 13, only the Atlantic Water content of the NIIC is considered which was computed with a  $T > 4.5^\circ\text{C}$  criterion applied to the sum of the iNIIC and oNIIC crossing Sect. 5. Our simulation shows a  $85\%$  increase of the multi-annual mean NIIC; the simulated flux of Atlantic Water across Sect. 5 was  $0.54\text{ Sv}$  during the period 1992 to 1999 and rose to  $1.00\text{ Sv}$  during 2001 to 2006.

The NIJ across Sect. 5 is negatively correlated to the NIIC. A period of rather high transport,  $1.03\text{ Sv}$  during 1992 to 1999, is followed by a phase of weaker transport,  $0.75\text{ Sv}$  during 2001 to 2006; a decrease of  $27\%$ . This points again towards the wind stress as the basic driving mechanism of the temporal variability or, to be more precise, the wind-induced barotropic pressure field, whereas the anti-correlation would be the consequence of the opposite flow directions. However, Fig. 13 also shows NIJ signals which are positively correlated with the NIIC, e.g. a 2003 maximum. This could be explained by the fact that an increased NIIC intensifies the density gradient between

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its Atlantic Water and the adjacent Arctic waters, with the resulting increase in the pressure gradient accelerating the NIJ.

The SIC across Sect. 13 shows the same “remarkably stable” behaviour, at least between 1995 to 2002, as that of the FC analysed by Hansen et al. (2003). This stability excludes the variable local wind stress from being a major driving mechanism of the SIC. The 2005 minimum of the SIC could therefore be related to weakened density gradients across the Iceland–Faroe Ridge caused by an intensified NIIC. The SIC transport through Sect. 13, which solely consists of Atlantic Water, was 0.69 Sv during the period 1992 to 1999, clearly above the NIIC Atlantic Water transport at that time.

## 4 Sensitivity experiments

In order to examine the forcing mechanism behind the different simulated currents, a series of sensitivity experiments was carried out. First, the data assimilation routines were deactivated, the model was restarted at 12 July 2003 and a simulation until the end of 2003 was performed. This output, not disturbed by the corrections towards observations but fully consistent with the physical model equations, was used as the reference. A comparison of this solution with the original, including data assimilation, showed only minor deviations (experiment ADJ in Table 3) which ensures that the reference run is still realistic with just the NIJ being intensified by 29.5%. This reflects both the sensitivity of the NIJ regarding the density field and the discard of numerous CTD profiles recorded in August and November 2003 north of Iceland within the reference run.

The “local area” was then defined, i.e. the area where different forcing terms were switched off within various sensitivity experiments. We decided on a circular area having its centre at 64°36′ N, 20°56′ W, a radius of 512 km and a transition ring with the width of 64 km at its boundary where the abnormal inner conditions were linearly led back to normality (see Fig. 14).

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The following six model runs, simulating the same time period as the reference run, were carried out:

1. NOWI – no wind stress in the local area
2. locRHO – no horizontal density gradients in the local area
3. gloRHO – no horizontal density gradients in the entire model domain
4. NORO – no Icelandic river runoff
5. NOTI – no tidal forcing in the entire model domain
6. NONL – no momentum advection in the entire model domain

For each model run the August to December 2003 mean flow field and the corresponding difference of volume flux at each section relating to the reference run was computed. We assumed that a significant reduction of a current's flow rate, caused by the deactivation of a specific term, points towards an important role of the related physical process in forcing the current.

We have listed a selection of these computations in Table 3 where the most significant results are marked with bold numbers. These indicate that:

- None of the currents are primarily driven by the local wind stress.
- The ICC and ICUC are primarily driven by the coastal density gradients caused by river run-off but tide-induced residual currents are also important.
- The EGC in Denmark Strait is mainly driven by density gradients related to the Polar Front.
- The iNIIC over the north-eastern shelf, the NIJ and the ISC are forced by the local density field, mostly related to the Arctic Front or to the overflow.

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- Not more than 10 % of the NIIC in Denmark Strait can be explained by the inertia of the IC along its curved path south of the strait.
- The NIIC and SIC are predominantly driven by the barotropic pressure field related to the Arctic Front. To a lesser extent this is also valid for the EIC.

5 This last conclusion – the NIIC and SIC being mainly driven by the Arctic Front barotropic pressure field – was drawn when observing the immediate shutdown of the currents when horizontal density gradients were removed from the entire model domain (experiment gloRHO), whereas both currents increased when only the local density gradients were removed (experiment locRHO). Hence, our sensitivity experiments pointed towards the basin-scale pressure field, i.e. the difference of the sea surface height between the colder and denser waters to the north and the warmer waters to the south of Iceland, being the main forcing factor of the currents. In order to further illuminate this point an additional model experiment was carried out.

### 4.1 NIIC/SIC forcing experiment

15 In order to understand the nature of the NIIC and SIC forcing, we set up a very simple hydrodynamic scenario:

- a rectangular ocean basin at the reference latitude of 65° N with closed boundaries and side lengths of 1600 km × 1600 km
- an undisturbed ocean depth of 3000 m and a circular island of the radius of 210 km in the centre of the basin described by:

$$D(r) = 500 \text{ m} - 1750 \text{ m} (1 - \tanh(1.0472 \times 10^{-5} \text{ m}^{-1} r - \pi)) \quad (11)$$

with  $r$  being the distance from the basin centre (see Fig. 15a),

- a zonal, stationary density front separating denser water with  $1028.4 \text{ kg m}^{-3}$  in the north from less dense water with  $1027.9 \text{ kg m}^{-3}$  in the south, roughly describing

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the conditions around Iceland. The meridional density profile is given by:

$$\rho(y) = 1027.9 \text{ kg m}^{-3} + 0.25 \text{ kg m}^{-3} \left( 1 + \tanh \left( \frac{y - 800 \text{ km}}{30 \text{ km}} \right) \right) \quad (12)$$

with  $y$  being the meridional distance from the southern boundary (see Fig. 15a).

The stationary solution of this problem was determined with a simplified version of the CODE model, using a homogenous horizontal grid with a spacing of 10 km and 37 z-levels with a vertical spacing from 10 m near the sea surface to 160 m close to the sea floor. Using a time step of 30 s the model was spun up by linearly raising the density gradients from zero to the prescribed values during the first simulated week. Quasi-stationary conditions were achieved shortly afterwards (Fig. 15b–d).

Figure 15b shows the difference of sea surface height between the northern (lower level) and southern (higher level) part of the basin caused by the density difference. Like the density the sea surface height forms a front which is, distant from the island, on top of and parallel to the density front. The resulting pressure gradient force leads to an upper layer geostrophic eastward flow along the front (Fig. 15c). A counter-current is found in deeper layers (Fig. 15d).

However, close to the island, this structure is distorted. When hitting the island, the upper eastward flow causes a zone of high pressure at the island's western (windward) coast, and a low pressure zone at the eastern (lee side) coast. The consequences are two geostrophic northward currents along the west and the east coast, extending to the north and south coast respectively. These two currents have a clear similarity to the NIIC and SIC.

These results point to an NIIC/SIC forcing which consists of a topographically induced distortion of the Arctic Front barotropic pressure field. This would explain why both currents could only be stopped by switching off the Arctic Front within experiment gloRHO. It would also explain the NIIC flowing against the mean wind stress (Fig. 14) and in addition the zone of low sea surface height along the south-east coast recently detected by satellites (Hunegnaw et al., 2009) which cannot be attributed to

local density or wind stress structures (Fig. 17). Note also, the similarity between the observed and modelled sea surface elevation fields shown in Fig. 17 and those of the idealised case discussed here (Fig. 15b).

## 5 Discussion and conclusions

In this paper we have analysed a hydrodynamic simulation of Icelandic waters covering the time period 1992 to 2006. Thereby, we have concentrated particularly on the temporal mean state derived from the model output. However, we also presented the simulated temporal variability of the involved ocean currents which partly contains considerable inter-annual fluctuations. Furthermore, we know that the regional marine climate occasionally has undergone dramatic changes (Malmberg and Kristmannsson, 1992). Hence, this paper should not be understood as an attempt to specify ever-fitting structures of a stationary system, but rather as a proposed description of the current state.

The model results indicate that within the long-term mean flow field a distinct ICC exists only to the south-west of Iceland. Only in this coastal region between the Westman Islands to the south and the Látrabjarg tongue to the north, are the coastal waters sufficiently protected from a direct flushing of Atlantic Water and the freshwater discharge sufficiently large to enable the almost persistent formation of the coastal freshwater-induced density front. North of Látrabjarg and further downstream along the north-west and north coast, the North Icelandic Irminger Current (NIIC) dominates the near-shore circulation and erodes most of the coastal freshwater signatures. However, in more shielded areas like the Húnaflói Bay or within the large western and northern fjords, the ICC shows sporadic appearances which is in agreement with observations (Ólafsson et al., 2002). This also applies, to a lesser extent, to the south-east coast. Here, a counter-directed, north-eastward flow of Atlantic Water, the South Icelandic Current (SIC), not only erodes the coastal freshwater signature but also the related south-westward momentum.

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Hence, our model results offer a solution to the ICC quandary, which is defined by two opposing schemes of the coastal circulation around Iceland: (a) the classical view of a freshwater-induced current flowing clockwise around the island (e.g. Stefánsson and Ólafsson, 1991; Halldórsdóttir, 2006); and (b) the assumption that freshwater-induced near-shore dynamics do not form a separate current, with the coastal circulation instead thought to derive from the offshoots of the larger ocean currents further off-shore (e.g. Astþórsson et al., 2007). Our findings point to the possibility that both views are correct when applied to different coastal sections. They illustrate the transport of freshwater along the south-west coast in accordance with the measurements of Stefánsson and Guðmundsson (1978) and Ólafsson et al. (1985, 2008), but also explain the sparse occurrence of polar driftwood at south-eastern beaches which is in sharp contrast to the large deposits often found at north-eastern beaches (Eggertsson, 1994) – an observation that indicates the absence of a steady southward current connecting these areas.

Another result of this study is the possible existence of an undercurrent below the ICC, which we have called the *Icelandic Coastal Under-Current* (ICUC). Unfortunately there are no long-term current measurements from the depth range within the shallow near-shore waters along the south-west coast where we predict the ICUC to occur. We are therefore unable to confirm or refute our model predictions; however, the simulated structure is compatible with the theoretical predictions of ocean physics. These predict a counter-directed undercurrent if an along-shore density front exists which reaches down to the bottom-boundary layer (Chapman and Lentz, 1994; Pickart, 2000).

One might wonder whether the simulated undercurrent could have been caused by a numerical error which appears along the boundary between domains of different mesh refinement. The background of this question forms the widespread assumption of trapped or reflected kinetic energy at those boundaries in adaptive mesh ocean models (Griffies et al., 2000). However, in accordance with Popinet and Rickard (2007) we found the main reason for this problem to be the formulation of the discrete spatial operators, i.e. their accuracy and smoothness, across the resolution boundaries. Furthermore, regarding the model solution discussed here, the ICUC as a numerical



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(Hansen et al., 2010; Richter et al., 2012; Sandø et al., 2012). Herein, the meridional gradient of sea surface height across the Arctic Front, caused by the density gradient or even by the removal of dense water by the overflow (Hansen et al., 2010), is clearly identified as the basic forcing of the FC. Therefore, the assumption of an analogous forcing of the SIC and NIIC appears to hold true.

The NIIC is simulated to bifurcate north of Denmark Strait into the iNIIC which flows eastward along the north Icelandic coast, and the oNIIC which follows the continental slope north of Iceland. Whereas the iNIIC can be traced downstream up to the north-east coast of Iceland, the oNIIC only reaches up to the Kolbeinsey Ridge. Here, the current, which has further ramified into three sluggish branches, finally mixes into the Arctic waters of the East Icelandic Current (EIC) which flows southward along the eastern flank of the ridge. Further downstream the EIC forms the third eastward surface current in Icelandic waters, which can be attributed to the Arctic Front.

The NIIC is the result of the signal of high dynamic sea level height south of Iceland which is led downstream along the west and north-west coasts. An analogous structure is found along the east and south-east coasts where the signal of low dynamic sea level height from north of Iceland is led southwards (Fig. 17). Here the result is the SIC, simulated to flow with high intensity over the southern and south-eastern shelf to the east and north-east, respectively. The model showed months with the SIC being stronger than the NIIC or EIC, and indicated that the SIC is a major source of the FC, and could even be interpreted as the FC preform.

We successfully reproduced the NIIC/SIC structure with an idealised model setup: a circular island being placed on a zonal density front. This experiment resembles those of Hsieh and Gill (1984) addressing the Rossby adjustment problem (Rossby, 1937, 1938). Considering a meridional channel with a zonal density front Hsieh and Gill pointed to the existence of a northward western boundary current south of the density front and a northward eastern boundary current north of it, both being accompanied by deep counter-currents. They also discussed the application of their results to the hydrography of the Iceland–Faroe Ridge.

However, are our model predictions of the SIC realistic? After all, a description of a specific eastward current over the southern and south-eastern Icelandic shelf, independent and separated from the North Atlantic Drift does not exist within the classical view of Icelandic hydrography.

5 On the one hand, the near-surface flow field of the northern Iceland Basin is assumed to be predominantly topographically steered and cyclonic; perhaps a remnant of the circulation scheme of Nansen (1912), though his hypothesis referred to deeper layers. Here a broad (> 100 km) and sluggish south-westward current of coastal and Atlantic Water is assumed along the south-east coast of Iceland (e.g. Stéfan­sson, 1962; see  
10 Fig. 16a). Furthermore, the source of the Atlantic Water of the FC is thought to stem mainly from the area north-west of the Faroe Bank, where the Atlantic Water flows north-westwards along the southern flank of the Iceland–Faroe Ridge until it crosses the ridge close to the Icelandic shelf to form the current (see Fig. 16c) (e.g. Hansen et al., 2003; Østerhus et al., 2005). However, Larsen et al. (2012) state that the Atlantic  
15 Water characteristics of the FC north of the Faroe Islands point towards a considerable admixture from south of Iceland, and Hansen et al. (2003), referring to Orvik and Niiler (2002), do mention the possibility of an “alternative”, north-eastward path of the source waters (Fig. 16c).

Such a dissentient path of the FC’s source waters was first proposed by Orvik and Niiler (2002) based on an analysis of surface drifter tracks south-west of Ice-  
20 land (Fig. 16d). This is in agreement with our findings, regarding the bifurcation of the Atlantic Water flow into an eastward and a westward branch, with the eastward branch being the source path of the FC, although Orvik and Niiler (2002) made no reference to different dynamics of the eastward branch compared to the “wider, eddy  
25 structured” flow through the Iceland Basin south of Iceland. Several other empirical studies also lend support to this alternative view of the source path. For example, Hermann and Thomsen (1946), who published a circulation scheme based on drift bottle measurements, showed a clear north-eastward drift south-east of Iceland and in combination with the eastward flux along the Arctic Front, their scheme showed an anti-

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cyclonic structure of the near-surface circulation in the northern Iceland Basin. These pattern has subsequently been re-verified by modern drifter experiments (Perkins et al., 1998; Valdimarsson and Malmberg, 1999; Jakobsen et al., 2003) (Fig. 16b, d) and the inclusion of CTD profiles and records of moored current meters from the area south-east of Iceland (Perkins et al., 1998). In accord with our results, Perkins et al. (1998) have described an intense north-eastward flow of Atlantic Water at the shelf break south-east of Iceland, being forced by the dynamic height gradients of the Arctic Front.

Hence, the CODE simulation clearly supports a scheme of anti-cyclonic near surface circulation in the northern Iceland Basin. Though it shows a distinct increase of the SIC between Sect. 13 (0.7 Sv) and Sect. 12 (1.7 Sv), i.e. a swelling of the current by absorption of Atlantic Water from the south shortly before hitting the Iceland–Faroe Ridge. And though deeper portions of this water, being part of a deep, topographically steered slope current, may indeed stem from north-west of the Faroe Banks. In our simulation, the majority of the near-surface Atlantic Water east of Iceland is steered by the barotropic pressure field of the Arctic Front, which implies an eastward flow component over the Iceland–Faroe Ridge.

With the exception of the extensive field work of Perkins et al. (1998), which was however restricted to the shelf east of 14° W, neither the drifter studies indicate a distinct SIC, i.e. an energetic dynamical structure over the south-east Icelandic shelf being independent from the North Atlantic Current further offshore. If we assume our simulation to be realistic, what could be the reason for the past invisibility of this current?

Our simulation shows very homogenous vertical current profiles of the SIC (Fig. 12), reflecting its forcing by a near-coastal signal of low sea level height, independent from any local density gradient. This means that the SIC remains invisible when the dynamic method, based on CTD profiles, is applied. On the other hand, it is difficult to deduce a boundary current structure like the SIC from a limited number of surface drifter tracks. In addition, if we consider the fact that, in the Northern Hemisphere an eastward flow along a south coast forms an upwelling-favourable situation, we should assume a divergent near-surface flow field within the SIC. Hence, surface drifter would virtually



be repelled from the current's core and most of the SIC would remain invisible when looking at the drifter tracks.

However, we found observational evidence for the SIC when comparing drifter tracks (Valdimarsson and Malmberg, 1999) with the simulated flow field. Figure 5 shows the striking similarity between the observed and simulated eastward flow vectors south of Iceland. Note that in Fig. 5 the longest red vector, i.e. the fastest observed drift, is located south-east of Iceland and points to the east.

Whereas the numerical simulation of Nilsen et al. (2003) already comprised a sparsely resolved SIC, the work of Hunegnaw et al. (2009) revealed further details. Their dynamic topography, calculated from marine, airborne and satellite gravimetry, combined with satellite altimetry, confirms our model results showing a strong SIC signal along major parts of the south-east coast and even a weak, probably just barely resolved, signal of eastward flow along the south coast (Fig. 17).

Hence, we assume that, in the absence of direct current measurements over the southern shelf, evidence for the SIC arose only after the emergence of high resolution numerical ocean modelling (this study) or satellite altimetry (Hunegnaw et al., 2009). Therefore, it may be a new challenge for observational oceanography to verify the SIC postulated here.

Another current in Icelandic waters which has just recently been discovered (Jónsson and Valdimarsson, 2004) is the North Icelandic Jet (NIJ). Knowledge on the structure of the NIJ is still limited. However, our model results are in good accordance with the observations (Jónsson and Valdimarsson, 2004; Våge et al., 2011), whereby the NIJ is predicted to flow from east of the Kolbeinsey Ridge as a deep undercurrent along the north Icelandic continental slope with a volume flux of 1.5 Sv when entering Denmark Strait. An analysis of the sources and pathways of the Denmark Strait Overflow Water was beyond the scope of this paper. However, the simulated temperature and salinity values of the NIJ east of the Kolbeinsey Ridge are very close to those of the densest part of the Denmark Strait Overflow as observed by Våge et al. (2011). Further downstream the simulated heat transfer from the overlaying NIIC into the NIJ is

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over-estimated, and the corresponding salt transfer may be under-estimated. This was probably caused by an insufficient vertical resolution (80 m) within the NIJ depth range, which could therefore not be effectively corrected by the CTD data assimilation. The consequence is a simulated NIJ that is too warm and too fresh when entering Denmark Strait. Thus, the simulated density over the sill is under-estimated and this may be the main reason of the under-estimated volume flux of the overflow in this study (simulated 1.33 Sv in contrast to  $2.9 \pm 0.5$  Sv observed by Våge et al., 2011).

We found the NIJ to be forced by local density gradients. These are caused by the collision of the warmer Atlantic Water of the NIIC with the Arctic waters north of Iceland, i.e. by the Arctic Front. Like the SIC, the NIJ leads to an up-slope secondary circulation. Convectively formed Arctic waters from deeper central parts of the Iceland Sea could be pumped up the North Icelandic continental slope and finally onto the sill of Denmark Strait. This would explain the NIJ's dominant role in providing the densest parts of the overflow (Våge et al., 2011).

In terms of the temporal variability in the NIJ, the simulation indicates two primary characteristics. First, a trend of decreasing volume flux during the period 1992–2006. This decrease was most pronounced during the years 1999 and 2000 when the 13-month moving average of the westward volume flux north-west of Iceland (Sect. 5) dropped from 1.5 Sv down to 0.6 Sv. This may be related to the observed decrease of the Denmark Strait Overflow during the years 2000 to 2003 (Macrander et al., 2005). The trend of NIJ decrease is accompanied by a trend of NIIC increase. Logemann and Harms (2006) showed that the latter is mainly wind-induced. Hence, we believe this to be valid as well for the anti-correlated, counter-directed NIJ.

Secondly, we also see positively correlated signals between the NIIC and NIJ, such as the 2003 maximum. An explanation for this could be an increased density contrast of the Arctic Front, caused by an increased NIIC forming a stronger NIJ forcing. Further studies should examine this mechanism and its impact on the variability of the Denmark Strait overflow as well as the formation processes of the NIJ water, which may become a key issue for our understanding of the Atlantic meridional overturning circulation.

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In conclusion, our numerical ocean model CODE, established on the basis of the differential equations of ocean physics together with hydrographic measurements, has given us a number of new insights into the circulation of Icelandic waters. We hope it could contribute to a further clarification of certain objects of the regional oceanography – the structure of the ICC, the primary forcing of the NIIC and the circulation patterns south-east of Iceland. We have extracted several detailed and previously unknown structures (e.g. the SIC, the ICUC or the NIIC bifurcations) and determined the primary forcing mechanism of all currents within Icelandic waters. Of course, these postulates require observational verification and an expansion of the simulation's temporal range. This would provide further insights on the relevance of our results to the Icelandic marine ecosystem, the local circulation's role within the Atlantic meridional overturning circulation and its behaviour in a changing marine climate.

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**Table 1.** Model temperature and salinity errors within Icelandic waters during the period 1992–2006 at the location and time of all available CTD profiles. Listed are the mean and the median model errors as well as the standard deviation  $\sigma$  of the mean error.

Depth range (m)	Number of obs.	$\Delta T_{\text{mean}}$ (K)	$\Delta T_{\text{median}}$ (K)	$\sigma_T$ (K)	$\Delta S_{\text{mean}}$ (psu)	$\Delta S_{\text{median}}$ (psu)	$\sigma_S$ (psu)
0–10	35 794	0.28	0.09	1.27	0.120	0.004	0.562
10–20	28 879	–0.07	–0.10	1.21	0.097	0.007	0.434
20–30	28 103	–0.26	–0.24	1.22	0.080	0.008	0.478
30–50	54 647	–0.27	–0.26	1.17	0.062	0.007	0.423
50–100	119 459	–0.24	–0.23	1.08	0.037	0.001	0.335
100–150	59 874	–0.28	–0.23	1.11	0.011	–0.006	0.317
150–200	51 926	–0.31	–0.25	1.21	0.004	–0.008	0.327
200–300	80 490	–0.39	–0.29	1.31	–0.007	–0.011	0.393

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**Table 2.** Simulated 1992–2006 mean volume flux, temperature and salinity of the currents in Icelandic waters across the 16 analysis sections. See Fig. 7 for the locations of the sections and for the abbreviations of ocean current names.

Section	Current	Flux [Sv]	$T$ [°C]	$S$ [psu]	Section	Current	Flux [Sv]	$T$ [°C]	$S$ [psu]
1	ICC	0.012	6.98	33.481	8	oNIIC3	0.92	1.45	34.824
1	ICUC	0.016	6.93	34.482	8	NIJ	1.39	−0.22	34.887
1	WIIC	0.20	6.98	35.043	9	iNIIC	0.33	4.05	34.816
1	IC	10.62	5.89	34.937	9	EIC	1.26	2.32	34.824
2	ICC	0.079	6.82	34.889	9	NIJ	1.04	−0.16	34.885
2	ICUC	0.049	6.35	35.060	10	iNIIC	0.23	3.88	34.774
3	ICC	0.021	5.93	34.736	10	EIC	0.90	2.55	34.802
3	NIIC	1.58	6.16	35.036	10	NIJ	2.20	−0.46	34.889
3	OF	1.33	1.18	34.894	11	iNIIC	0.07	4.28	34.678
4	iNIIC	0.30	5.95	34.963	11	EIC	0.57	2.29	34.780
4	oNIIC	1.07	5.34	34.986	11	NIJ	0.35	−0.19	34.889
4	NIJ	1.53	0.43	34.876	12	SIC	1.70	7.24	35.141
4	EGC	1.15	0.09	34.520	13	SIC	0.70	7.53	35.140
5	iNIIC	0.46	5.66	34.943	13	ISC	0.32	6.74	35.152
5	oNIIC	1.68	2.56	34.859	14	SIC	0.31	7.49	35.124
5	NIJ	0.96	0.21	34.881	14	ISC	1.13	7.04	35.161
6	iNIIC	0.42	5.14	34.905	15	ICC	0.010	6.74	34.585
7	oNIIC	2.02	2.32	34.858	15	ICUC	0.045	7.87	35.033
7	NIJ	1.23	0.09	34.868	15	SIC	0.43	7.62	35.169
8	iNIIC	0.12	4.75	34.862	16	ICC	0.033	6.86	35.007
8	oNIIC1	0.37	4.04	34.858	16	ICUC	0.009	7.73	35.026
8	oNIIC2	0.56	2.98	34.855					

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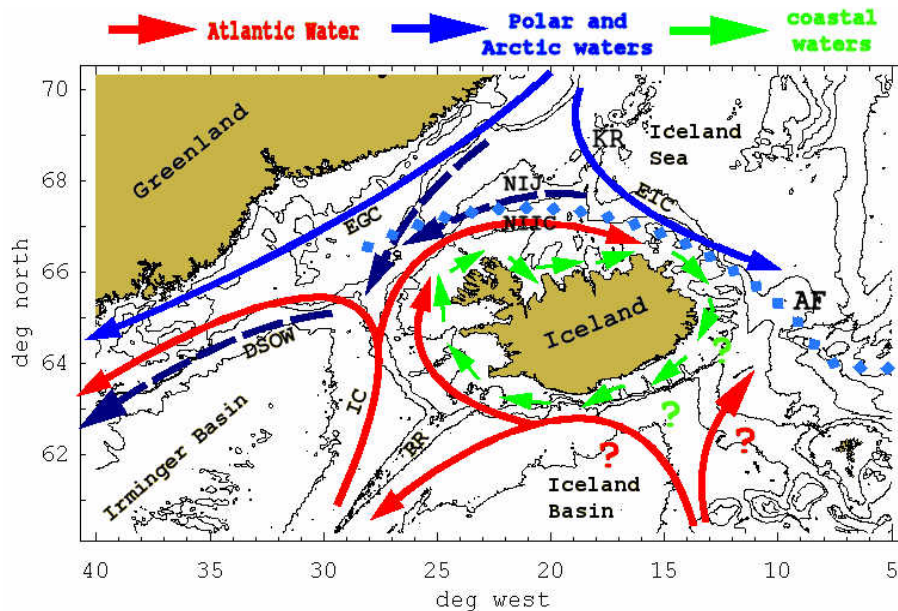
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**Fig. 1.** Bathymetry around Iceland and the classical view of the ocean circulation. The isobaths are: 200, 500, 1000, 2000 and 3000 m. The abbreviations are: AF – Arctic Front, DSOW – Denmark Strait Overflow Water, EGC – East Greenland Current, EIC – East Icelandic Current, IC – Irminger Current, KR – Kolbeinsey Ridge, NIIC – North Icelandic Irminger Current, NIJ – North Icelandic Jet, RR – Reykjanes Ridge. The question marks indicate questionable structures like the coastal current. Modified after Logemann and Harms (2006).

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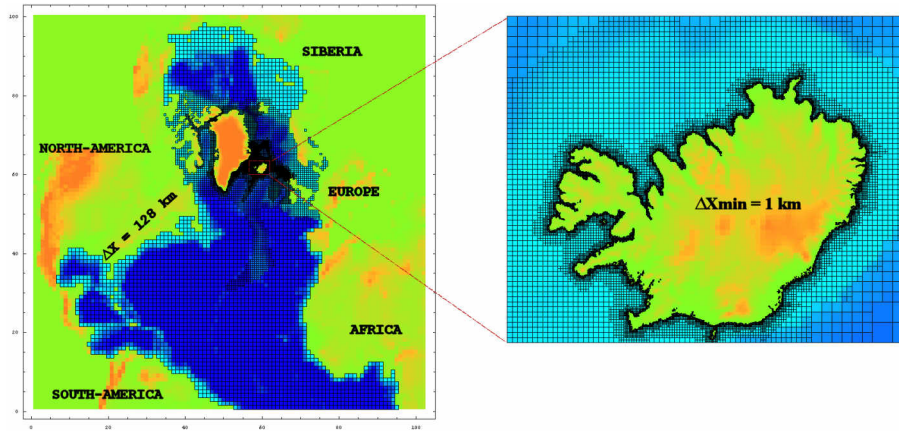
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**Fig. 2.** Left: the model domain in Cartesian coordinates with a basic mesh size of 128 km. Right: the highly resolving (up to 1 km) grid around Iceland.

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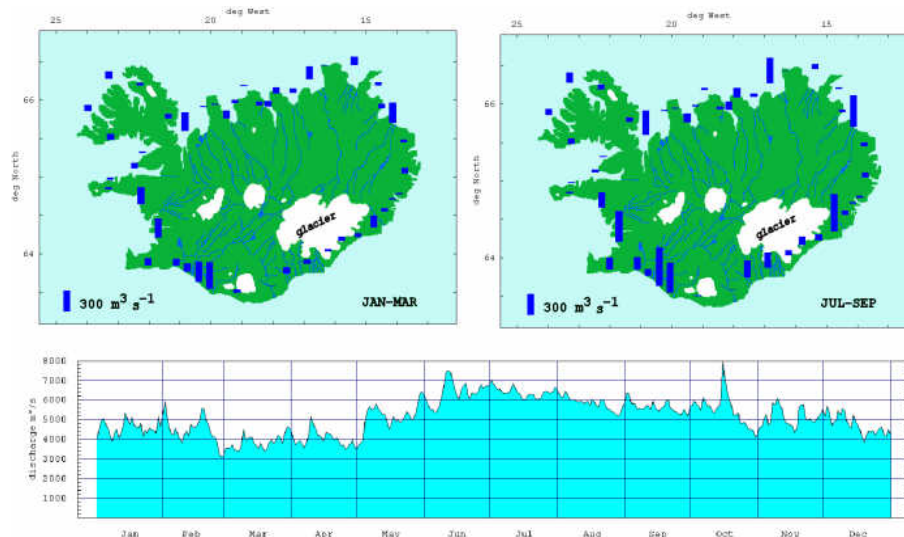
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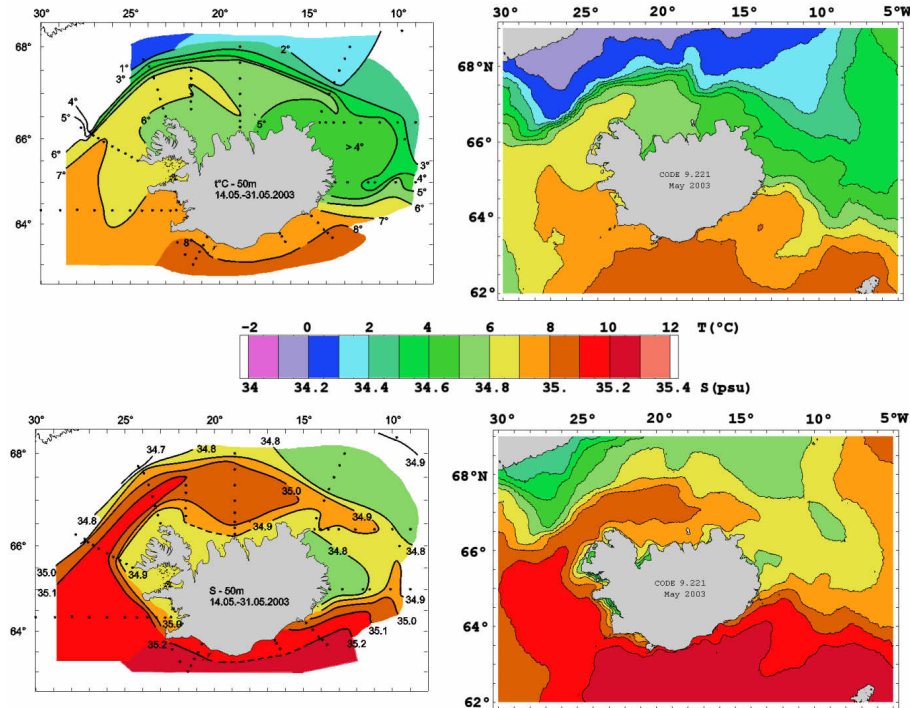


**Fig. 3.** Winter (left panel) and summer (right panel) mean discharge of 46 Icelandic watersheds for the time period 1992 to 2006 simulated with WaSiM. Below the simulated mean seasonal signal of the island's overall discharge for the same time period is shown.



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**Fig. 4.** Observed (left panels) and simulated (right panels) temperature (upper row) and salinity (lower row) in May 2003 at the depth of 50 m. Observational based charts are drawn after charts published by the Marine Research Institute, Iceland ([www.hafro.is/Sjora/](http://www.hafro.is/Sjora/)). The black dots show the location of CTD stations.

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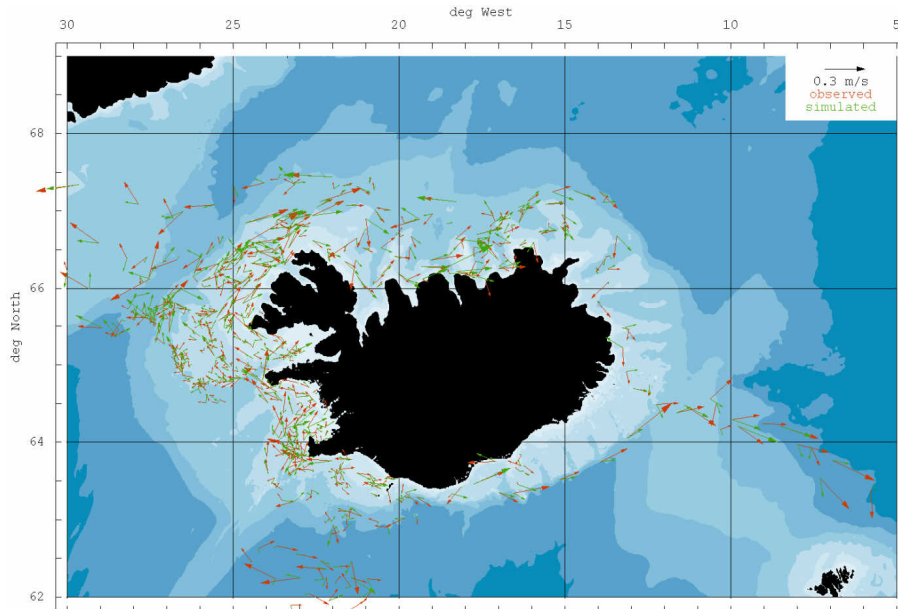
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**Fig. 5.** Observed (red) and simulated (green) drift vectors at 15 m depth between May 1998 and December 1999. Observed vectors are based on the surface drifter experiments by Valdimarsson and Malmberg (1999).

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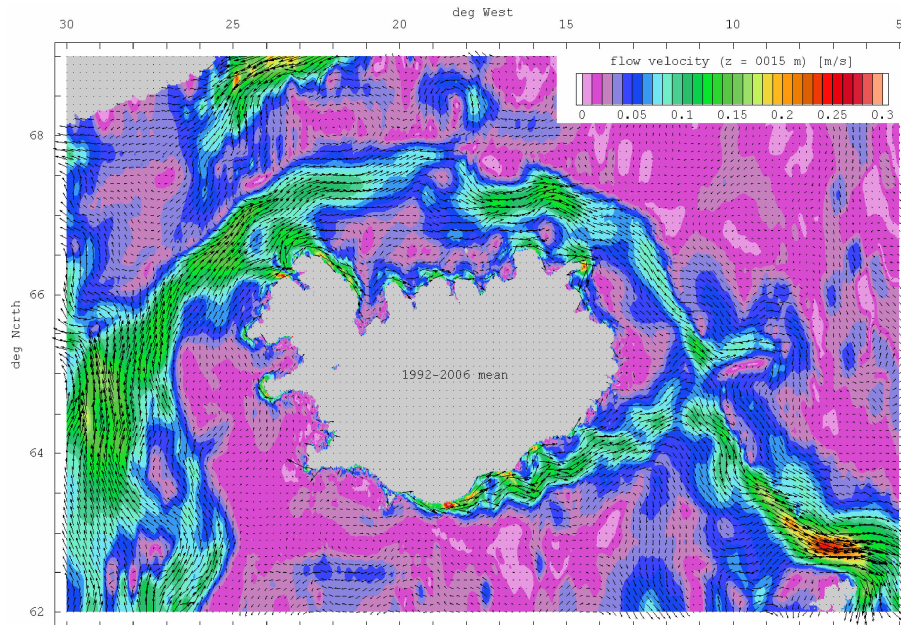
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**Fig. 6.** Simulated mean flow field around Iceland at 15 m depth, averaged over the period 1992 to 2006.

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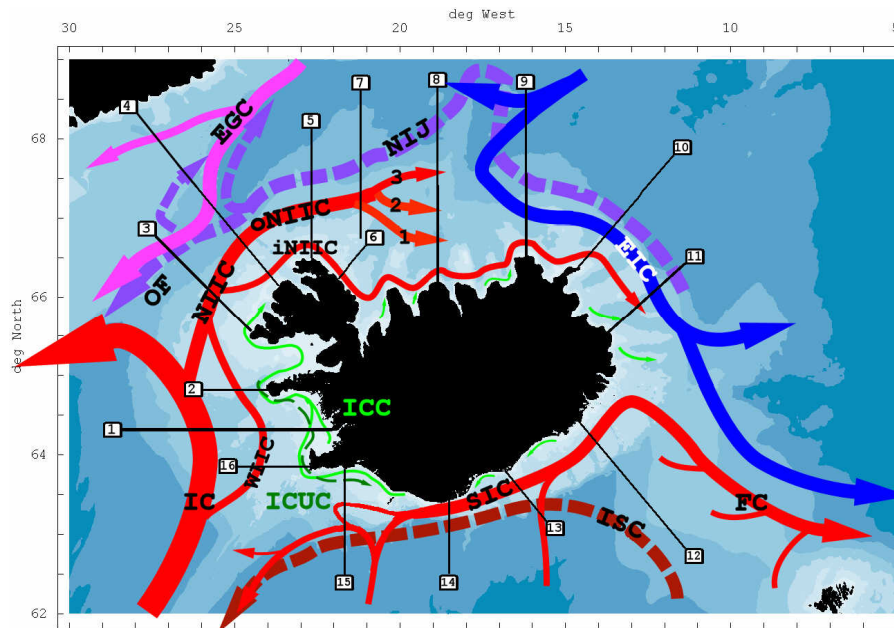
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**Fig. 7.** Proposed three-dimensional circulation scheme of Icelandic waters with the locations of the 16 analysis sections. Dashed arrows denote deep currents. The abbreviations are: EGC – East Greenland Current, EIC – East Icelandic Current, FC – Faroe Current, IC – Irminger Current, ICC – Icelandic Coastal Current, ICUC – Icelandic Coastal Undercurrent, iNIIC – inner NIIC, ISC – Icelandic Slope Current, NIJ – North Icelandic Jet, NIIC – North Icelandic Irminger Current, OF – Overflow, oNIIC – outer NIIC, SIC – South Icelandic Current, WIIC – West Icelandic Irminger Current.

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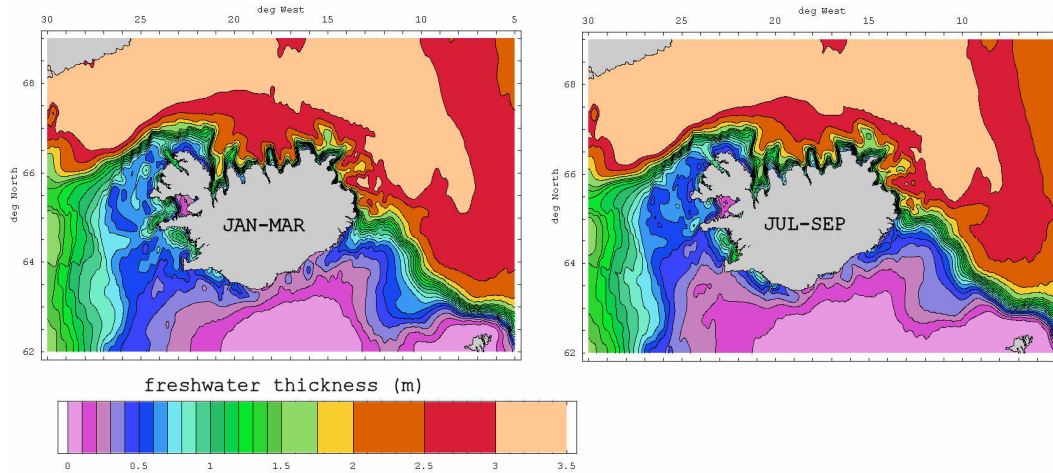
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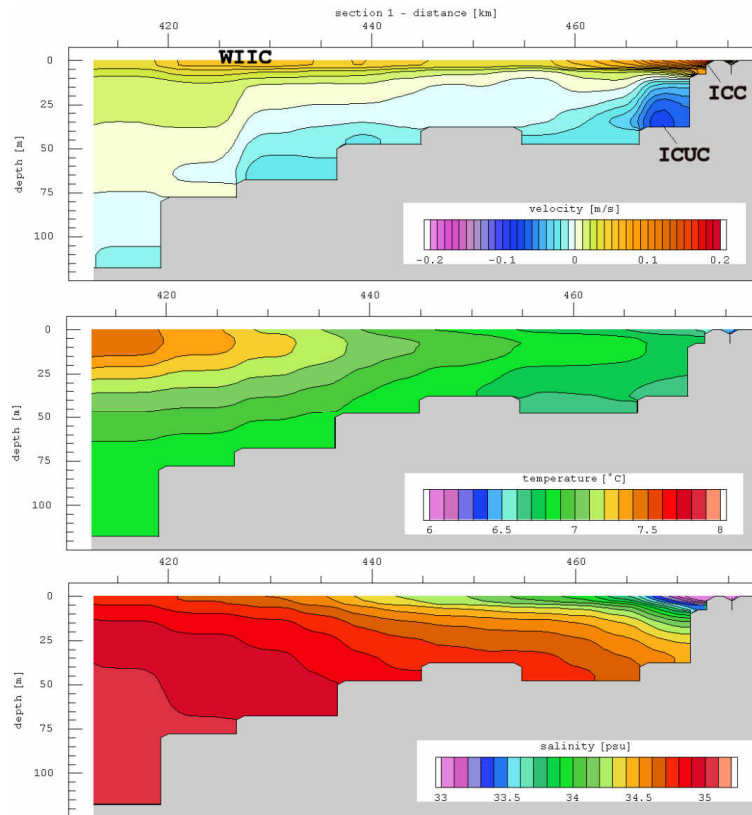


**Fig. 8.** Mean simulated wintery (left) and summerly (right) freshwater thickness of Icelandic waters for the time period 1992 to 2006.

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**Fig. 9.** Simulated 1992–2006 mean of flow (positive (red) values denote northward flow), temperature and salinity across Sect. 1 (eastern end). See Fig. 6 for section location.

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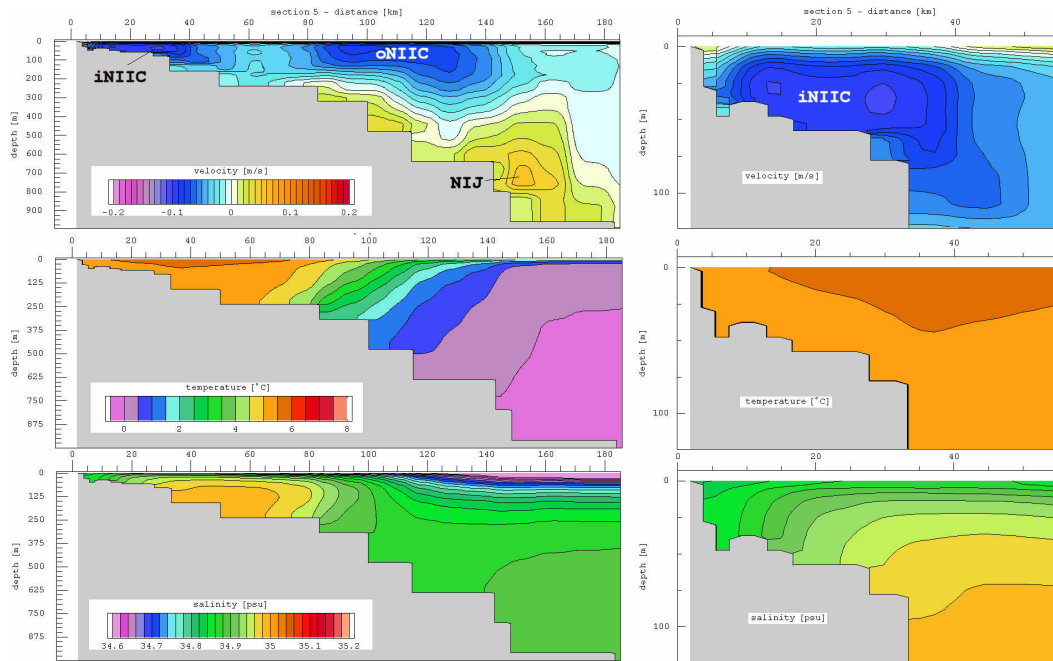
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**Fig. 10.** Simulated 1992–2006 mean of flow (positive (red) values denote westward flow), temperature and salinity across Sect. 5. See Fig. 6 for section location.

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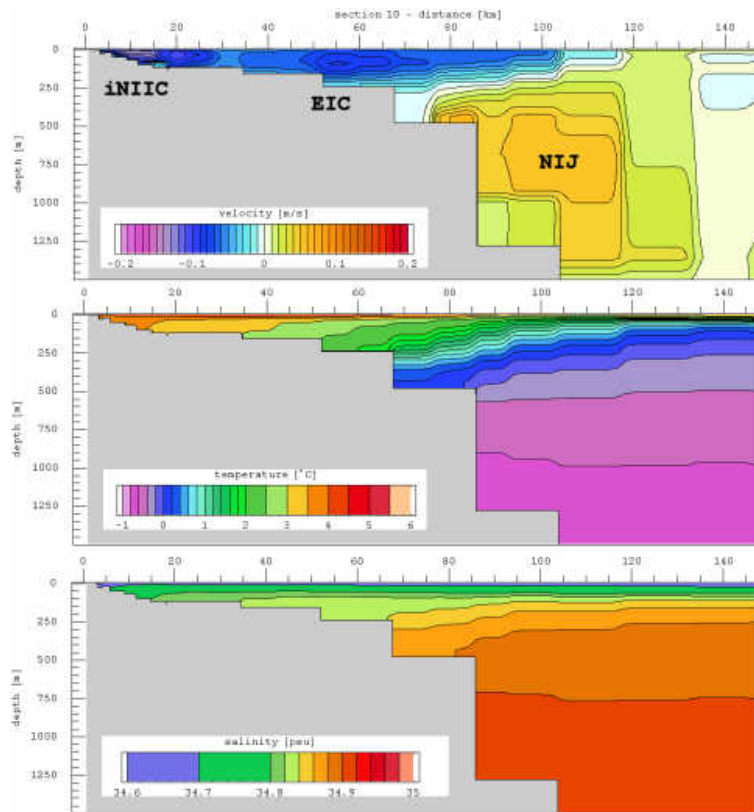
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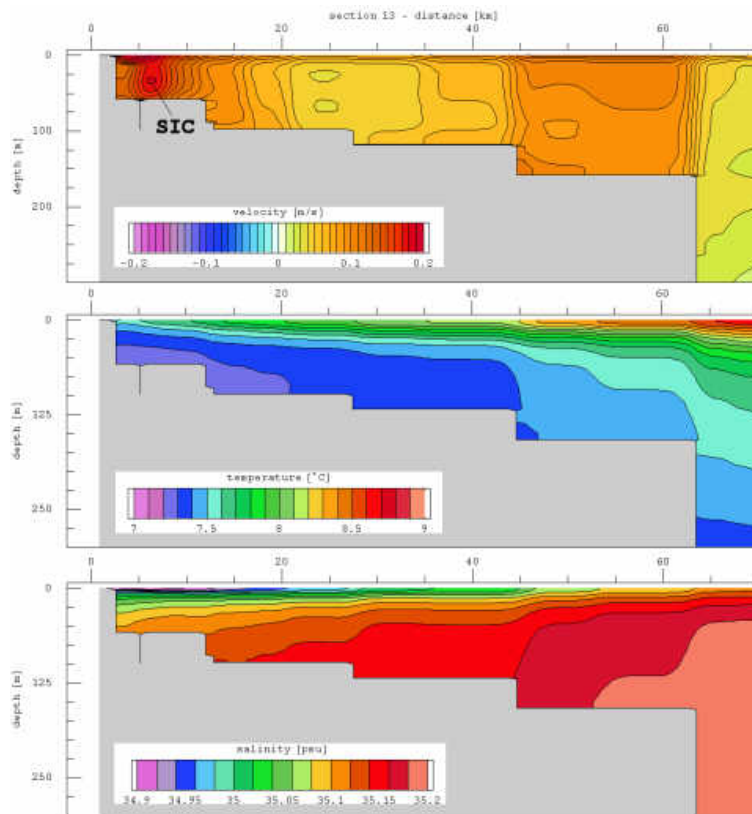
**Fig. 11.** Simulated 1992–2006 mean of flow (positive (red) values denote north-westward flow), temperature and salinity across Sect. 10. See Fig. 6 for section location.

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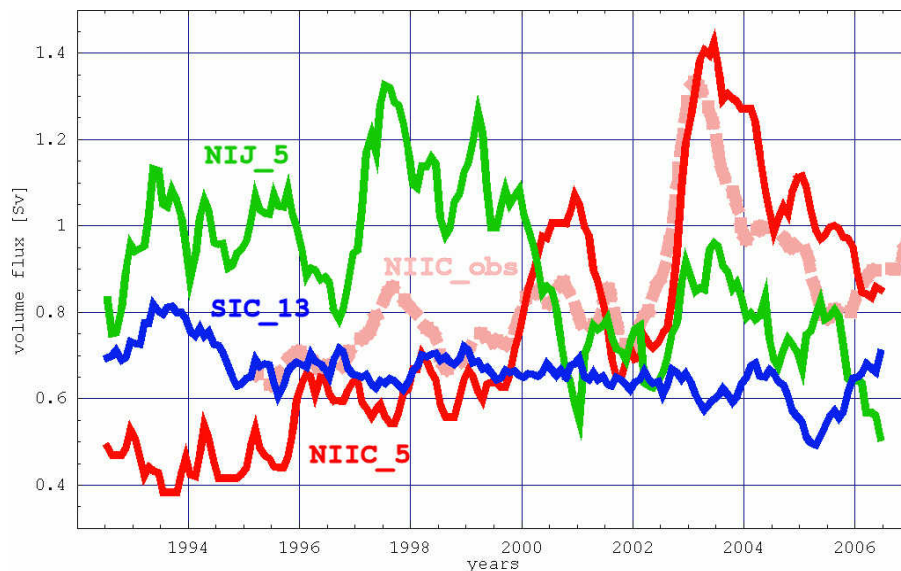


**Fig. 12.** Simulated 1992–2006 mean of flow (positive (red) values denote north-eastward flow), temperature and salinity across Sect. 13. See Fig. 6 for section location.

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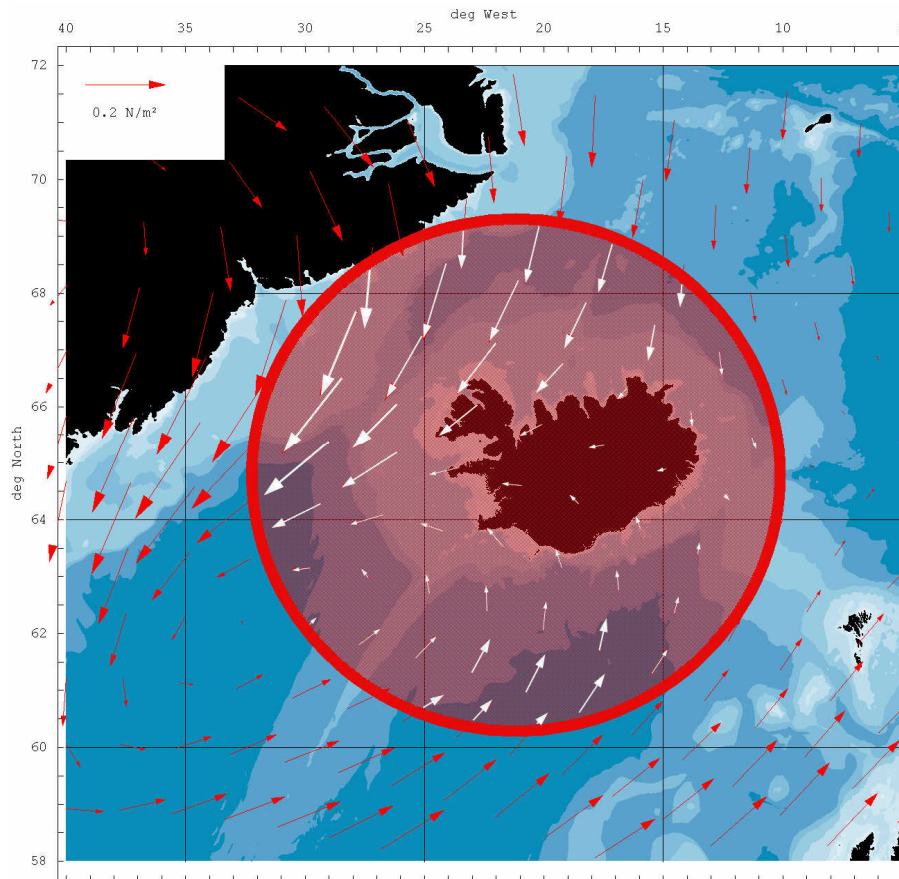
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**Fig. 13.** Time series (13 months moving average) of simulated and observed volume fluxes around Iceland. Green: simulated North Icelandic Jet (NIJ) across Sect. 5; red: simulated Atlantic Water ( $T > 4.5^{\circ}\text{C}$ ) transport of the North Icelandic Irminger Current (NIIC) across Sect. 5; light red dashed: Atlantic Water transport of the NIIC close to Sect. 5 derived from current meter data (after Jónsson and Valdimarsson, 2012); blue: simulated South Icelandic Current (SIC) across Sect. 13.

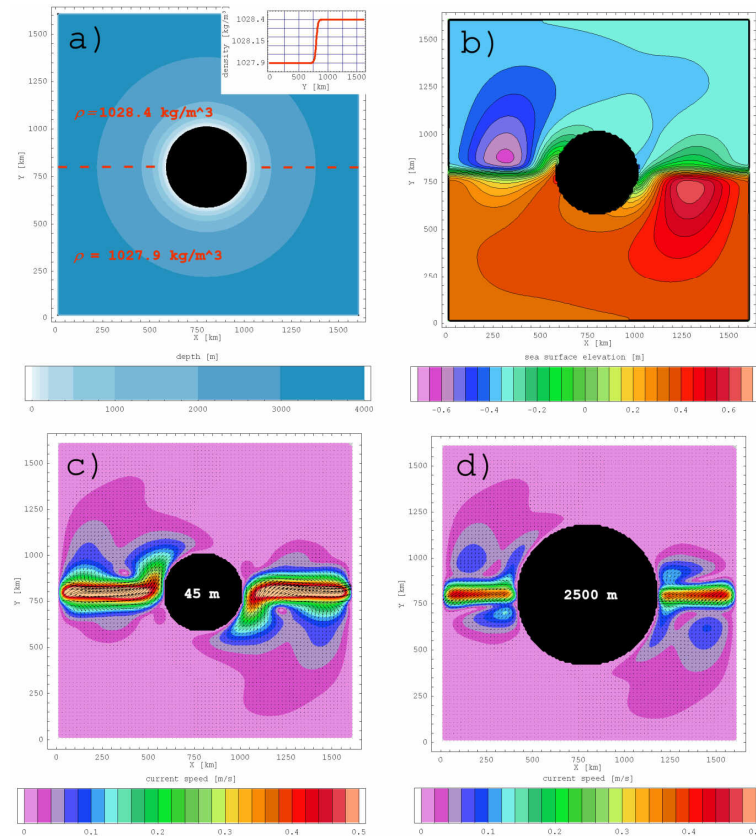
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**Fig. 14.** Bathymetry and mean surface wind stress averaged over the period 1992 to 2006. In the frame of the various sensitivity experiments different forcing terms were switched off within the red encircled area.

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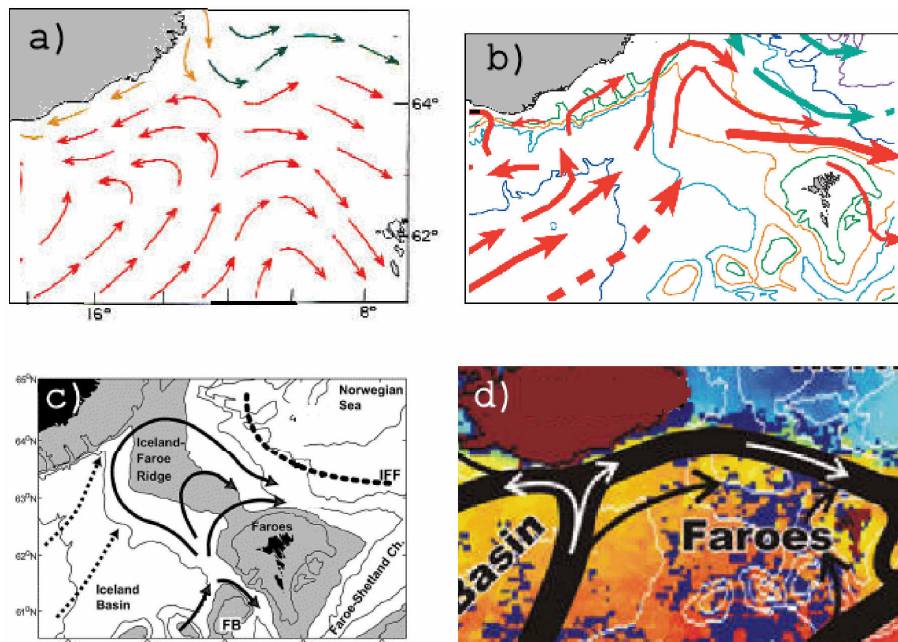
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**Fig. 15.** Setup and results of the NIIC/SIC forcing experiment. **(a)** Topography and prescribed stationary density field, **(b)** stationary sea surface elevation after the spin-up, **(c)** stationary flow field at the depth of 45 m, **(d)** stationary flow field at the depth of 2500 m.

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**Fig. 16.** Different interpretations of the Atlantic Water flow (red or black arrows) between Iceland and the Faroes: **(a)** from Stefánsson and Ólafsson (1991), **(b)** from Valdimarsson and Malmberg (1999) based on drifter data, **(c)** the classical view of Atlantic Water pathways (unbroken arrows) and “alternative suggestions” (broken arrows) from Hansen et al. (2003), **(d)** the Atlantic Water pathways suggested by Orvik and Niiler (2002) based on drifter data. Modified after Stefánsson and Ólafsson (1991), Valdimarsson and Malmberg (1999), Orvik and Niiler (2002), and Hansen et al. (2003).

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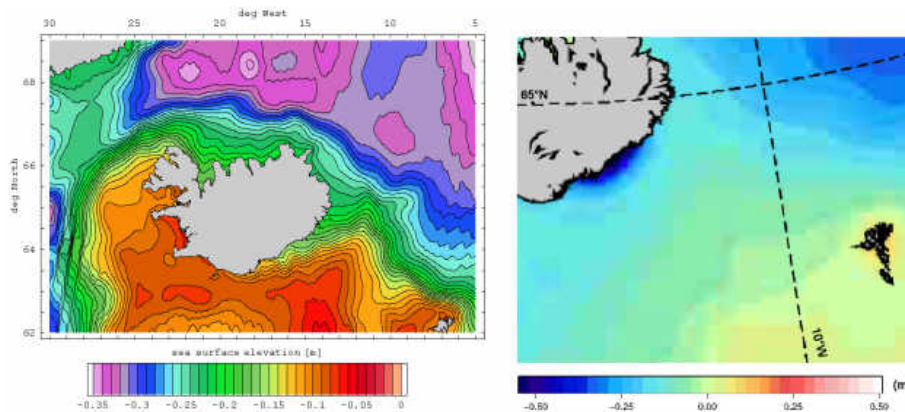
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**Fig. 17.** Left: the simulated 1992–2006 mean sea surface elevation around Iceland. Right: the mean dynamic topography model of Hunegnaw et al. (2009) calculated from marine, airborne and satellite gravimetry, combined with satellite altimetry. Modified after Hunegnaw et al. (2009).

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